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Quantifying the Known Unknown: Including Marine Sources of Greenhouse Gases in Climate Modeling

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ABSTRACT

Researchers have recently estimated that Arctic submarine permafrost currently traps 60 billion tons of methane and contains 560 billion tons of organic carbon in seafloor sediments and soil, a giant pool of carbon with potentially large feedbacks on the climate system. Unlike terrestrial permafrost, the submarine permafrost system has remained a “known unknown” because of the difficulty in acquiring samples and measurements. Consequently, this potentially large carbon stock never yet considered in global climate models or policy discussions, represents a real wildcard in our understanding of Earth’s climate. This report summarizes our group’s effort at developing a numerical modeling framework designed to produce a first-of-its-kind estimate of Arctic methane gas releases from the marine sediments to the water column, and potentially to the atmosphere, where positive climate feedback may occur. Newly developed modeling capability supported by the Laboratory Directed Research and Development (LDRD) program at Sandia National Laboratories now gives us the ability to probabilistically map gas distribution and quantity in the seabed by using a hybrid approach of geospatial machine learning, and predictive numerical thermodynamic ensemble modeling. The novelty in this approach is its ability to produce maps of useful data in regions that are only sparsely sampled, a common challenge in the Arctic, and a major obstacle to progress in the past. By applying this model to the circum-Arctic continental shelves and integrating the flux of free gas from *in situ* methanogenesis and dissociating gas hydrates from the sediment column under climate forcing, we can provide the most reliable estimate of a spatially and temporally varying source term for greenhouse gas flux that can be used by global oceanographic circulation and Earth system models (such as DOE’s E3SM). The result will allow us to finally tackle the wildcard of the submarine permafrost carbon system, and better inform us about the severity of future national security threats that sustained climate change poses.

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ACRONYMS AND TERMS

Acronym/Term	Definition
AOM	anaerobic oxidation of methane
BP	before present
DOE	Department of Energy
ESAS	East Siberian Arctic Shelf
E ³ SM	Energy Exascale Earth System Model
IGHCC2	2 nd International Gas Hydrate Code Comparison Study
GHG	greenhouse gas
kA	kilo annum (thousand years)
KNN	k-nearest neighbors
LDRD	Laboratory Directed Research and Development
LGM	Last Glacial Maximum
MAAT	mean annual air temperature
MIS	Marine Isotope Stage
RSL	relative sea level
TOC	total organic content

1. MOTIVATION

Researchers have recently estimated that Arctic submarine permafrost currently traps 60 billion tons of methane and contains 560 billion tons of organic carbon in seafloor sediments and soil [1], a giant pool of carbon with potentially large feedbacks on the climate system. For comparison, humans have released a total of ~500 billion tons of carbon into the atmosphere since the Industrial Revolution. Unlike terrestrial permafrost, the submarine permafrost system has remained a “known unknown” because of the difficulty in acquiring samples and measurements. Consequently, this potentially large carbon stock never yet considered in global climate models or policy discussions, represents a real wildcard in our understanding of Earth’s climate.

“Arctic methane is underestimated in current carbon budgets, increasing its already significant global climate impact. Arctic methane should play an important role in global climate policy decisions and actions. However, the Global Methane Pledge does not fully address permafrost sources of methane release and warming. Permafrost methane has significant potential to drive accelerated warming. We currently have no integrated, strategic approach to track and predict carbon emissions from thawing permafrost. Major scientific and technical gaps in measurement, monitoring and modeling capabilities produce large uncertainties in current and future trajectories of Arctic methane emissions. Current Earth system models have very limited capability for modeling permafrost thaw and associated methane emissions. Important recommendations emerging from this Arctic methane workshop include: Acceleration of collaboration among scientists and policymakers on the needed upgrades in spatial and temporal coverage of various monitoring and modeling assets; Rigorous assessment requiring improved modeling, both in terms of fundamental methane physical/chemical/biological processes and at multiple spatial and temporal resolutions.”

- *Arctic Methane — Situational Awareness, Assessment & Policy Directions* Workshop hosted by the Wilson Center Polar Institute and Sandia National Laboratories, June 23, 2022.

This report summarizes our group’s effort at developing a numerical modeling framework designed to produce a first-of-its-kind estimate of Arctic methane gas releases from the marine sediments to the water column, and potentially to the atmosphere, where positive climate feedback may occur. Newly developed modeling capability supported by the Laboratory Directed Research and Development (LDRD) program at Sandia National Laboratories now gives us the ability to probabilistically map gas distribution and quantity in the seabed by using a hybrid approach of geospatial machine learning, and predictive numerical thermodynamic ensemble modeling [2]. The novelty in this approach is its ability to produce maps of useful data in regions that are only sparsely sampled, a common challenge in the Arctic, and a major obstacle to progress in the past. By applying this model to the circum-Arctic continental shelves and integrating the flux of free gas from *in situ* methanogenesis and dissociating gas hydrates from the sediment column under climate forcing, we can provide the most reliable estimate of a spatially and temporally varying source term for greenhouse gas (GHG) flux that can be used by global oceanographic circulation and Earth system models (such as DOE’s E³SM). The result will allow us to finally tackle the wildcard of the submarine permafrost carbon system, and better inform us about the severity of future national security threats that sustained climate change poses.

1.1. Greenhouse Gases Sequestered in Arctic Marine Sediments

As summarized in [1], during the last glacial period (~115,000–11,700 years BP), permafrost formed on exposed portions of the continental shelves surrounding the Arctic Ocean [3],[4]. Unglaciaded portions of the exposed continental shelves accumulated billions of tons of undecomposed plant material in frozen sediment [5],[6],[7]. Methane (CH₄) from biogenic and thermogenic sources accumulated within and below permafrost deposits, with some in gas hydrate form [8][9],[10]. After the Last Glacial Maximum (LGM, ~26,500 years BP), climate warming melted ice sheets and glaciers, which increased global sea level by ~134 m on average [5][11], inundating more than 3 million km² of permafrost [7][12], roughly 1/5th the size of the present-day terrestrial permafrost region [13],[14] and containing similar density in soil organic carbon [1]. This marine transgression changed the thermal conditions of inundated permafrost, initiating warming and thawing of submarine permafrost that continues today [15],[16][17]. Like terrestrial permafrost, warming of the submarine permafrost makes the carbon stocks available to methanotrophic communities in the marine sediments which release CH₄ as a metabolic product of organic matter degradation. Warming of the sediments may also destabilize relict methane hydrate deposits which dissociate into water and methane gas. Because neither the amount nor the climate sensitivity of subsea carbon stocks is truly known, the submarine permafrost domain remains one of the least-constrained ecosystem feedbacks in the Earth's climate system [18][14][9].

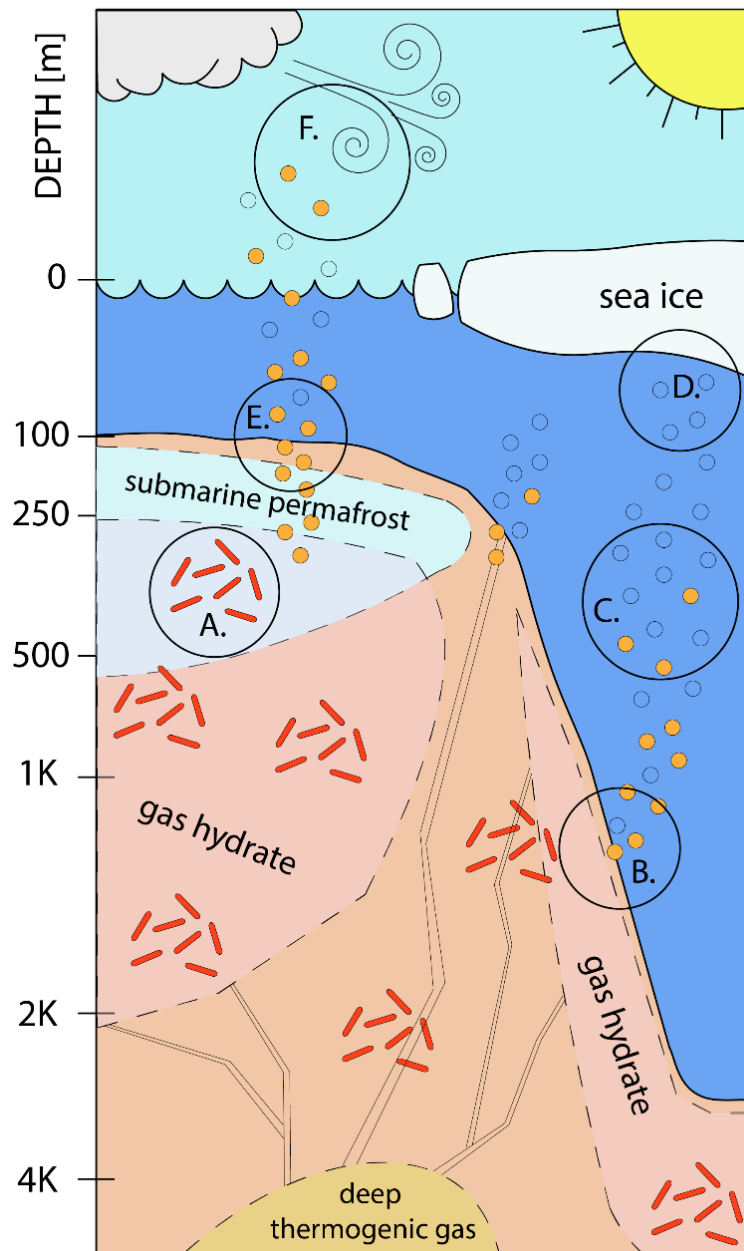
1.2. Gas Escape at the Sediment-Water Interface

Not all gas produced in the marine sediments from methanogenesis or gas hydrate dissociation will make its escape into the water column. Geochemical [19][20] and microbiological evidence [21], together with mass balance calculations [22], nonetheless suggest that up to 90-100% of the dissolved methane produced in marine sediments is consumed *in situ* before reaching the water column by a process known as the anaerobic oxidation of methane (AOM) [23]. AOM is mediated by a consortium of methane-oxidizing Archaea and sulfate-reducing bacteria [21] and requires an alternate electron acceptor other than oxygen. In marine systems this electron acceptor is usually sulphate which comes from the overlying water column [22]. The reaction produces sulphide and the dissolved methane is oxidized to CO₂, which contributes to the formation of carbonates in the sediments near methane seeps [24].

The efficiency of the AOM community may vary under certain conditions, such as in active sediments where upward fluid fluxes are very high, or in areas with high sedimentation rates [27]. Such scenarios may occur after the onset of sudden submarine permafrost thaw, sudden methane hydrate dissociation, or near river mouths. However, for the majority of shelf sediments, as exemplified in a study at the Laptev Sea, [25] show the seawater intrusion and supply of sulphate is sufficient to keep pace with the release of dissolved methane from the thawing submarine permafrost below, fueling an efficient microbial biofilter that prevents much of the dissolved methane from reaching the water column.

On the other hand, methane is much more likely to pass through this biofilter entirely if it is rapidly advected towards the surface in the form of a bubble [26], also known as ebullition. Because methane in gaseous form is not directly accessible to the AOM community [27], it is assumed that the entire amount of methane transported within the gas phase can reach the water column. Such bubbling gas seeps are not uncommon, although quantifying their distribution in space and time is still the focus of ongoing research [28].

The Fate of Methane Produced in the Marine Environment

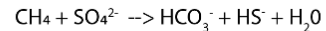


Original figure by J. M. Frederick.

A. Sources of methane

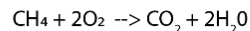
Organic matter in the sediments is consumed by methanotrophs, producing methane that is sequestered in gas hydrate. Deep thermogenic gas can also feed hydrate reservoirs.

B. Anaerobic oxidation of methane (AOM) in sediments



About 90% of dissolved methane flux is taken up by AOM, an effective biofilter. However, rapid advection as bubbles can bypass this filter.

C. Aerobic oxidation of methane in water column



As bubbles rise in the water column, methane is lost as it dissolves into the water and undergoes oxidation.

D. Seasonal gas emission due to sea ice cover

Sea ice cover reduces any gas exchange by an order of magnitude, creating seasonal patterns in gas emissions.

E. Vigorous gas seep in shallow water

Ebullition of methane bubbles in shallow water have the best chance of bypassing the AOM biofilter and reaching the atmosphere.

F. Atmospheric transport and chemical reactions

Methane that reaches the atmosphere is transported and reacts, contributing to the greenhouse gas effect.

Figure 1. A visual summary of the fate of methane produced in the marine environment. Methane sources and reservoirs in the sediment include: deep thermogenic reservoirs, in-situ methanogenesis of organic matter, and gas hydrate reservoirs. Gas can migrate from the sediments to the water column by diffusion, or ebullition (bubbling). Anaerobic oxidation of methane, mediated by sulfate-reducing bacteria, acts as an effective biofilter for dissolved gas, but vigorous advective transport (ebullition) can bypass this filter and enter the water column. As bubbles rise in the water column, methane is lost to diffusion, and gets aerobically oxidated. Only bubbles released from less than ~100m depth have a chance of surviving to the atmosphere with some methane left. Once in the atmosphere, transport and chemical reactions occur. Sea ice coverage can reduce gas exchange by more than an order of magnitude, creating seasonality to methane emissions. The greenhouse gas effect can lead to a positive climate feedback, potentially accelerating methane emissions.

1.3. Gas Transport in the Ocean Water Column

Methane released from marine sediments can reach the ocean surface or surface mixed layer via turbulent diffusion or rising bubbles. While turbulent transport dominates in deep systems, the release of methane bubbles (ebullition) is the most important pathway in shallow waters, such as the pan-Arctic continental shelf (average water depth ~ 50 m), where they are most likely to reach the atmosphere [29],[30]. As previously described, ebullition is the most effective way for gas to pass through the sediment biofilter as well.

Once in the water column, the larger the bubble, or the shallower the release depth, the more CH_4 will survive the journey to the ocean surface. This is because as the bubble rises, dissolution (methane gas escaping the bubble) and exsolution (other gas species present in the water invading the bubble) will eventually change the composition of gas species inside the bubble [31]. For small bubbles or deep release depths, it is entirely possible for a bubble (consisting entirely of methane when released) to be composed of nothing but “air” when it finally reaches the water surface (air being predominantly nitrogen, oxygen, and carbon dioxide). In a seminal study, [32] have quantified methane gas bubble dissolution using a combination of bubble modeling and acoustic observations of rising bubbles to determine what fraction of the methane transported by bubbles will reach the atmosphere as a function of bubble release depth and bubble diameter. This study can provide bounds on the ebullitive flux of gas required for CH_4 to reach the atmosphere. For example, they found (for a 50 m release depth) that bubbles larger than 4 mm in diameter will reach the atmosphere with some CH_4 left [32]. Simple yet highly predictive 1D models of free and dissolved gas transport in the shallow water column toward the atmosphere have recently been developed [31] and can also be adopted to all of the pan-Arctic continental shelves to provide estimates for how much methane released at the seafloor can survive the journey to the ocean surface.

At bubble seep sites, bottom water is often highly concentrated in methane relative to the surroundings. It has been traditionally understood that dissolved methane in these bottom waters does not influence the atmosphere. For example, in a study conducted west of Svalbard in the summer of 2014, [33] observed high concentrations of dissolved CH_4 in the ocean above the seafloor with a sharp decrease above the pycnocline and very little CH_4 in the atmosphere directly above the seep site. They suggest that dissolved methane captured below the pycnocline may only be released to the atmosphere when physical processes remove this dynamic barrier. Storms that cause large mixing events are one mechanism which may facilitate transport of methane-bearing bottom waters to the surface, resulting in episodic but potentially large methane emissions to the atmosphere [34].

1.4. Gas Emission from the Ocean Surface to the Atmosphere and The Ongoing Scientific Debate

Although the Arctic Ocean constitutes just 4.3% of the global ocean surface area, according to [1] the conservative ocean-atmosphere flux of CH_4 from the submarine permafrost domain is $\sim 15\%$ of CH_4 emissions from global ocean estimates [35]. This flux is dominated by shallow environments, where CH_4 released from the seafloor can more efficiently escape to the atmosphere before oxidation. Moreover, fluxes from submarine permafrost represent 10%–40% of CH_4 release from the five-fold larger terrestrial permafrost zone [36]. Clearly, if these estimates are reliable, methane is being released from the thawing Arctic Ocean seabed at disproportionately larger quantities than from terrestrial permafrost (albeit the proportion that reaches the atmosphere is certainly lower but still highly debated).

Estimates of the annual methane emissions from the ocean to the atmosphere vary widely, with the most hotly contested area being the East Siberian Arctic Shelf (ESAS). The largest shelf in the world and comprising ~25% of all the Arctic shelves, methane emissions have been reported as high as 17 Tg yr⁻¹ from the ESAS alone, attributed to degrading submarine permafrost and gas hydrates [37][38]. A large community continues to question the magnitudes of methane emissions previously reported. For example, based on a comprehensive statistical analysis of atmospheric observations and simulations, annual methane emissions from the ESAS were estimated to range from 0.0 to 4.5 Tg yr⁻¹ [39]. The same study reported isotopic observations that suggest a biogenic origin (either terrestrial or marine) of the methane in air masses originating from the ESAS. However, a more recent study based on triple-isotope fingerprinting suggests, on the contrary, that ESAS methane may not primarily originate directly from the submarine permafrost (which would have a biogenic origin) but rather the continuous leakage of an old geological reservoir to the water column through degraded regions of the submarine permafrost, serving as conduits of deeper methane to gas-charged shallow sediments [40]. These are just a few examples of the scientific controversy surrounding seabed methane emissions in the ESAS.

The debate is ongoing among the rest of the Arctic too, as nicely summarized by [27], in terms of the significance of submarine permafrost thaw and hydrate destabilization for GHG emissions from Arctic sediments [41],[10], the methane concentrations in Arctic Ocean waters [42][44] and, ultimately, the methane emissions from the Arctic waters [10][43]. Yet methane emissions from warming Arctic shelves remain poorly quantified today [45],[46] and the truth is still far from being known.

Recent satellite observations provide a different perspective in our understanding of GHG emissions from the Arctic Ocean. Using three thermal infrared space-borne spectrometers, [47] estimates the annual CH₄ emission from the Arctic seas is roughly 18 Tg yr⁻¹, with the ESAS contributing 1.3–3.5 Tg yr⁻¹. Because this study was based on remote sensing data, available year-round, it was able to highlight the seasonality associated with methane emissions. [47] found that ocean-atmosphere exchange was correlated to deep mixing events in the Arctic seas, which tended to happen in winter (a time field work is rarely done) and depended on sea ice coverage. Sea ice cover effectively puts a lid on most gas exchange, reducing fluxes by at least an order of magnitude [48], and can cause build-up of gas beneath the ice, which escapes only after break-up, or through large fractures such as polynyas. Similar patterns in modeled methane emissions have been reported by [49], who found that spatial variability of methane fluxes into the atmosphere were primarily due to the peculiarities of the water circulation and ice conditions. Furthermore, increasing periods and areas of ice-free water and decreasing ice concentration have contributed to a steady increase in methane emissions since the middle of the first decade of the current century

1.5. Global Climate Modeling

Global climate models allow us to better understand how the Earth's climate will evolve as an interconnected, complex system. Models such as the DOE's E³SM, an advanced Earth system model, do this through the coupled simulation of land, ocean, and atmospheric physics, and require the world's most powerful supercomputers to run. While E³SM does not currently have the capability to accept spatially varying methane sources, researchers are currently working on several critical improvements for the next phase (V3). These include the capability to: (i) simulate methane based on surface emissions and climate feedbacks, and (ii) simulate methane gas in the stratosphere, which is critical to simulating the climate impact of its evolving emissions. Once this capability is in place, the framework developed in this LDRD study will be well positioned to deliver a methane source term from the Arctic Ocean to the atmosphere that varies spatially and temporally.

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2. A HYBRID MODELING FRAMEWORK FOR ESTIMATING MARINE SOURCES OF GREENHOUSE GASES FROM SHALLOW ARCTIC SHELVES

Newly developed modeling capability supported by the LDRD program now gives us the ability to probabilistically map gas distribution and quantity in the seabed by using a hybrid approach of geospatial machine learning, and predictive numerical thermodynamic ensemble modeling [2]. The novelty in this approach is its ability to produce maps of useful data in regions that are only sparsely sampled, a common challenge in the Arctic, and a major obstacle to progress in the past. By applying this model to the circum-Arctic continental shelves and integrating the flux of free gas and dissociating gas hydrates from the sediment column under climate forcing, we can provide the most reliable estimate of a time-varying source term for GHG flux that can be used by global oceanographic circulation, chemistry-climate, and Earth system models (such as DOE's E³SM). This section describes the components of the workflow.

Our workflow consists of 3 main software packages: (1) the Global Predictive Seabed Model which provides us with machine-learned seabed parameters used in geophysical modeling, (2) Dakota which statistically samples on the seabed parameters provided by GPSM in order to create and drive ensemble simulations in the geophysical model, and (3) the geophysical numerical model, PFLOTTRAN, which is a powerful subsurface flow and reactive transport simulator for porous media.

2.1. The Global Predictive Seabed Model (GPSM)

The objective of seafloor prediction is to obtain optimal estimates of any given quantity where it has not been directly measured. Because the seafloor is vast and very sparsely sampled in most quantities, we require a prediction almost everywhere. Conventional interpolation techniques (e.g. splines in tension), that rely only on geospatial proximity, ignore almost everything we know about the seabed are ill-suited for our purpose.

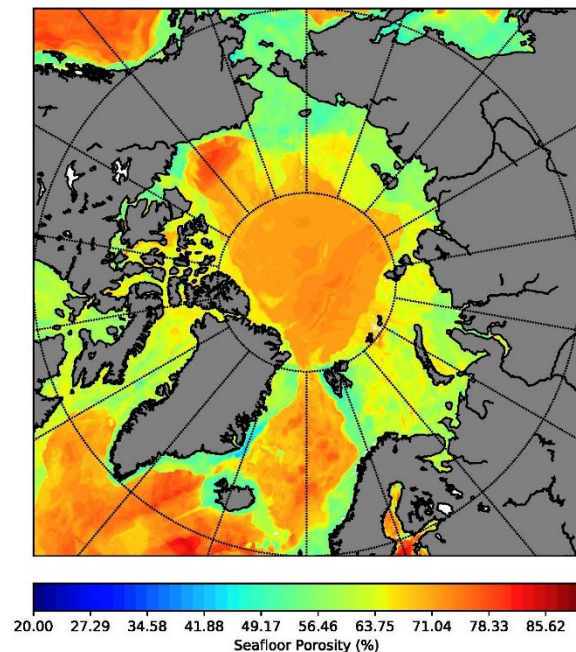


Figure 2. Seafloor porosity predicted by GPSM using geospatial machine learning. (Martin et al., 2015)

Geospatial Machine Learning (GML) predictions are based on the proximity in multi-dimensional geologic predictor space, rather than just geospatial proximity. With GML we predict, or interpolate intelligently, taking advantage of areas that may be geographically very distant, but geologically very similar. GML can use all the information we have about the seabed to predict what we do not know. Not relying strictly on geographical proximity allows prediction in geographical areas that are denied, as well as simply difficult-to-access, such as the Arctic.

GML uses predictors and a predictand. The quantity we wish to predict (the predictand) must also be observed at multiple locations (these observations are the sparse data). Increasing numbers of observations generally correspond to a denser sampling of predictor space, and more certain predictions.

The geologic predictors used in GML are quantities that are known both at the geospatial locations of the

observations, as well as the locations of sites at which we want to predict. These can be quantities that have been measured (e.g. bathymetry), calculated (e.g. distance from shore), or predicted from other GML applications (e.g. porosity). Spatial statistics of these parameters may also be used as predictors. One must clearly be careful to avoid circularity; e.g. using predicted porosity as a predictor to predict a quantity, then subsequently using that same quantity to upgrade the prediction of porosity.

The relationship between the predictors and the predictand need not be known a priori – this is a strength of the GML approach, but also means that for any given application we do not know which predictors will be the most useful. Bad predictors (those poorly correlated with the predictand), can degrade the prediction. For this reason, we typically employ several strategies of predictor (also called feature) selection – determining the best predictors for any given seafloor quantity.

Quantitatively determining the best predictors, as well as quantifying the accuracy of the final prediction, is done through validation. To validate any given group of predictors, we withhold a random fraction of the observed data (usually 10%, for 10-fold validation), make predictions at those locations, and then compare the predicted with the observed values. This is repeated with a different random fraction, until each observation has, at some point, been withheld and predicted. The error between observed and predicted is a reliable, quantitative measure of predictive skill of the prediction method. Validation proves quantitatively and decisively the superiority of machine learning interpolation to conventional, geospatial-only interpolation.

The Naval Research Laboratory's Global Predictive Seafloor Model (GPSM) is a practical implementation of GML designed to provide near real-time estimates of Navy-relevant quantities from continuous seafloor property fields generated by machine learning algorithms (K-Nearest Neighbor, Random Forests, etc.) given often sparse measurements or direct observations compiled from widely available sources.

GPSM has been in use and continuously improved since providing the first global seafloor porosity field, published in 2015 ([50]; see **Figure 2**). More recent publications include predictions of total organic carbon [51], seep probability [28], seafloor sediment accumulation rates [52], recent bathymetry change [53], sediment thickness vs. geologic age [54], density [55], and gas hydrate accumulation [56].

In this work, GPSM is a critical source of model input parameters. Because the seafloor is vast and very sparsely sampled in most quantities, it is impossible to use a measured value everywhere when modeling large regions, such as the Arctic continental shelves. GPSM provides needed model parameters that are spatially continuous and at a high resolution (5 arc minute x 5 arc minute). Because calculated quantities are obtained using machine learning techniques rather than interpolation based on spatial proximity, they provide us with the most reliable values in the absence of observations. Moreover, as a result of GPSM's validation methods, predicted parameter values provided at each seabed location have a standard deviation associated with them (used as a measure of uncertainty from the validation method). The standard deviation is used to construct a distribution for the seabed parameter, allowing us to sample across the distribution and create model ensemble members for geophysical modeling rather than being limited to deterministic simulations. Specifically, we use GPSM-derived seabed values for sedimentation rate (m yr^{-1}), total organic content (TOC) (%), heat flux (W m^{-2}), sediment porosity, seafloor depth (m), seafloor temperature ($^{\circ}\text{C}$), among many other characteristics of interest.

Estimates for TOC and heat flux were made using the k-nearest neighbors (KNN) algorithm [54] based on datasets of TOC [57] and heat flux [58]. Preliminary predictions for sedimentation rate come from taking sedimentary thickness values from GlobSed [59] and dividing by their reported ages [60].

Seafloor temperature estimates are provided through the World Ocean Atlas [61] and depths were taken from SRTM15+ down-sampled to 5 minutes [62]. Porosity profiles were determined by the Random Forest method using data from the Deep Sea Drilling Project, Ocean Drilling Program, Integrated Ocean Drilling Program, and Pangaea data sets [63],[64]. Further detail on how GPSM is used in our simulation framework will be given in the following sections.

2.2. Dakota

We use the statistical software package, Dakota, to sample on the distribution associated with each uncertain seabed parameter at each simulation location within the computational domain used in geophysical numerical modeling. Dakota was developed to optimize simulation codes and provide uncertainty quantification of parameter studies [65] and has a wealth of capability that we do not use in this study but can be included for additional analysis of simulation results (see [2] for examples, such as determining correlation coefficients between input variables and output values). Values for the mean, μ , and standard deviation, σ , of each seafloor parameter from GPSM are written as input to Dakota to sample on for each simulation location. Dakota uses the Latin hypercube sampling method to produce a set of values for any uncertain variable used as input in the simulations by dividing the range of each variable into N distributions of probability $1/N$ and then selecting N values for each variable in order to accurately represent the variability of all parameters [66]. In this work, we assume the distribution is normal for all uncertain parameters and allow for the truncation of this distribution to prevent unphysical values from being chosen (i.e., negative values for parameters that can only be positive, etc.). Dakota is also used as a driver for the ensemble geophysical simulations it has partially defined through the statistical sampling method used.

2.3. PFLOTTRAN HYDRATE Mode

The open source, massively parallel, multiphase flow and reactive transport finite volume simulator PFLOTTRAN [105] is used as the geophysical numerical model in this workflow and was chosen both for its high-performance computing capabilities and its extensive in-house user and developer community at Sandia National Laboratories. Previous support by the LDRD program at Sandia National Laboratories allowed the development of gas hydrate modeling capability to model methane production along with associated free gas and/or gas hydrate formation in marine sediments (PFLOTTRAN's HYDRATE Mode) [2]. This capability was verified through an international code comparison effort where results very closely matched those of simulators like TOUGH+HYDRATE and STOMP [74]. Moreover, it was applied to a well-studied gas hydrate province offshore North Carolina called Blake Ridge [56], and the results from this study showed strong agreement with previous modeling studies at this site and with data from Ocean Drilling Program expeditions [104].

Under the scope of the current work, further development was undertaken to fully treat the evolution of multiphase free gas, gas hydrate, ice, and liquid water systems in marine sediments under significant thermal and pressure perturbations and over geologic length- and time- scales. Specifically, the work presented here will discuss how ice fits into the governing equations solved by HYDRATE Mode, the 15 “phase states” of the simulator in HYDRATE Mode, capillary pressure function non-uniqueness when considering the presence of multiple non-wetting phases like ice, and an overview of all options available to a user of HYDRATE Mode in PFLOTTRAN.

2.3.1. Governing Equations

PFLOTTRAN's HYDRATE Mode solves a nonlinear system of partial differential equations describing transient conservation of water mass, component mass (e.g., methane), and energy. The system of

equations is tightly coupled and solved fully-implicitly using a two-point flux finite volume spatial discretization and backward Euler time discretization, ensuring unconditional stability of the solution. The system of equations is solved using a Newton-Raphson solution search technique. The code is fully parallelized and can run on structured or unstructured grids. The governing equations for mass and energy conservation are:

$$\frac{\partial}{\partial t} \phi (\sum_{\alpha=l,g,h,i} s_{\alpha} \rho_{\alpha} x_{w}^{\alpha}) + \nabla \cdot (q_l \rho_l x_w^l + q_g \rho_g x_w^g - \phi s_l D_l \rho_l \nabla x_w^l - \phi s_g D_g \rho_g \nabla x_w^g) = Q_w \quad \text{Equation 1}$$

$$\frac{\partial}{\partial t} \phi (\sum_{\alpha=l,g,h,i} s_{\alpha} \rho_{\alpha} x_m^{\alpha}) + \nabla \cdot (q_l \rho_l x_m^l + q_g \rho_g x_m^g - \phi s_l D_l \rho_l \nabla x_m^l - \phi s_g D_g \rho_g \nabla x_m^g) = Q_m \quad \text{Equation 2}$$

$$\sum_{\alpha=l,g,h,i} \left(\frac{\partial}{\partial t} (\phi s_{\alpha} \rho_{\alpha} U_{\alpha}) + \nabla \cdot (q_{\alpha} \rho_{\alpha} H_{\alpha}) \right) + \frac{\partial}{\partial t} ((1 - \phi) \rho_r C_p T) - \nabla \cdot (\kappa \nabla T) = Q \quad \text{Equation 3}$$

where ϕ is the porosity of the medium, s_{α} is the saturation of phase α , ρ_{α} is the density of phase α , x_w^{α} is the mass fraction of water in phase α , q_l and q_g are the respective advective liquid and gas fluxes, D_l and D_g are the respective liquid and gas diffusion coefficients, Q_w and Q_m are respective source/sink terms for water and the solute component, U_{α} is the specific internal energy of phase α , H_{α} is the specific enthalpy of phase α , C_p is the heat capacity of the rock, T is temperature, and κ is the composite thermal conductivity of all phases. $-\partial \phi / \partial t (\sum_{\alpha=l,g,h,i} s_{\alpha} \rho_{\alpha} x_w^{\alpha}) + \nabla \cdot (q_l \rho_l x_w^l + q_g \rho_g x_w^g - \phi s_l D_l \rho_l \nabla x_w^l - \phi s_g D_g \rho_g \nabla x_w^g) = Q_w$ Equation 1 describes conservation of the water mass component (w) and $-\partial \phi / \partial t (\sum_{\alpha=l,g,h,i} s_{\alpha} \rho_{\alpha} x_m^{\alpha}) + \nabla \cdot (q_l \rho_l x_m^l + q_g \rho_g x_m^g - \phi s_l D_l \rho_l \nabla x_m^l - \phi s_g D_g \rho_g \nabla x_m^g) = Q_m$ Equation 2 describes conservation of the solute component (m), which for the purpose of constitutive relationship development in this work is considered to be methane but which could be developed more generally in the future. $\sum_{\alpha=l,g,h,i} \left(\frac{\partial}{\partial t} (\phi s_{\alpha} \rho_{\alpha} U_{\alpha}) + \nabla \cdot (q_{\alpha} \rho_{\alpha} H_{\alpha}) \right) + \frac{\partial}{\partial t} ((1 - \phi) \rho_r C_p T) - \nabla \cdot (\kappa \nabla T) = Q$ Equation 3 describes conservation of energy across all phases, including the rock matrix and considering heat conduction through the composite material (rock matrix and pore constituents).

In this formulation, a given simulated porous medium can contain any combination of liquid (l), free gas (g), gas hydrate (h), and ice (i) phases. The gas hydrate and ice phases are considered immobile and can only move through the domain if the sedimentation feature is turned on, which is described in a later section. Therefore, the governing equations only contain advective flux terms for the liquid and gas phases. Similarly, diffusion is only modeled through the fluid phases and is neglected in the solids.

Advective fluxes are modeled using Darcy's Equation as follows:

$$q_{\alpha} = -\frac{k k_{\alpha}^r}{\mu_{\alpha}} \nabla (P_{\alpha} - \rho_{\alpha} g z) \quad \text{Equation 4}$$

where k is the intrinsic permeability of the porous medium, k_{α}^r is the relative permeability of mobile phase α , μ_{α} is the phase viscosity, and P_{α} is the phase pressure. In multiphase systems, the nonwetting phase pressure is constrained through a capillary pressure function relationship. HYDRATE Mode can access PFLOTTRAN's standard capillary pressure relationships which include Van Genuchten and Brooks-Corey formulations, among others. Likewise, HYDRATE Mode also can use all of PFLOTTRAN's standard relative permeability functions.

The Fourier heat conduction term uses a composite thermal conductivity term, κ , which includes all phases in the pore space and the solid rock phase. HYDRATE Mode can use any of the standard thermal conductivity options in PFLOTTRAN, including temperature-dependent thermal conductivity, but because HYDRATE Mode models gas hydrate and ice it is recommended that the user use the default formulation for effective thermal conductivity, κ_{eff} :

$$\kappa_{eff} = \kappa_{dry} + \phi \sum_{\alpha=l,g,h,i} S_{\alpha} \kappa_{\alpha} \quad \text{Equation 5}$$

where κ_{dry} is the dry thermal conductivity, an input parameter equivalent to $(1 - \phi)k_{rock}$, $\kappa_l = 0.59$ W/m-K, $\kappa_m = 0.03$ W/m-K, $\kappa_h = 0.58$ W/m-K, and $\kappa_i = 2.2$ W/m-K.

Furthermore, internal energy is related to enthalpy via phase pressure as follows:

$$U_{\alpha} = H_{\alpha} - \frac{P_{\alpha}}{\rho_{\alpha}} \quad \text{Equation 6}$$

The internal energy of the ice phase is represented via the formulation of [67]:

$$\begin{aligned} U_i &= L_w + 185T + 3.445(T^2 - 273.15^2), \quad T \geq 90K \\ U_i &= L_w + 4.475(T^2 - 273.15^2), \quad T < 90K \end{aligned} \quad \text{Equation 7}$$

where L_w is the latent heat of fusion for water, set at -6017.1 J/mol.

The enthalpy of the hydrate phase is given by the following [68]:

$$H_h = C_{ph}(T - T_q) + \frac{H_{h0}}{N_{hyd}+1} \quad \text{Equation 8}$$

where H_{h0} is the hydrate reference enthalpy, set to -54,734 J/mol, N_{hyd} is the hydration number (default is 6), T_q is the quadruple point (0C in the absence of salt), and C_{ph} is the heat capacity of hydrate, computed as follows:

$$C_{ph} = \frac{1620(M_w N_{hyd} + M_m)}{1000} \quad \text{Equation 9}$$

where M_w is the molecular weight of water and M_m is the molecular weight of methane.

To form a free gas phase from dissolved methane in an aqueous solution, methane must accumulate until its concentration exceeds its solubility in the aqueous phase. In the current equilibrium formulation, Henry's Law is used to describe the partitioning between the dissolved aqueous and free gas phases. The Henry's constant for methane is computed as follows [69]:

$$K_h = 1000e^{5.1345 + \frac{7837}{T} - \frac{1.509 \times 10^6}{T^2} + \frac{2.06 \times 10^7}{T^3}} \quad \text{Equation 10}$$

where K_h is the Henry's constant in Pa per dissolved mole fraction of methane. When the concentration of dissolved methane in aqueous solution surpasses its equilibrium concentration at a given pressure, the system transitions from the aqueous state to the two-phase gas-aqueous state. More information on phase state combinations in HYDRATE Mode can be found in the following sections.

Likewise, if enough methane is available in the dissolved phase and the pressure and temperature conditions are favorable for gas hydrate formation, excess methane will form a new methane hydrate phase. Formation of the methane hydrate phase consumes both water and methane in constant proportion, by default set to 0.1429 mol CH₄ per mol H₂O. There is also an associated exothermicity of hydrate formation, which is accounted for in the energy balance. Free gas, gas hydrate, and liquid can coexist at specific pressures and temperatures defined by the phase boundary. The default hydrate phase boundary is the representation from [70], and the user also has the option to use a more accurate (and more complex) Moridis (2003) formulation [71]. For this work, a simplified (and importantly, invertible) version of the Moridis (2003) equation was developed:

$$P_e = e^{0.03349409997T - 8.1938174346}, \quad T < T_q \quad \text{Equation 11}$$

$$P_e = e^{0.1100383278T - 29.1133440975}, \quad T \geq T_q$$

where P_e is the equilibrium pressure in MPa.

An ice phase forms whenever the temperature drops below T_q . If there is no salinity in the system, T_q is set to 0C. For this work, a salinity adjustment was added to shift the freezing temperature both of water and of gas hydrate. For gas hydrate, the equilibrium phase boundary is adjusted by a temperature correction as a function of salinity as follows in [72]:

$$dT = dT_0 \frac{\ln(1-x_m^l)}{\ln(1-x_{m,ref}^l)} \quad \text{Equation 12}$$

where dT_0 is the reference temperature correction, set to -0.37C , and $x_{m,ref}^l$ is the reference salt mole fraction, set to 0.03311.

For ice formation, T_q is adjusted as a function of and is computed as follows [73]:

$$dT = -57.5x_m^l + 112(x_m^l)^2 \quad \text{Equation 13}$$

where x_m^l is the dissolved salt mass fraction.

2.3.2. Representation of Ice and Hydrate Phase Behavior in HYDRATE Mode

When gas hydrate forms out of aqueous solution, its formation involves a density change, exothermic energy release, and consumption of the water and methane components. Therefore, further methane hydrate growth is constrained by several mechanisms, including pressure change, temperature change, and limited availability of the methane component to take from the aqueous phase.

In the case of ice formation, the ice phase only consumes the water component and not the methane component. If methane is dissolved in the pore space, the process of ice formation tends to concentrate methane in the remaining liquid water. As ice formation proceeds, the methane can supersaturate and eventually form a separate free gas or gas hydrate phase. Ice formation is typically only constrained by the resulting exothermic temperature change on freezing, meaning that if heat is sufficiently dissipated through the system over the timescales of interest it is very easy to completely fill the pore space with ice. This results in significant numerical challenges (especially in 1D) as the pore space becomes occluded with a solid and pressure in adjacent grid cells with fluid in them have no means of dissipating excess pressure buildup due to volume changes associated with ice formation. It is also physically unrealistic, as a liquid phase is known to exist in the small pores of sediments well below 0°C . Therefore, in the current framework the option was included to treat both ice and gas hydrate formation in small pores as following a formulation of the Gibbs-Thomson equation that transforms pore radius to nonwetting phase capillary pressure vis-à-vis the Young-Laplace equation and all relevant assumptions as follows:

$$dT = -T_m \frac{P_c}{H_f \rho_\alpha \cdot 1000} \quad \text{Equation 14}$$

where dT is the freezing point depression of either ice or hydrate, T_m is the bulk melting temperature of a given phase (ice or hydrate), P_c is the phase capillary pressure, evaluated as a function of effective phase saturation (described in detail in a later section), H_f is the bulk enthalpy of fusion of ice (6033.54 J/mol) or hydrate (7161 J/mol), and ρ_α is the density of ice (920 kg/m^3) or hydrate (900 kg/m^3). The resulting effect is that for a solid ice or gas hydrate phase to precipitate in progressively smaller pores at a given pressure requires progressively lower temperatures. This ultimately adds an additional constraint to both hydrate and ice formation which can practically prevent full blockage of the pore space by a solid.

2.3.3. Capillary Pressure Function Non-Uniqueness

In the formulation developed here, liquid water is considered the only wetting phase. All other phases (free gas, gas hydrate, ice) are considered nonwetting phases. In order to apply a freezing point depression for ice phase or hydrate phase growth, the phase saturations are plugged into a capillary pressure function to obtain a phase capillary pressure and thus a freezing temperature offset. This is straightforward when transitioning from single phase liquid to two-phase (liquid-ice or liquid-hydrate). It is more complicated, however, when dealing with 3 or 4 phases coexisting in the pore space.

For instance, take the example scenario where dissolved methane exists in sediment pore water at high pressure, low temperature, and close to methane solubility in equilibrium with a gas hydrate phase. If the top of the sediment column were to begin to freeze, ice formation in the pore space would consume water but not methane. This would result in a concentration of the remaining methane in the pore water. As the ice phase grows, dissolved methane concentrations in the remaining pore water increase (methane mass stays constant and liquid water mass decreases). At some point, methane concentration exceeds its solubility, and a hydrate phase can form. But hydrate and ice cannot occupy the same pores at the same time, so using absolute phase saturations to compute their respective capillary pressures will result in underpredicting the extent of subcooling required to grow ice or hydrate. If, say, ice exists at 40% saturation and gas hydrate exists at 30% saturation. In this case, 70% of the pore space is occupied by non-wetting phases, so the wetting phase (water) will only be present in the smallest 30% of pores. If capillary pressure were evaluated independently for ice and for gas hydrate, the capillary pressure function would use an inconsistent value for wetting phase saturation (60% for the ice function and 70% for the gas hydrate function), and both values would be much greater than the actual amount. Thus, the capillary pressure function evaluation must consider all phases present.

If the behavior of the system is to be reversible, the resulting state of the system that includes gas hydrate and ice must be solely a function of phase saturations and not on the state of the system at some prior time. To achieve reversibility, the current approach assumes that the ice and hydrate phases partition the largest pore space in a consistent manner that is independent of simulation history. For simplicity, HYDRATE Mode assumes that all nonwetting phases partition the largest pore space in equal amounts. What this means is that for a 3-phase system with ice, gas hydrate, and water, the water remains in the small pores while the ice and gas hydrate partition the large pore space. If gas hydrate forms in pores occupied by ice, the ice must either melt or form in smaller pores to maintain its phase saturation. The resulting capillary pressure function evaluation proceeds as follows:

$$P_{c\alpha} = P_c(2s_\alpha), \quad s_\alpha \leq s_\beta \quad \text{Equation 15}$$

$$P_{c\alpha} = P_c(s_\alpha + s_\beta), \quad s_\alpha \geq s_\beta$$

where $P_{c\alpha}$ is the capillary pressure of phase α , β denotes the other nonwetting phase in the system, and P_c denotes the function for capillary pressure evaluated at an arbitrary saturation, s .

Expanding this formulation, in the instance where four phases are present in the pore space (water, free gas, gas hydrate, and ice), assuming equal partitioning of the co-occupied pore space by the non-wetting phases, capillary pressure of each phase is computed as follows:

$$P_{c\alpha} = P_c(3s_\alpha), \quad s_\alpha \leq s_\beta \leq s_\gamma$$

Equation 16

$$P_{c\alpha} = P_c(2s_\alpha + s_\beta), \quad s_\beta \leq s_\alpha \leq s_\gamma$$

$$P_{c\alpha} = P_c(2s_\alpha + s_\gamma), \quad s_\gamma \leq s_\alpha \leq s_\beta$$

$$P_{c\alpha} = P_c(2s_\alpha + s_\gamma), \quad s_\gamma \leq s_\alpha \leq s_\beta$$

2.3.4. Phase States and Primary Variables

Table 1. Primary Variables for Each Phase State in Hydrate Mode. P_l: liquid pressure, P_g: gas pressure, P_a: air pressure, RH: relative humidity x_m^a: mole fraction of methane in the aqueous phase, x_m^h: mole fraction of methane in the hydrate phase, x_mⁱ: mole fraction of methane in the ice phase, T: temperature, S_l: liquid saturation, S_g: gas saturation, S_h: hydrate saturation, S_i: ice saturation. In instances where P_g is solved but no gas phase exists, this is treated as the pressure of a gas phase were it to be present in the system (often also equal to the liquid pressure).

Phase State	Primary Variables
Aqueous (Liquid) State	P _l , x _m ^a , T
Gas State	P _g , P _a (RH), T
Hydrate State	P _g , x _m ^h , T
Ice State	P _g , x _m ⁱ , T
Gas-Aqueous State	P _g , S _g , T
Hydrate-Gas State	P _g , S _g , T
Hydrate-Aqueous State	P _g , S _h , T
Hydrate-Ice State	P _g , S _h , T
Gas-Ice State	P _g , S _i , T
Aqueous-Ice State	P _l , x _m ^a , S _i
Hydrate-Gas-Aqueous State	S _l , S _h , T
Hydrate-Aqueous-Ice State	P _g , S _l , S _i
Hydrate-Gas-Ice State	S _l , S _h , T
Gas-Aqueous-Ice State	P _g , S _l , S _i
Hydrate-Gas-Aqueous-Ice State	S _l , S _g , S _i

PFLOTTRAN's HYDRATE Mode uses primary variable switching to evaluate state properties of the system in any given grid cell at a given simulation time. For example, if a given grid cell is in the liquid state, then that cell is fully saturated with liquid water and contains no other phases in the pore space. In this case, HYDRATE Mode solves for liquid pressure, mole fraction of dissolved methane, and

temperature as its primary variables. While the primary variable switching occurs internally when a phase transition occurs (e.g., from the liquid state to liquid-hydrate state), if a user intends to apply an initial or boundary condition to the model, the user must set such a condition with the primary variables that are appropriate for the intended state of the condition. For example, if the user initializes the system at hydrostatic pressure with only dissolved methane in the system, the initial condition of the simulation would be a liquid state condition with primary variables as described above. If the user were to then set a top boundary condition representing exposure to the air, then the boundary condition would be in the gas state and the primary variables would be total gas pressure, air partial pressure (or relative humidity), and temperature (see Table 1).

There are a few phase state combinations for which reading in primary variables by themselves does not immediately identify the phase state. Due to the structure of model initialization in the code, these phase states require the user to identify the intended state in the specification of the flow condition. Those states are Hydrate, Ice, Hydrate-Ice, and Hydrate-Aqueous-Ice.

In total, there are 15 combinations of phases that can be present in HYDRATE Mode. When considering the possibility of ice formation, there are 8 potential phase states in which ice could exist; these phase states were more robustly developed as part of this work. To check initialization for all phase states and internal phase state consistency, a suite of single-cell test problems was designed where each simulation is initialized and run to 1 day with no flux boundary conditions. By taking multiple timesteps, the simulator not only initializes the system at a given state with a given set of primary variables, it also must check whether the state changes given the set of primary variables. For a given material, **Table 2** describes the series of tests conducted to verify flow condition setting and internal consistency.

Table 2. Test Cases for Phase State Initialization and Internal Consistency. In bold are the primary variables used to define the state. Pure Hydrate and pure Ice states currently yield numerical convergence difficulties due to fixed methane and ice component mass fractions.

Phase State	T (C)	P _l (Pa)	P _g (Pa)	S _l	S _g	S _h	S _i	x _m ^a	x _m ^g
Liquid	10	1E6	--	1	0	0	0	1E-4	--
Gas	25	-9.9E7	1E6	0	1	0	0	--	0.998
Hydrate	--	--	--	--	--	--	--	--	--
Ice	--	--	--	--	--	--	--	--	--
Gas-Aq	30	9.682E5	1E6	0.5	0.5	0	0	2.24E-4	0.996
Hydrate-Gas	18	--	1E8	0	0.7	0.3	0	--	0.195
Gas-Ice	-0.05	--	2E5	0	0.8	0	0.2	--	0.997
Aq-Ice	-5.8E-3	1E6	--	0.9	0	0	0.1	1E-5	--
Hyd-Gas-Aq	18	1.948E7	1.949E7	0.8	0.14	0.06	0	5.47E-3	0.999
Hyd-Aq-Ice	-1.40E-2	1E7	--	0.65	0	0.05	0.3	1.18E-3	--

Phase State	T (C)	P _l (Pa)	P _g (Pa)	S _l	S _g	S _h	S _i	x _m ^a	x _m ^g
Hyd-Gas-Ice	-1E-2	--	2.62E6	0	0.65	0.25	0.1	--	0.999
Gas-Aq-Ice	-3.76E-3	1.91E5	2E5	0.8	0.15	0	0.05	8.99E-5	0.997
Hyd-Gas-Aq-Ice	-9.66E-3	2.59E6	2.62E6	0.5	0.1	0.2	0.2	1.18E-3	0.999

2.3.5. PFLOTRAN HYDRATE Mode Test Cases

Updated software capabilities were tested on a series of transient simulation test cases to as concept demonstration and to identify areas for performance improvement. The first test case is a modification of a 1D radial thermal stimulation test case that was part of the 2nd International Gas Hydrate Code Comparison Study (IGHCC2) to include capillarity, a simplified 3-phase equilibrium pressure formulation, and 3-phase effective saturation considerations following the formulation of $P_{c\alpha} = P_c(2s_\alpha)$, $s_\alpha \leq s_\beta$ Equation 15. The second test case tracks ice formation in a vertical water column which uses a new Dirichlet Seepage boundary condition and is designed to compare the effect of different pore size distributions on the freezing behavior of water in soils. The third test case is perhaps the most numerically complex and is an extension of the freezing test case which adds a gas source term and also uses a standard Dirichlet top boundary condition for comparison.

2.3.5.1. IGHCCS2 Thermal Test with Capillary Pressure

This test is a modification of the Benchmark Problem 3, Case 1 presented in [74] in which a 1D cylindrical model simulates the dissociation of gas hydrate by thermal stimulation. Benchmark Problem 3, Case 2, models gas hydrate dissociation by depressurization with similar sediment physical properties. The model is radially symmetric with a log-spaced discretization starting with $\Delta r = 0.02$ m and ending with a max $\Delta r = 12.6$ m for a total domain radius of 1 km. The model starts with a hydrate saturation of 50%, a uniform pressure of 4.6 MPa, uniform temperature of 3°C, and applies a heat source at the center of 150 W. In the initial formulation, the thermal stimulation problem does not consider capillary pressure.

For this test, we incorporated the capillary pressure function from the depressurization benchmark problem as well as a simplified version of the Moridis (2003) phase boundary, Gibbs Thomson effects on hydrate growth due to capillarity, and scaling of the mobile phase saturations to send the relative permeability functions values for phase saturations expressed as fractions of only the mobile pore phases (excluding solid phases like hydrate; this would result in, e.g., an effective water saturation of 1 if the system had only gas hydrate and water in the pore space).

In the original problem, all teams showed numerical oscillations in the hydrate and gas saturation solutions. This is seemingly an unavoidable component to these kinds of problems, which involve solving numerically challenging problems related to coupled thermodynamics and phase behavior. Interestingly, incorporating capillary pressure appears to dampen this effect. It also results in lower gas saturations even though total mass in the system should be constant, which is likely due to gas compressibility effects. This also verifies that the simplified phase boundary produces similar results; small differences in the leading edge of the gas hydrate propagation front could be due to the changes to relative permeability computation and small differences in the phase boundary computation. Both simulations run in comparable times (roughly 1.5 min). Results are shown in **Figure 3**.

2.3.5.2. Freezing with Seepage Boundary Condition, P_c smoothing

In this test problem, ice formation was modeled in a 1D vertical column for 2 different pore systems. The column is 1 km thick with 1-m grid spacing in the vertical dimension. It is initialized at hydrostatic, water-saturated conditions at 20 MPa at the bottom of the domain. The initial temperature everywhere is 2°C, and a small concentration of methane is initialized in the system (1E-04 mol CH₄/mol total). The bottom boundary is a Dirichlet BC set to the same conditions as the initial condition, and the top boundary is set to a gas state with a temperature below freezing. The simulation is run for 10 kyrs.

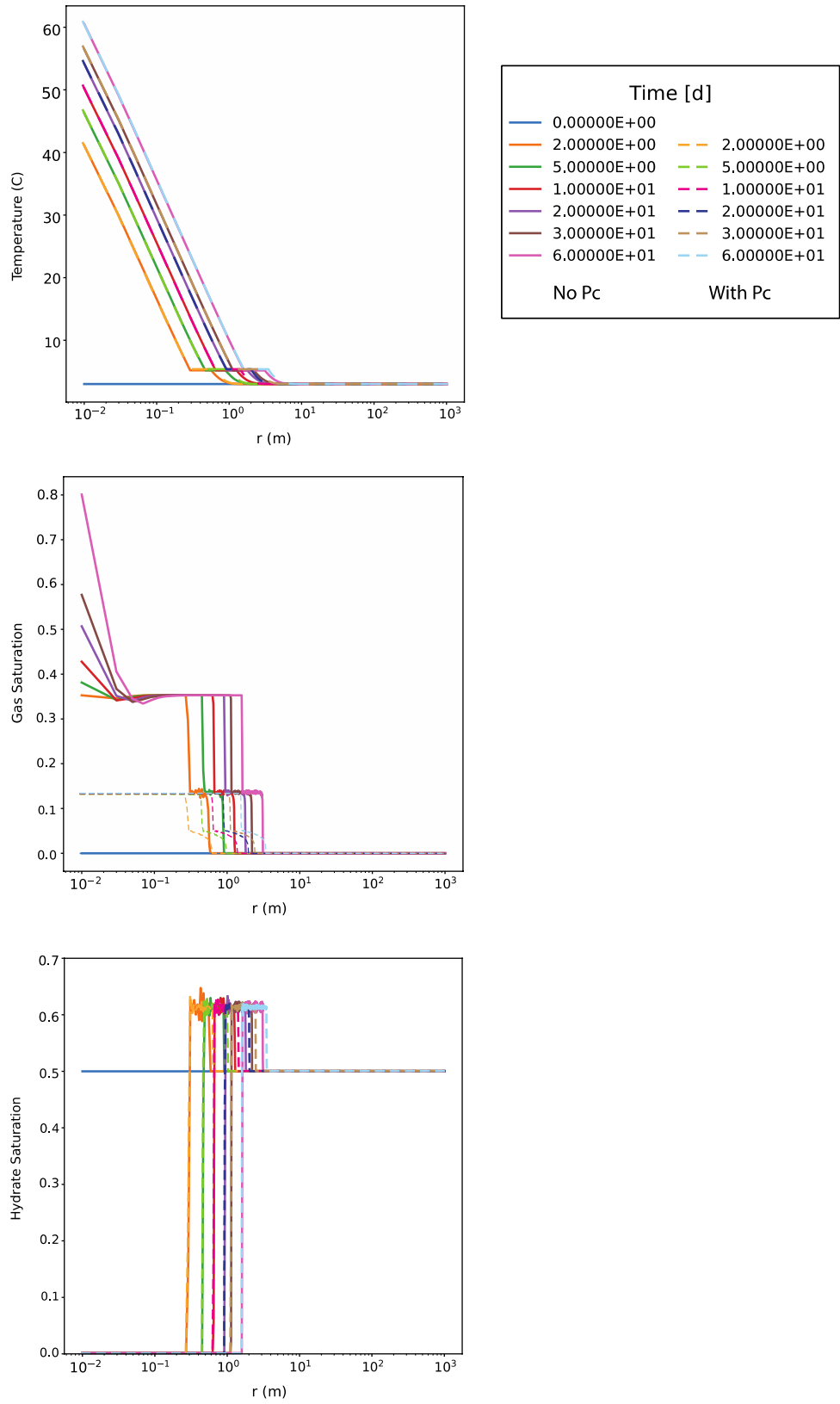


Figure 3. Thermal stimulation benchmark problem including updated HYDRATE mode features in dashed lines (labeled “no Pc”) compared to the original in solid lines (labeled “with Pc”).

This simulation is meant to model exposure of the water-saturated soil to air temperatures below freezing and thus prompt freezing of the soil from top to bottom. The test also demonstrates that freezing is inhibited as ice enters progressively smaller pores, requiring greater and greater subcooling below the bulk freezing temperature of water. The test makes use of a new boundary condition called DIRICHLET_SEEPAGE (modeled after a similar style boundary condition available to RICHARDS mode in PFLOTRAN), in which the simulator only allows flux of gas into the domain if the pressure across the boundary is greater than a BC_REFERENCE_PRESSURE. This is useful for a situation where the modeled gas component, in this case methane, is not the same as the gas that would be represented by a gas phase boundary and thus no significant influx of that component would be expected, such as if the boundary were a standard atmospheric air composition. Finally, a new feature of PFLOTRAN is invoked in this example whereby the unsaturated end of the capillary pressure curve is smoothed (instead of capped), which can allow for better numerical convergence (see **Figure 4**). The capillary pressure functions used here generally represent a pore system with bigger pores (generally lower capillary pressures) and a lower capillary entry pressure (“larger pores” in **Figure 4**) as compared to a system with smaller pores (larger capillary pressures) and a higher capillary entry pressure (“smaller pores” in **Figure 4**). The Van Genuchten parameters used to describe these two pore systems are described in **Table 3**.

Table 3. Van Genuchten parameters for each test for the freezing example.

	Larger Pores	Smaller Pores
Alpha (Pa ⁻¹)	5.8E-04	1E-04
m	0.5	0.2
S _d	0.1	0.1
Max P _c (Pa)	1E+08	1E+08

The influence of the pore network on ice phase formation is apparent in **Figure 5**. In the system with larger pores, ice is less inhibited from forming and it therefore can accumulate in higher saturations at higher temperatures. The system with smaller pores requires a greater degree of subcooling to precipitate ice in the pore space, resulting in overall lower ice saturations. The ice front penetrates deeper at 10 kyrs in the simulation with smaller pores though, and correspondingly the smaller pore simulation sees slightly more subcooling at depth at 10 kyrs than the large pore simulation. This could be because less ice has formed in total in the simulation with smaller pores and therefore there is less heat released due to the phase change, which could allow the temperature pulse to propagate faster. Finally, in both simulations there is an increase in dissolved methane concentration in the liquid pore water upon ice formation because ice formation only consumes the water component and excludes the methane component. The degree to which methane becomes concentrated is a function of the ice saturation.

2.3.5.3. Freezing with Gas Production

In this test, the same domain as described above is simulated with a source term of methane turned on for the first 1,000 years of the simulation and is distributed between 200- and 300-meters sediment depth. The top boundary is a standard Dirichlet boundary condition where gas flux is allowed across the top boundary. This leads to accumulation of dissolved methane in the upper part of the column, some of which is likely erroneously generated by influx from the top. In addition, the initial and

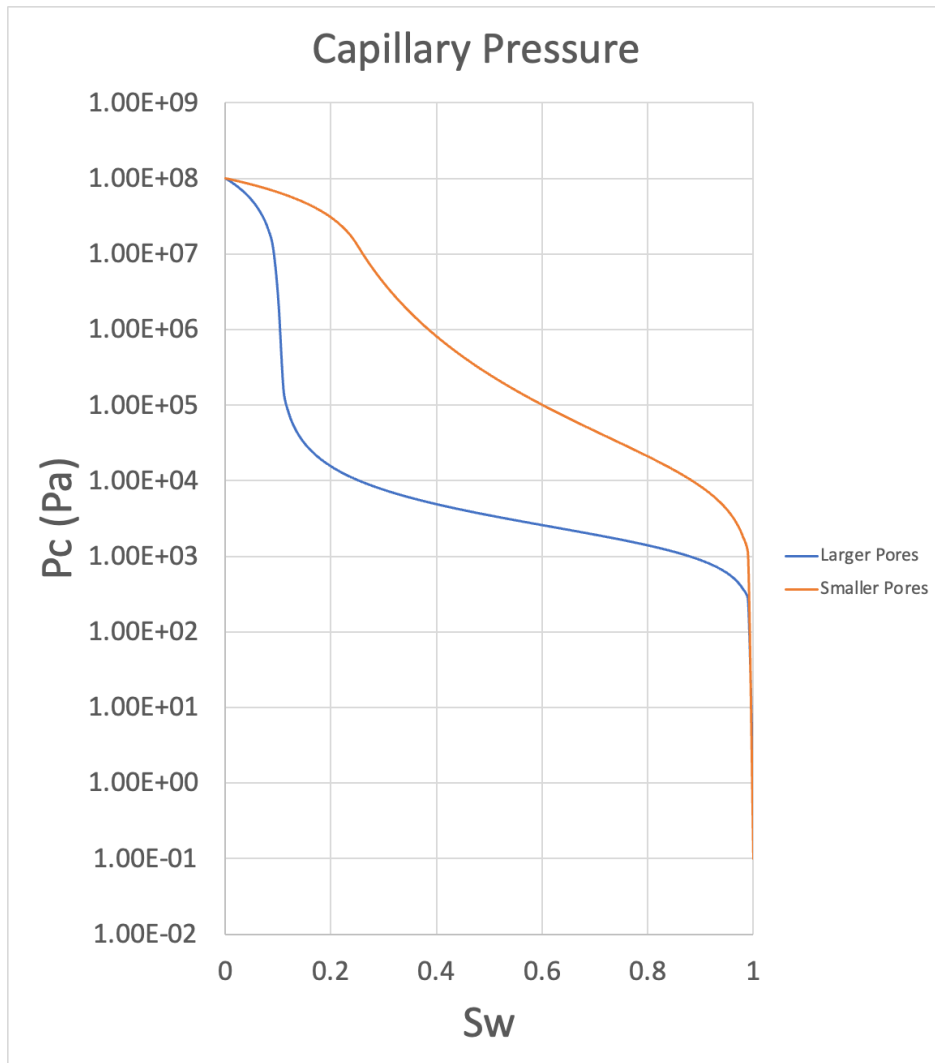


Figure 4. Capillary pressure (a proxy for pore size distribution) as a function of liquid water saturation (wetting phase), for 2 different pore systems.

bottom boundary dissolved methane mole fraction is set to 5E-04 instead of 1E-04 in this simulation in order to speed up hydrate formation. The initial conditions and boundary conditions are otherwise the same as the previous simulation. The sediment capillary pressure function is slightly different and is summarized in **Figure 6** and **Table 1**.

This test is significantly more numerically challenging than the previous example because it involves two solid phases forming in the pore space at the same time. The ice phase forms downward from the top, and the hydrate phase forms in place over the course of the simulation. After about 3,000 years, the ice front meets the hydrate front and **Figure 7** demonstrates the importance of including accurate phase partitioning vis-à-vis $P_{c\alpha} = P_c(2s_\alpha)$, $s_\alpha \leq s_\beta$

Equation 15. At 1,000 years, ice saturations are at around 40% of the pore space as the ice front reaches 200 m depth while hydrate saturations at 200 m exceed 60% saturation. If ice were to keep forming as if no hydrate existed in the pore space at 200 m, the simulator would end up with

over 100% saturation of solid ice + hydrate in the pore space. Not only is greater than 100% saturation physically impossible, but it is also numerically challenging in 1D to have grid

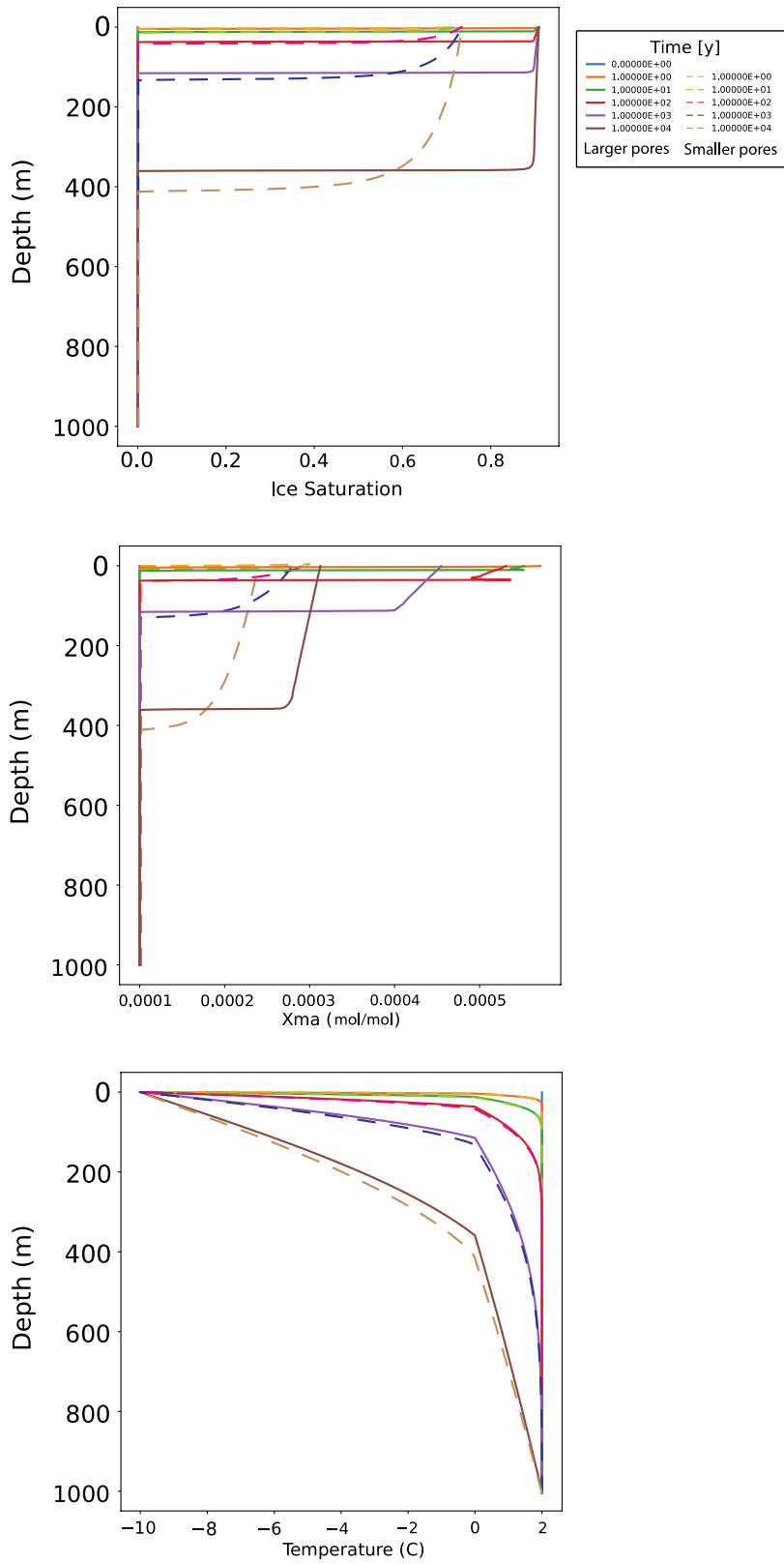


Figure 5. Ice saturation, methane mole fraction, and temperature profiles over time for simulations using 2 different soils.

cells that are completely occupied by solid phases as there is no way to diffuse pressure buildup and the primary variables governing the phase state of the system have less of an influence on the system Jacobian. When considering co-occupancy of the large pore space by both hydrate and ice, there is a noticeable drop in ice saturation in the region already containing hydrate. In this region, the sum of ice and hydrate saturations does not exceed 100%. There are some numerical oscillations apparent in the hydrate saturation profile (and some in the ice saturation profile), which are likely because the combined ice + hydrate saturation is approaching the residual liquid saturation and water is starting to become immobile. This can be problematic because it takes away response of the system to pressure perturbations via invoking fluid flow, which can make the problem significantly more difficult to solve. This suggests that an area for future investigation could be into different mechanisms for scaling the inputs to the relative permeability functions. Also, significant assumptions were made in the model about how the nonwetting phases partition the large pore space and in how both the ice phase and hydrate phase are inhibited from forming in small pores. If the inhibiting force were stronger than the functions currently implemented, for a given simulation we might expect to see less hydrate and ice formation. At lower hydrate and ice saturations, the simulator would likely have a much easier time converging since an appreciable amount of the liquid phase would still exist in the pore space.

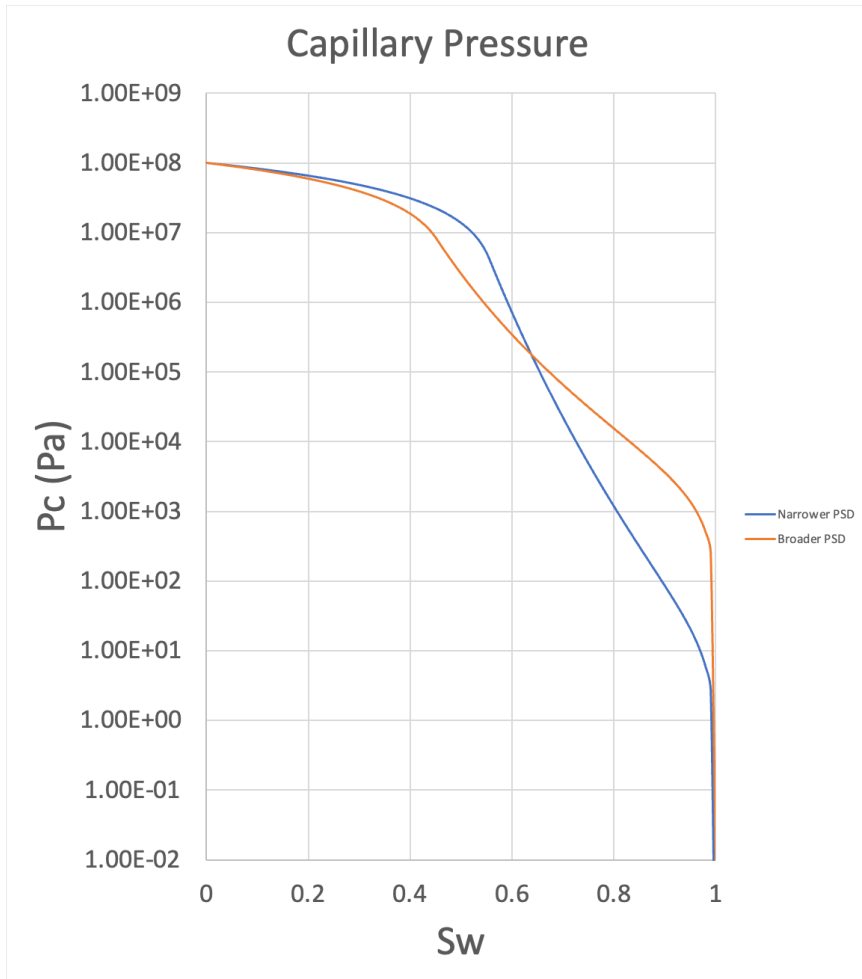


Figure 6. Capillary pressure function for the “narrower” and “broader” pore size distribution (PSD) used in the test of freezing with an internal methane source.

Table 4. Van Genuchten parameters for the “narrower” and “broader” pore size distribution cases of the test of freezing with an internal methane source.

	Narrower PSD	Broader PSD
Alpha (Pa^{-1})	0.1	5.8E-04
m	0.05	0.1
S_{ri}	0.1	0.1
Max P_c (Pa)	1E+08	1E+08

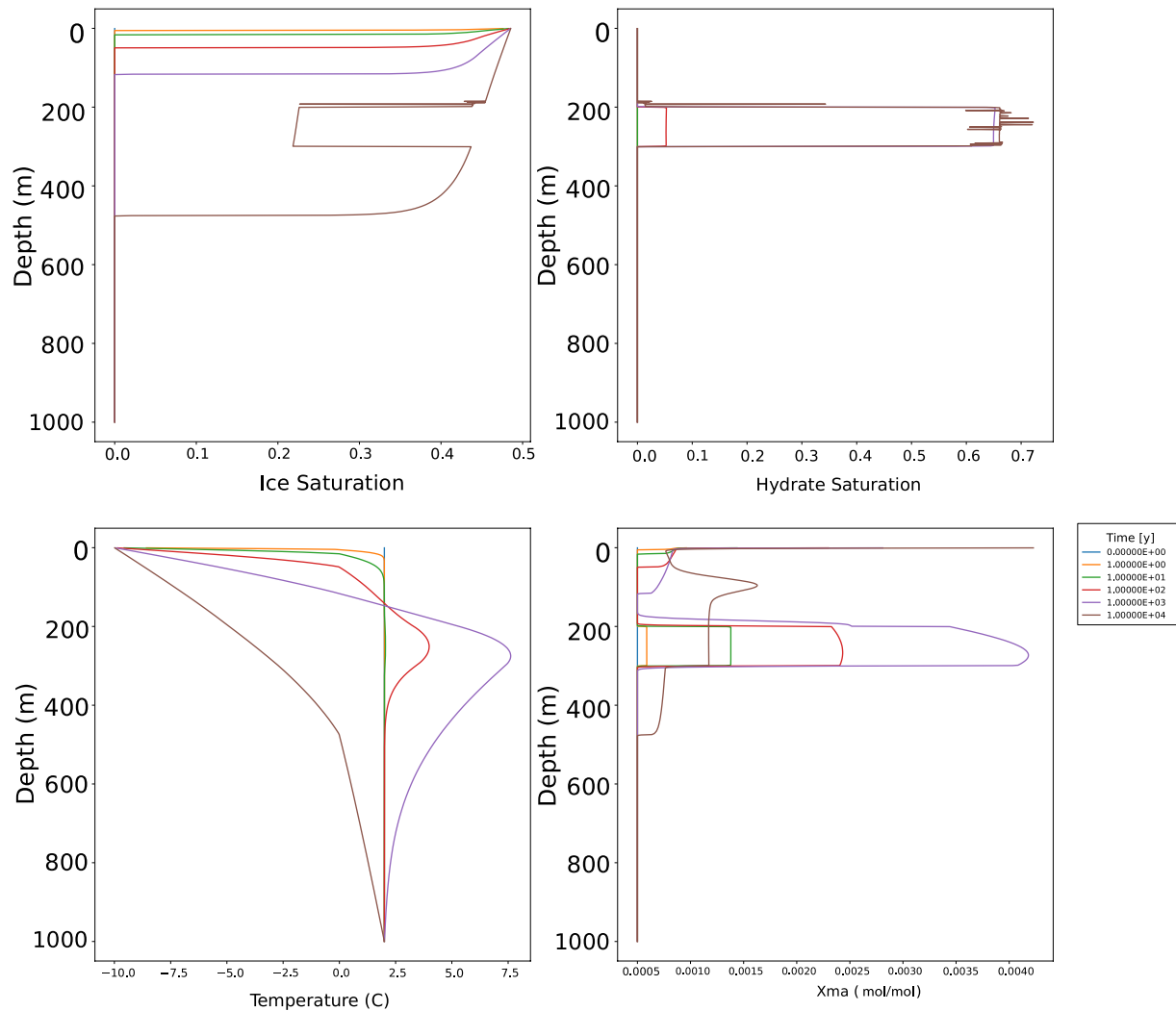


Figure 7. Ice saturation, hydrate saturation, temperature, and aqueous methane mole fraction evolution over time for a test involving freezing from the top and methane production within the model domain.

In addition to the saturation distribution, the interaction between hydrate and ice formation influences the thermal state of the system in this simulation. While the temperature distribution at 10 kyr is very

similar to the simulation example with only freezing, in earlier years while hydrate is forming there is a noticeable increase in temperature associated with exothermic hydrate formation over that period. Correspondingly, the methane mole fraction profile reflects concentration due to the downward propagation of the ice front as well as changes in methane solubility due to changes in temperature within the zone containing gas hydrate.

Like the simulations with only ice formation, these simulations were also run with two different pore size distributions expressed via the capillary pressure function (**Figure 6**). However, to achieve simulation convergence both pore systems had to be characterized by higher capillary pressures than in the previous example. This indicates that the conditions imposed on the model for this example, which result in significant formation of both ice and hydrate, require significant inhibition of ice and hydrate growth in order to keep enough of the pore space saturated with liquid to achieve convergence. A future area of investigation should be into whether this issue is resolved by running in 2D and 3D, if phase growth inhibition needs to be expressed in a more advanced way, or if more advanced numerical techniques need to be developed to overcome this hurdle.

Results of these simulations confirm the phenomenon that was observed in the previous test. In this test, the broader pore size distribution is associated with lower capillary pressures at ~50% phase saturation. It is in this range where we see separation in ice saturations in **Figure 8**, where the broader pore size distribution example exhibits higher ice saturations than the narrower pore size distribution. At total nonwetting phase saturations above 60%, however, the capillary pressures for both curves become similar, which is why the hydrate saturation profiles at the same time for both examples look similar.

The simulations presented in this section illustrate the need to constrain the growth of hydrate and ice in porous media as a function of pore size distribution. Without significant capillary inhibition to nonwetting phase growth, these solids can completely occlude the pore space and cause significant numerical challenges. These challenges can be addressed through manipulation of pore size distributions and through scaling the liquid and gas phase saturations as fractions of the mobile phases instead of as fractions of total pore space, but as the total hydrate + ice saturations approach 100%, there is clearly still work to be done to address numerical challenges. Future work will explore new numerical techniques to address this issue and also to potentially more realistically represent phase growth inhibition at high nonwetting phase saturations.

2.3.5.4. HYDRATE Mode Options

The following is a recap of some of the new options that are available in HYDRATE Mode that were discussed in previous sections and how to invoke them in PFLOTRAN. For examples of their usage, see the Appendix. The keywords available in HYDRATE Mode along with a brief description of their usage is included in **Table 5**.

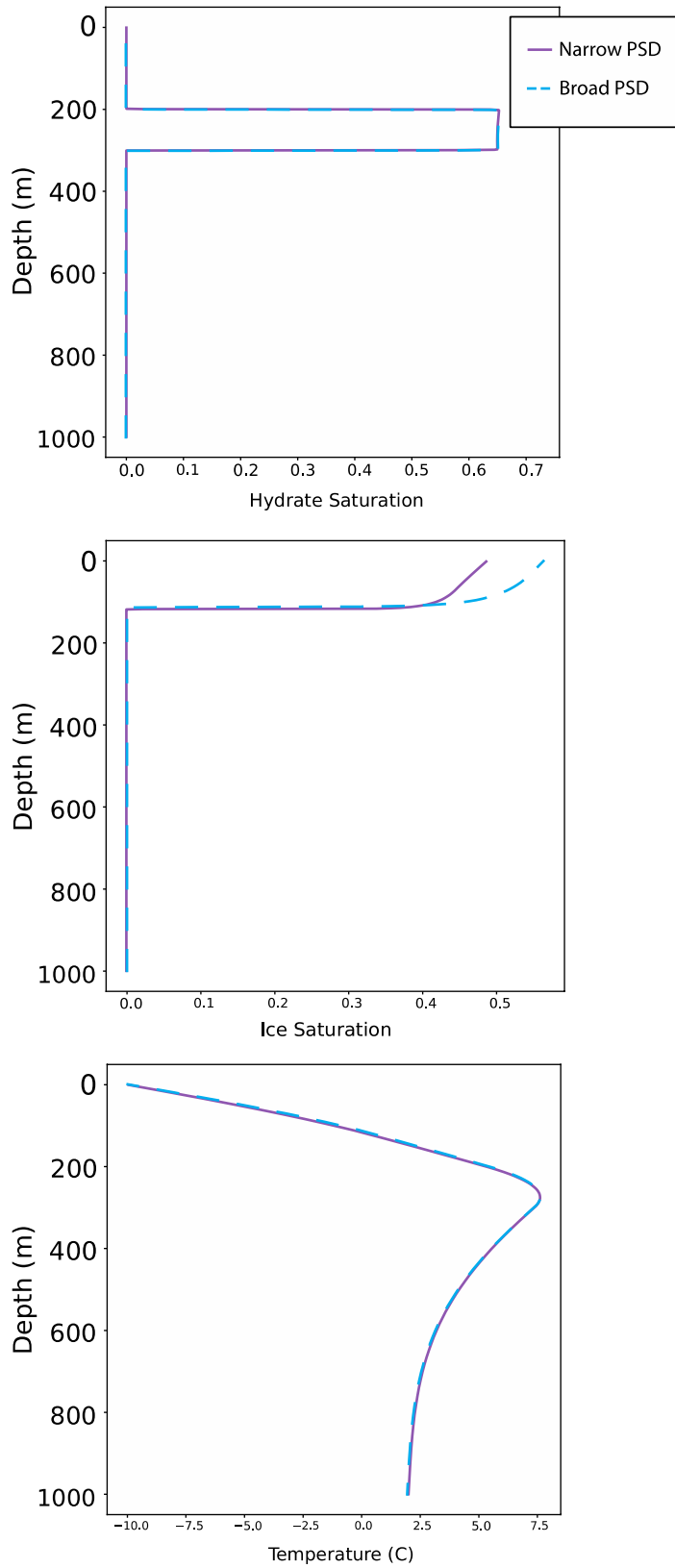


Figure 8. Hydrate saturation, ice saturation, and temperature at 1,000 years simulation time for the “narrower” and “broader” pore size distribution examples.

Table 5. Keywords available in HYDRATE mode and their usage.

Keyword	Usage
HYDRATE_PHASE_BOUNDARY	Set the equation for the phase boundary. Available options are MORIDIS, MORIDIS_SIMPLE, and KAMATH
WITH_GIBBS_THOMSON	Turns on the Gibbs-Thomson effect
GT_3PHASE	Turns on Gibbs-Thomson effect for 3 (or more) phases
SCALE_PERM_BY_HYD_SAT	Scales the permeability by hydrate saturation using the [75] formulation
EFFECTIVE_SAT_SCALING	Scales the liquid and gas phase fractions by the total saturation of mobile phases for input to relative permeability functions
ADJUST_SOLUBILITY_WITHIN_GHSZ	Adjusts the solubility of methane within the gas hydrate stability zone according to [77] formulation
BC_REFERENCE_PRESSURE	Sets the reference pressure for the DIRICHLET_SEEPAGE flow condition
SALINITY	Sets a constant salinity (mass fraction) to shift freezing point and hydrate phase boundary
WITH_SEDIMENTATION	Turns on sedimentation flux
METHANOGENESIS	Opens the METHANOGENESIS block, which currently sets a source term following [76]
ALPHA	Within the METHANOGENESIS block, sets the organic carbon mass fraction
K_ALPHA	Within the METHANOGENESIS block, sets the conversion factor from organic carbon to methane (should be set to 2241)
LAMBDA	Within the METHANOGENESIS block, sets the microbial reaction rate
V_SED	Within the METHANOGENESIS block, sets the sedimentation rate
SMT_DEPTH	Within the METHANOGENESIS block, sets the depth of the SMT (m), above which no methanogenesis occurs

3. PALEOCEANOGRAPHIC CONDITIONS OF THE ARCTIC OCEAN

Although the field of paleoceanography traces its origins to the 1950s [78], quantitative understanding of the late Pleistocene near-surface temperature history of the Arctic Ocean remains rudimentary [79]. Several factors contribute to this state of affairs, including: an overall paucity of sediment cores compared to other ocean basins, due in part to the difficulty of ship operations in sea ice; low sedimentation rates, commonly with hiatuses during glacial maxima (particularly in the western Arctic [79]); and small anticipated temperature signals relative to proxy noise, given that Holocene sea surface temperatures are already close to freezing. More paleotemperature work has been done in the high latitude North Atlantic south of Fram Strait (e.g. [80][81][82]), but here we treat that region as heavily influenced by the Atlantic Ocean and outside of the Arctic Ocean proper.

The oxygen isotopic composition ($\delta^{18}\text{O}$) of planktic foraminifera (cosmopolitan protozoans that secrete calcium carbonate shells) can serve as a paleotemperature proxy if the $\delta^{18}\text{O}$ of past seawater can be constrained. There are at least 20 existing Pleistocene-Holocene $\delta^{18}\text{O}$ records of varying quality and temporal scale in the Arctic [83], and about 40 Last Glacial Maximum (LGM, ~ 19 -23 kilo annum (ka)) observations [84][85]. Seawater $\delta^{18}\text{O}$ is affected by global ice volume, which is relatively easy to correct for, and the local hydrologic cycle. Unfortunately, the large dynamic range of salinity and hence seawater $\delta^{18}\text{O}$ in the Arctic makes foraminiferal $\delta^{18}\text{O}$ very difficult to interpret in terms of temperature. Instead, Arctic $\delta^{18}\text{O}$ records are best interpreted as qualitative salinity histories (e.g., [84][86]).

Aside from $\delta^{18}\text{O}$, the most commonly applied sea surface temperature proxies in paleoceanography are based on: biotic assemblages (foraminifera, radiolaria, diatoms, dinoflagellates); the magnesium content (Mg/Ca) of planktic foraminifera; and the saturation index of algal lipids (alkenone U^k_{37}). In the MARGO LGM compilation, the Arctic Ocean contains only foraminiferal assemblage data, which extend to 84°N in the North Atlantic sector and imply a warming of roughly 1°C [87]. We are not aware of any Late Pleistocene-Holocene Mg/Ca or alkenone sea surface temperature records from the Arctic Ocean proper.

Our best insights into Arctic Ocean near-surface temperatures come from Mg/Ca in benthic foraminifera and in ostracods, which are microscopic benthic crustaceans. While we would like to have benthic records from continental shelf depths, the published data are from deeper waters. In the Canadian Basin, [88] presented an 8 ka Holocene ostracod Mg/Ca record from 203 m water depth, and two 7 ka Holocene benthic foraminiferal Mg/Ca records from 369 m. Together these records vary between about -1 and 1.5°C , but with no clear trends over the Holocene. [89] compiled 50 ka of ostracod Mg/Ca data from 31 Arctic cores from a range of water depths and latitudes. During Marine Isotope Stage 3 (MIS 3, ~ 25 -50 ka), water depths of between 1 and 3 km warmed by $\sim 2^\circ\text{C}$, while the sole core from shallower depths (~ 250 m) was close to freezing. This pattern was interpreted as a warming and downward expansion of the Atlantic Layer, which is today found below the colder and fresher Polar Mixed Layer and Pacific Water and centered around 500 m water depth. Informed by numerical modeling, the cause of the intermediate depth warming was inferred to be a deepening of the glacial halocline in response to reduced runoff, preventing the Atlantic Layer (now dubbed Glacial Arctic Intermediate Water) from losing heat to the surface. Although [89] had fewer intermediate depth observations from the LGM, that interval appears to have warmed even more than MIS 3. These observations make some predictions for the glacial-interglacial history of the shallow subsurface, but data are sorely lacking.

3.1. Research Strategy

Conversations with Arctic paleoceanographers Tom Cronin (USGS, Reston) and Jesse Farmer (Princeton University) suggested that on the one-year timescale of this project, it would be very difficult to find shelf-depth cores with benthic foraminifera or ostracods. We therefore decided to target Mg/Ca in the ubiquitous planktic foraminifer *Neogloboquadrina pachyderma*, rendering the water depth of cores mostly irrelevant. In the modern Arctic, *N. pachyderma* calcifies mainly below the mixed layer, with larger (ontogenetically older) specimens calcifying deeper [90]. We selected the smallest size fraction that could be reliably handled, 150-212 μm , which is estimated to occupy water depths around 150 m on average. This habitat is in the upper half of the Arctic pycnocline, and it is possible that these organisms would migrate vertically if the pycnocline depth changed over time. However, the fact that *N. pachyderma* $\delta^{18}\text{O}$ and Mg/Ca do record past climate variations implies that these foraminifera don't simply track an isopycnal surface. Nevertheless, habitat uncertainty is an unavoidable complication when using planktic foraminifera.

N. pachyderma was one of the first planktic foraminifera to be calibrated for the Mg/Ca paleotemperature proxy [91]. Several subsequent *N. pachyderma* calibration studies consistently show exponential Mg/Ca sensitivities of 8 to 9% per $^{\circ}\text{C}$, with estimated pre-exponential constants ranging from ~ 0.4 to 0.6 [92]. In cold waters, the pre-exponential constant has a strong influence on reconstructed temperatures. Here we use the following approximation, which yields realistic Holocene temperatures for our measurements to date:

$$\frac{\text{Mg}}{\text{Ca}} = 0.7e^{0.09T} \quad \text{Equation 17}$$

Additional (underway) Mg/Ca measurements will help us determine whether this equation is broadly appropriate for Arctic Ocean *N. pachyderma*.

Mg/Ca-based paleotemperatures can theoretically be combined with *N. pachyderma* $\delta^{18}\text{O}$ to remove the temperature signal from the $\delta^{18}\text{O}$ of calcite, leaving an estimate of seawater $\delta^{18}\text{O}$. Although it is not straightforward to convert seawater $\delta^{18}\text{O}$ into a quantitative salinity history (due to uncertain $\delta^{18}\text{O}$:salinity relationships in the past, and the effect of brine rejection during sea ice formation), it should give us a better constraint on paleosalinity than foraminiferal $\delta^{18}\text{O}$ alone. We therefore leverage our Mg/Ca measurements by also analyzing $\delta^{18}\text{O}$ (and $\delta^{13}\text{C}$) on splits of the same samples.

3.2. Core Selection

Suitable sediment cores were identified with the help of Tom Cronin, who maintains a large collection of Arctic sediment samples. We surveyed 29 cores on the basis of their coordinates, temporal span, sedimentation rates, and evidence of hiatuses. Nearly all of the cores have radiocarbon-based age models. We selected four cores that provide a good balance between desirable temporal characteristics and broad spatial coverage. PS2163-1 is located in the Barents/Kara sector on the flank of Gakkel Ridge (86°N , 59°E , 3040 m water depth). It is 27 cm long and spans from ~ 3 to 43 kA with no apparent LGM hiatus, and sedimentation rates of ~ 0.6 - 0.8 cm/kA. It therefore provides a low-resolution continuous record from MIS 3 until the late Holocene. 96-12-P1 is close to the North Pole in the Laptev Sea sector on the Lomonosov Ridge (87°N , 145°E , 1003 m water depth). We sampled the core to 240 cm which corresponds to ~ 160 kA. 96-12-P1 has a mean sedimentation rate of ~ 1.5

cm/kA and no apparent LGM hiatus, so it provides a relatively high-resolution record of the entire past glacial-interglacial cycle (MIS 6 to Holocene). AOS94-B8 is close to the East Siberian and Chukchi Seas on the Mendeleev Ridge (78°N, 177°W, 1031 m water depth). It spans from ~4 to 52 kA but with a 14 kA hiatus across the LGM. Sedimentation rates are relatively high at ~1-2 cm/ky, so this core provides a relatively high-resolution look at MIS 3 and the Holocene. Finally, P1-92-AR25 sits on Northwind Ridge in Canada Basin (75°N, 157°W, 1625 m water depth). Although this core is missing the Holocene, it has an average sedimentation rate close to 5 cm/kA and no apparent LGM hiatus. It therefore provides a high-resolution record of MIS 3 (from ~42 kA) through the deglaciation.

3.3. Analytical Methods

Previously washed and sieved (sand fraction) sediment samples were dry-sieved to isolate the 150-212 μm particle sizes. *N. pachyderma* specimens were then individually picked using a wetted paintbrush under a binocular microscope. There are at least five recognized morphotypes of *N. pachyderma* [93],[94], and we targeted the quadrate Nps-2. From each sample we picked 50 specimens for trace metal analysis (~300 μg CaCO_3) and 25 specimens for stable isotope analysis (~150 μg). Given a total of 200 sediment samples across the four cores, this amounted to 15,000 foraminifera picked.

Prior to trace metal analysis, the foraminifera must be cleaned of contaminants. The first step in this process is to crush the shells between two glass slides under a binocular microscope, to expose interior chambers to the chemical cleaning process. Crushed samples are then transferred to polypropylene microcentrifuge tubes to undergo a three-stage cleaning process [95][96]. The first stage removes fine-grained extraneous material (clays) through repeated sonication and siphoning of supernatant. This is followed by reductive cleaning in a buffered anhydrous hydrazine solution, to remove iron/manganese oxide coatings. Finally, organic matter is oxidized using a buffered hydrogen peroxide solution. Once samples are crushed, 30 samples can be cleaned in about six hours. Cleaned samples are then dissolved in nitric acid and analyzed on a ThermoFinnigan Element2 sector field inductively-coupled plasma mass spectrometer (ICP-MS) at the University of Colorado Boulder Trace Metal Lab [97]. A total of 11 elemental ratios are measured against liquid standards, with our focus here on Mg/Ca. Analytical precision (1σ RSD) is $\pm 0.5\%$ for Mg/Ca.

Foraminifera do not need to be chemically cleaned for $\delta^{18}\text{O}$ analysis, only weighed on a microbalance. $\delta^{18}\text{O}$ and $\delta^{13}\text{C}$ are measured on a Thermo Delta V continuous-flow stable isotope ratio mass spectrometer attached to a GasBench II gas preparation device, at the University of Colorado Boulder Earth Systems Stable Isotope Laboratory. Repeated measurements of in-house standards yield precisions of $\pm 0.1\text{‰}$ or better for both $\delta^{18}\text{O}$ and $\delta^{13}\text{C}$.

3.4. Results

As of this writing, only core PS2163-1 has a complete Mg/Ca record. For the other three cores, sample preparation for trace metals (picking, crushing, and cleaning) is about 85% complete, but analysis by ICP-MS awaits. Sample weighing for stable isotopes is also substantially complete, but analysis cannot occur until October because the global helium shortage prevents operation of the mass spectrometer until then.

At the core site of PS2163-1 (86°N, 59°E), *N. pachyderma* Mg/Ca documents subsurface (150 m) temperatures of $\sim 1^\circ\text{C}$ during the mid-Holocene, warming to $\sim 2^\circ\text{C}$ during the early Holocene and deglaciation, and peaking at $\sim 3^\circ\text{C}$ during the LGM (**Figure 9**). Coldest temperatures, close to the

freezing point of seawater, occurred during late MIS 3. Earlier in MIS 3, temperatures were around 0°C.

The MIS 3 cooling is consistent with a thickening of the polar surface/mixed layer as hypothesized by [89]. The LGM temperatures suggest that the maximum intermediate depth warming documented by [89] at that time extended close to the sea surface, implying a shallow halocline, perhaps due to extensive ice cover. During deglaciation, excessive runoff from ice sheet melting likely maintained a shallow mixed layer and halocline [86], but by then the intermediate water was cooler [89].

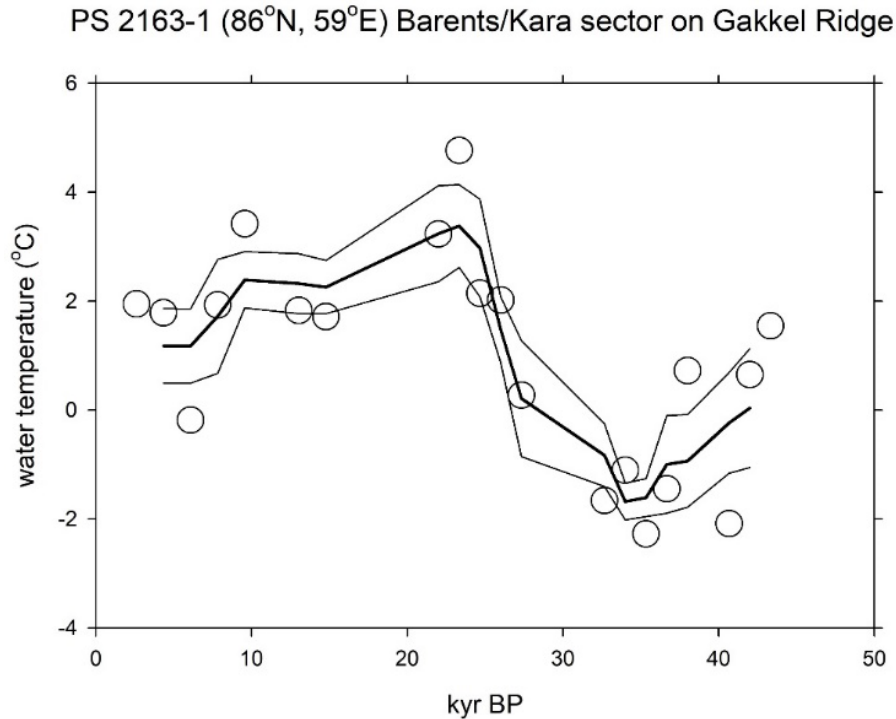


Figure 9. Subsurface (~150 m) paleo-temperatures recorded by *N. pachyderma* Mg/Ca in core PS2163-1. Heavy line is a three-point running mean and light lines are ± 1 standard error uncertainty.

We use the Mg/Ca record from core PS2163-1 to generate a preliminary scenario for the temperature history of the Arctic near surface (**Figure 10**). This scenario is informed by prior work and the paleoceanographic interpretation described above, but is currently quite speculative until the other three cores are analyzed. Reconstructed Late Holocene temperatures of 1°C are assigned to the last interglacial (MIS 5), except for the early portion of the peak interglacial MIS 5e. The early Holocene and deglacial temperature of 2°C is assigned to early MIS 5e and the penultimate deglaciation, as well as (more speculatively) the mini-deglaciation following MIS 4 at ~53-62 ka [86]. The glacial peaks of the LGM, MIS 4, and MIS 6 are assigned the LGM temperature of 3°C. Finally, MIS 3 currently stands alone as the coldest interval. To apply this history in the numerical model, we propose adding the temperature anomaly (red line in **Figure 10**) to the 150 m modern temperature in each grid cell. For simplicity, sea surface temperatures can be held constant at the modern values, since they are typically close to freezing already. Then anomalies can be interpolated with water depth, from the assigned value at 150 m to zero at the sea surface.

PS 2163-1 (86°N, 59°E) Barents/Kara sector on Gakkel Ridge

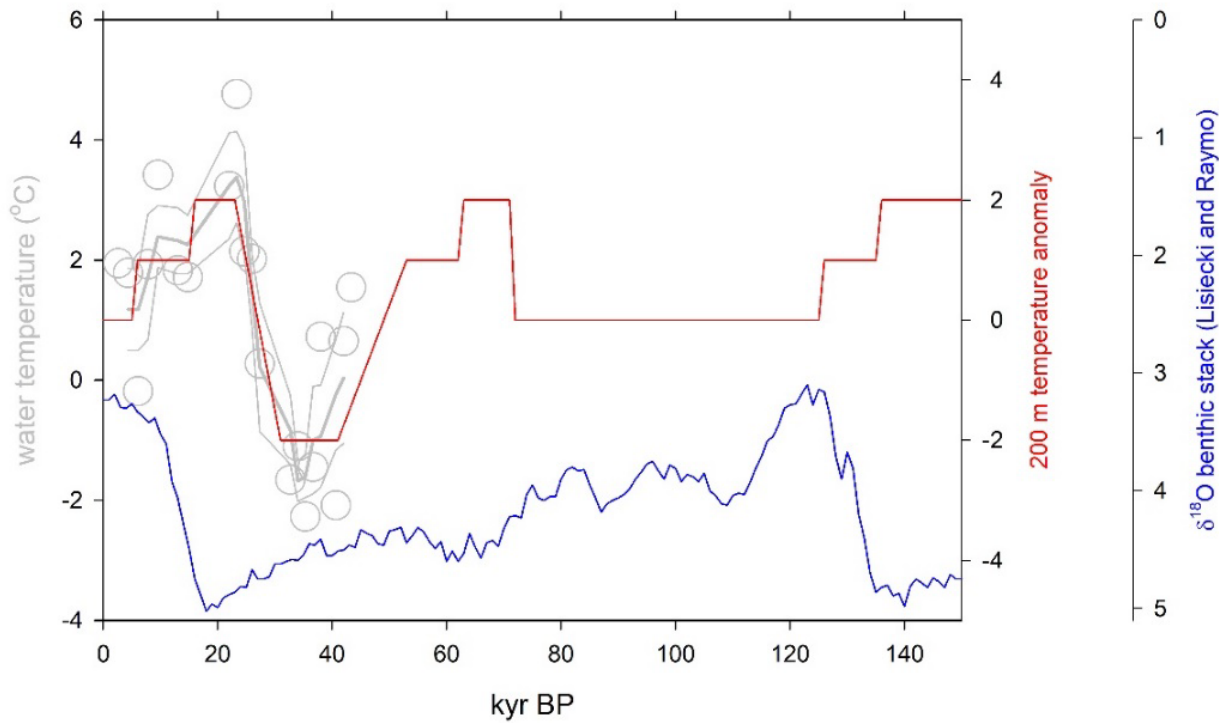


Figure 10. Mg/Ca paleotemperatures as shown in Figure 9 (gray), with the addition of our preliminary scenario for temperature anomalies through time (red). The global benthic $\delta^{18}\text{O}$ stack of Lisiecki and Raymo (2005) is shown for reference (blue).

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4. DEMONSTRATION OF THE HYBRID MODELING FRAMEWORK ON THE ALASKA DOMAIN

In this section, we present a demonstration of the developed framework that applies a hybrid modeling technique using geospatial machine learning and geophysical numerical modeling to produce probabilistic predictions of methane emissions from the Arctic seabed. Because the majority of *dissolved* methane flux is thought to be taken up by the anaerobic oxidation of methane (AOM) [23], mediated by a consortium of methane-oxidizing Archaea and sulfate-reducing bacteria (as described in **Section 1.2**), we focus instead on *ebullition*, or gas phase methane flux from the seabed to the water column. Because methane in gaseous form is not directly accessible to the AOM community [27], it is assumed that the entire amount of methane transported within the gas phase can reach the water column. Furthermore, to estimate how much of the modeled ebullitive flux may reach the atmosphere, simple yet highly predictive 1D models of free and dissolved gas transport in the shallow water column toward the atmosphere [31],[32] can be adopted. In general, these models show that bubbles released only from shallow water depths (e.g., ~ 100 m or less) may have appreciable amounts of methane left inside of them by the time they reach the ocean surface. Therefore, we have reduced the simulation domain to treat seabed areas with a present-day water depth less than 175 m. This limits the simulation domain to the Arctic continental shelf. Finally, for the demonstration presented here, we further limit the computational domain to the region surrounding Alaska (with further details given in the following subsections).

We note that the modeling workflow demonstration presented in this section is still a work in progress, but at the time of this writing, has been developed to an advanced stage capable of being exercised through its entire workflow – from model set-up to model prediction. Here we present the details of this workflow in their current advanced stage and point out sections of the framework’s needed improvements where applicable. As such, the results presented here are not meant to be interpreted as new scientific evidence, but rather, a demonstration of how new scientific modeling predictions using this framework can be made and used to better understand marine sources of methane emissions under climate warming and how this framework can be used to provide a source term to global climate models over space and time.

4.1. Geographic Location

The modeling demonstration site was chosen surrounding Alaska’s modern-day shoreline. This region experienced large variations in sea level across past glacial cycles as did the entire Arctic coastline (as described in **Section 1.1**). As shown in **Figure 11** (in red), the relative sea level (RSL) for the Beaufort Sea varied from ~ 30 m above present day sea level 120 kA BP to ~ 100 m below present day sea level at the LGM [98],[99]. In blue, **Figure 11** shows the mean annual air temperature (MAAT) over the last glacial cycle, which was at times $\sim 10^\circ\text{C}$ colder than modern temperatures [100],[101]. As a consequence, vast regions of the presently submerged Alaskan continental shelf would have been exposed during portions of the last glacial cycle. In North America, Alaska is unique because the state encompasses the largest contiguous land area at high latitude that was never glaciated over its lower elevations [102], therefore, it is understood that the entire exposed seabed would have formed permafrost (where MAAT was cold enough, such as in the northern portions of the state). After the LGM, as the environment warmed, sea level began to rise back up again. This marine transgression changed the thermal conditions of inundated permafrost, initiating warming and thawing of submarine permafrost that continues today.

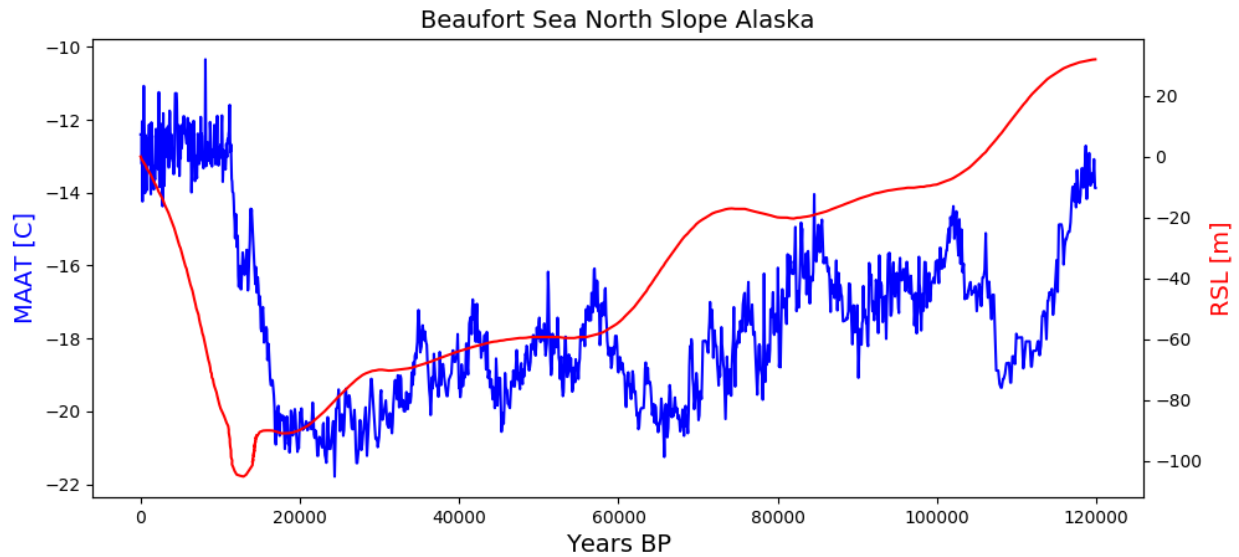


Figure 11. Evolution of the mean annual air temperature (MAAT – blue, left axis) and relative sea level (RSL – red, right axis) over the last glacial cycle. RSL records were made based on [98][99] and personal communication with J. X. Mitrovica. MAAT records were made based on the temperature anomaly from [100] applied to modern-day temperatures [101].

To demonstrate the current capabilities of the hybrid modeling framework developed in this study, we modeled probabilistic submarine permafrost and methane gas distribution and gas flux in the seabed off the coast of Alaska in response to environmental forcing. The simulations were set up using a flexible version of the Dakota-PFLOTRAN workflow that runs on Sandia’s Common Engineering Environment’s (CEE) High Performance Compute Servers. The computational domain is situated between 170°W & 140°W and 60°N & 74°N and includes only the locations where present-day water depths are less than 175 m within these coordinates. Furthermore, we down-selected 2,009 locations within this area to run 1D simulations at an approximate spatial resolution of 5 arc minutes in longitude and 10 arc minutes in latitude, resulting in a computational domain that covers the majority of the shallow continental shelf of Alaska (Error! Reference source not found.). 1D simulation locations are shown as black dots.

4.2. Model Setup and Methods

GPSM parameters within the geographic boundaries (170°W & 140°W and 60°N & 74°N) were obtained at a resolution of 5 arc minutes. The available parameters included seafloor depth (**Figure 13A**), seafloor temperature (**Figure 13B**), heat flux (**Figure 13C**), sedimentation rate (**Figure 13D**), and total organic carbon (**Figure 13E**). Of these parameters, the latter four were applied as conditions/constraints at the top of the down-selected 1D simulated sediment columns (also referred to as “profiles”). Seafloor depth was used to calculate evolving top boundary conditions for pressure and temperature.

The RSL history accounts for changes in sea level due to glaciation (**Figure 11**). RSL history is used to calculate the overburden pressure at each simulation location over time associated with changing water depth relative to the present-day depth. In addition, it is also used to calculate the temperature at the top of the 1D sediment column according to when the profile would be exposed to atmospheric conditions as opposed to submerged oceanic conditions. When profiles are exposed to the

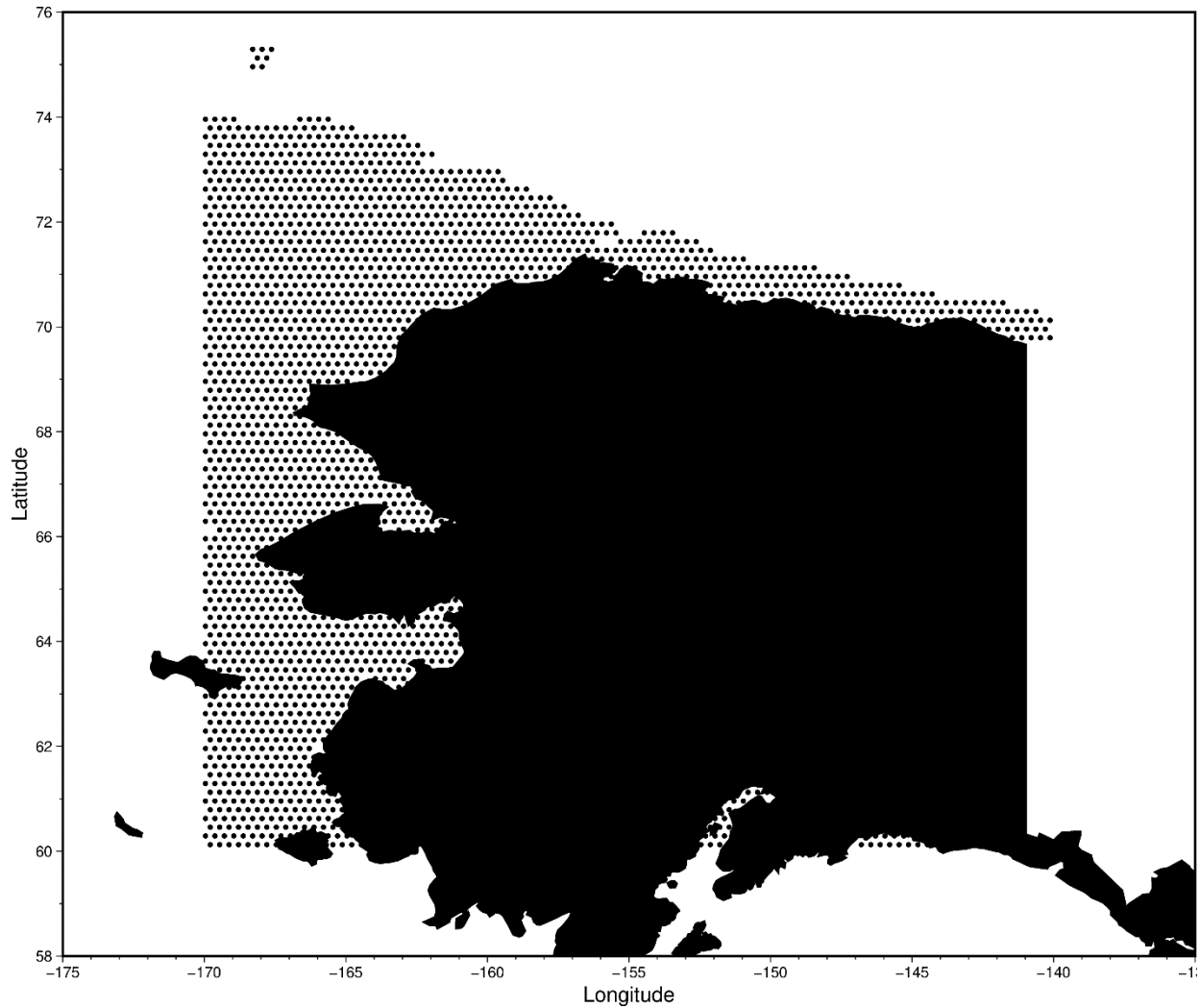


Figure 12. Geographic location of study site and 1D simulation profile locations.

atmosphere, the MAAT value from **Figure 11** is applied. When profiles are submerged by the ocean, the modern-day seabed temperature is applied.

The method to determine boundary conditions used here employs two simplifications that should be improved with future use. First, the RSL curve used is site-specific to the nearshore Beaufort Sea along the North Slope, yet we applied it to the entire computational domain. This is a simplification we made because this was the only RSL curve available to us at the time of this writing. However, results can be improved by using many site-specific RSL curves and applying them in their respective regions within the larger computational domain. Similarly, this should be done for MAAT also. Second, when the profiles are submerged by the ocean, the modern-day seabed temperature is applied. This is another simplification we made because we did not have a historic record of the bottom water temperature, which should be applied instead. However, as informed by prior work and the paleoceanographic interpretation described in **Section 3** of this study, we used the Mg/Ca record from core PS2163-1 to generate a preliminary scenario for the temperature history of the Arctic near surface (**Figure 10**). To apply this history in the numerical model, we propose adding the temperature anomaly (red line in **Figure 10**) to the 150 m modern temperature in each grid cell. For simplicity, sea

surface temperatures can be held constant at the modern values, since they are typically close to freezing already. Then anomalies can be interpolated with water depth, from the assigned value at 150 m to zero at the sea surface. This work is planned for FY23 by the authors.

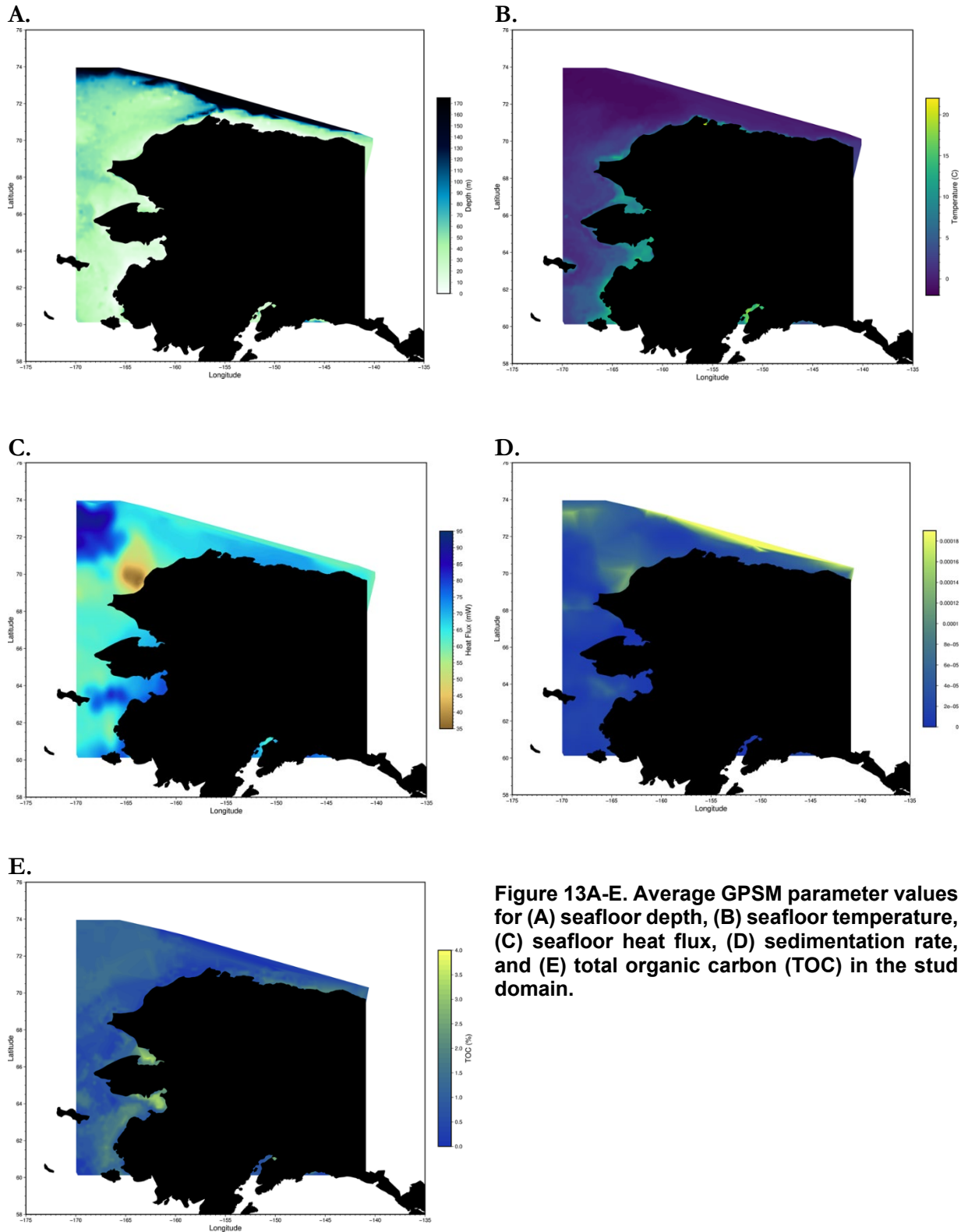


Figure 13A-E. Average GPSM parameter values for (A) seafloor depth, (B) seafloor temperature, (C) seafloor heat flux, (D) sedimentation rate, and (E) total organic carbon (TOC) in the stud domain.

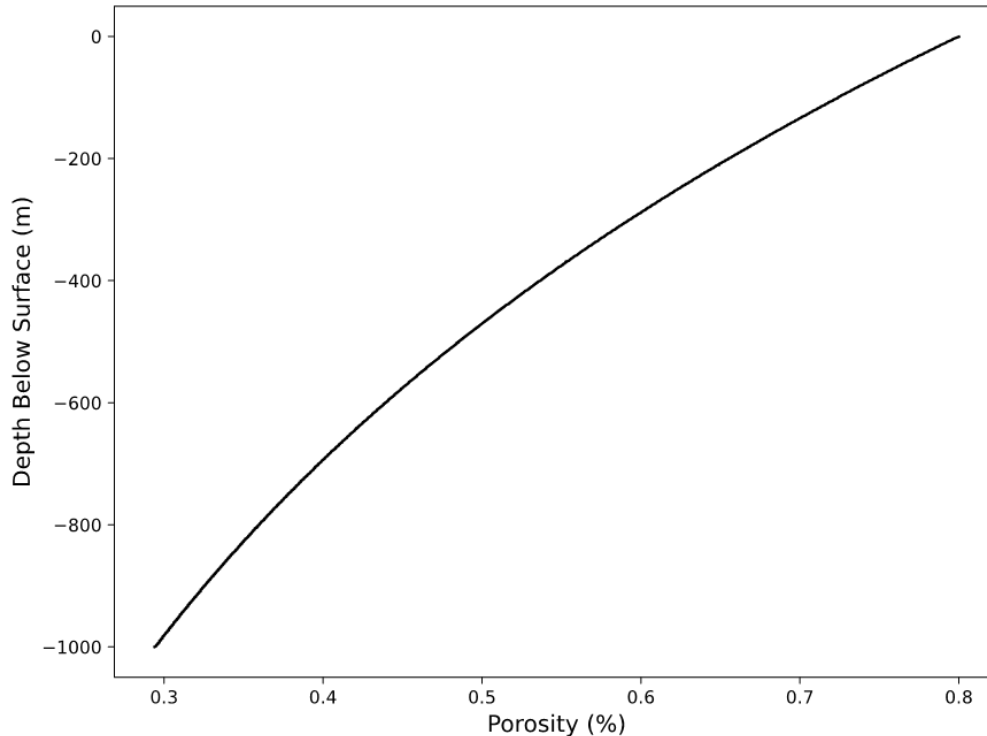


Figure 14. The porosity (%) profile as a function of depth (m). The same profile is used for both ice and gas simulations.

Porosity at each 1D simulation profile location decreases with depth accounting for compaction and is shown in **Figure 14**.

Simulations for each 1D profile location occur in two phases – the first phase simulates submarine formation (dubbed “ice simulations”) while the second phase simulates both permafrost and methanogenesis (dubbed “gas simulations”). All input decks and associated external text files are included in the Appendix, **Sections A.4, A.5, and A.6**.

4.2.1. Ice Simulations

Each 1D profile location begins with an ice simulation to estimate submarine permafrost distribution. For each ice simulation, Dakota takes the average value of the seafloor depth, seafloor temperature, and heat flux from the provided GPSM dataset. The sedimentation rate and total organic carbon from GPSM are not used for the ice simulations. In addition, the methanogenesis model within PFLOTTRAN Hydrate Mode is turned off, and as a result, no gas is simulated during this phase.

All ice runs simulate 405 kA of time, representing three full glacial cycles which begin 525 kA BP. Each glacial cycle is a repeat of the most recent glacial cycle (which occurred 0 – 120 kA BP in reality) and is shown in **Figure 11**. Three glacial cycles were simulated so as to give the model “spin-up time,” allowing the evolving permafrost depth to come to a quasi-steady state before gas simulations were started. Due to numerical issues associated with the sharp discontinuity in RSL between cycles, the discontinuity is smoothed by extending the 120 kA cycle duration to a duration of 135 kA, and a smooth connection is made between the end of one cycle and the beginning of the next. This was done purely for numerical reasons. A better approach would be to use historical records for each past glacial cycle so that the data is naturally continuous. Because we only had information on RSL for the

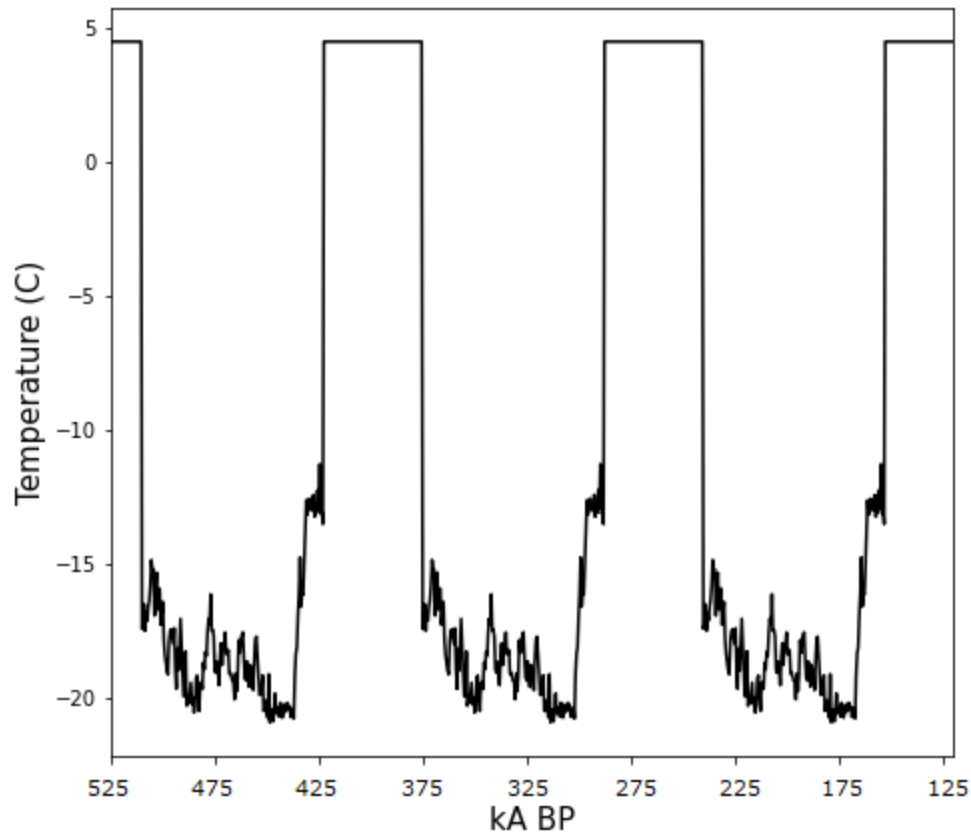


Figure 15. Top temperature boundary condition for an example ice simulation, Profile 332811, located at 165.792°W and 66.4583°N, where present day water depth is 14 m.

most recent glacial cycle (0 – 120 kA BP), we used the information we had and repeated it three times in lieu of actual records.

The simulation is driven by the evolving boundary conditions for temperature and pressure over the span of the 3 back-to-back glacial cycles. **Figure 15** and Figure 16 show the top temperature and pressure boundary condition for one of the 1D simulations, Profile 332811, located at 165.792°W and 66.4583°N, where present day water depth is 14 m. Temperature alternates between the seabed temperature ($\sim 3^{\circ}\text{C}$) and the MAAT as applied from **Figure 11**. Similarly, the pressure alternates between the hydrostatic pressure and atmospheric pressure. This evolving boundary condition is specified in PFLOTRAN HYDRATE Mode via the LIST feature in an external file called from the main input deck (see Appendix **Section A.6**). Initial conditions are made physically consistent with the boundary conditions applied at 525 kA BP. An example PFLOTRAN HYDRATE Mode input deck for an ice simulation is provided in the Appendix **Section A.4**.

Model output is given every 5 kA, and includes variables of interest such as ice saturation, liquid pressure, temperature, etc. The particular output frequency and variable list is adjustable by the user. Additionally, a checkpoint/restart file is requested at the end of the ice simulation run, at a simulation time of 405 kA, equivalent to 120 kA BP. In this checkpoint/restart file, all of the variables PFLOTRAN needs to restart the simulation again are saved. The restart file is used to begin the last glacial cycle of the simulation, which includes methanogenesis (described in next section).

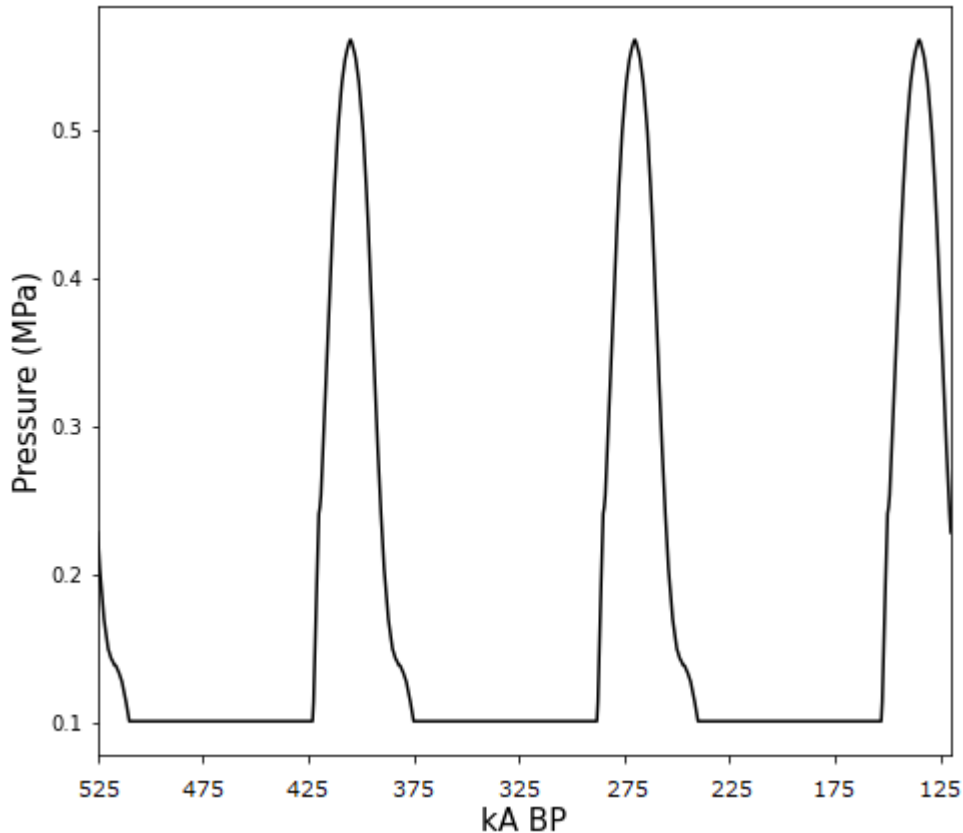


Figure 16. Top pressure boundary condition for an example ice simulation, Profile 332811, located at 165.792°W and 66.4583°N, where present day water depth is 14 m.

4.2.2. Gas Simulations

The ice simulations are followed by gas simulations designed to model methanogenesis and methane flux from the seabed to the water column at each location. Each gas simulation begins by utilizing the restart file generated by the ice simulation at 405 kA of simulated time (equivalent to 120 kA BP). This allows the overall workflow to pick up exactly where it left off. Each gas simulation runs for 150 kA, which includes an additional (4th) glacial cycle plus 30 kA into the future, putting the end point of each simulation at 555 kA of simulation time. Present day is represented at 525 kA of simulation time, or 0 kA BP, while the LGM is represented at 507 kA of simulation time, or 18 kA BP. Example pressure and temperature data for the gas simulation at the same location represented in **Figure 15** and **Figure 16** are shown in red in **Figure 17**. In addition, the gas simulations also utilize the same external flow conditions file (**Section A.6**) as the ice simulations.

We note that the 30 kA future years from the present day have not been designed to represent projected warming scenarios, as ideally should be the case. Instead, we simply used the extended portion of the glacial cycle that was designed to smooth the RSL curve. In effect, this extends the present-day conditions for temperature while simulating a small sea level rise. Future work will create more physically meaningful future scenarios for the modeling workflow. We include modeling results for 10 kA in the future as an example of the workflow capability, but these results should not be considered or interpreted as new scientific evidence.

Previously, the ice simulations used GPSM average parameter values for seafloor depth, seafloor temperature, and heat flux. The methane simulations also rely on these data, but in addition, sample on sedimentation rate and total organic carbon (TOC). For each 1D simulation profile, Dakota selects 25 samples among combinations of sedimentation rate and TOC Latin applying Hypercube methodology. These samples make up 25 ensemble members (or alternatively, “realizations”) at each profile location

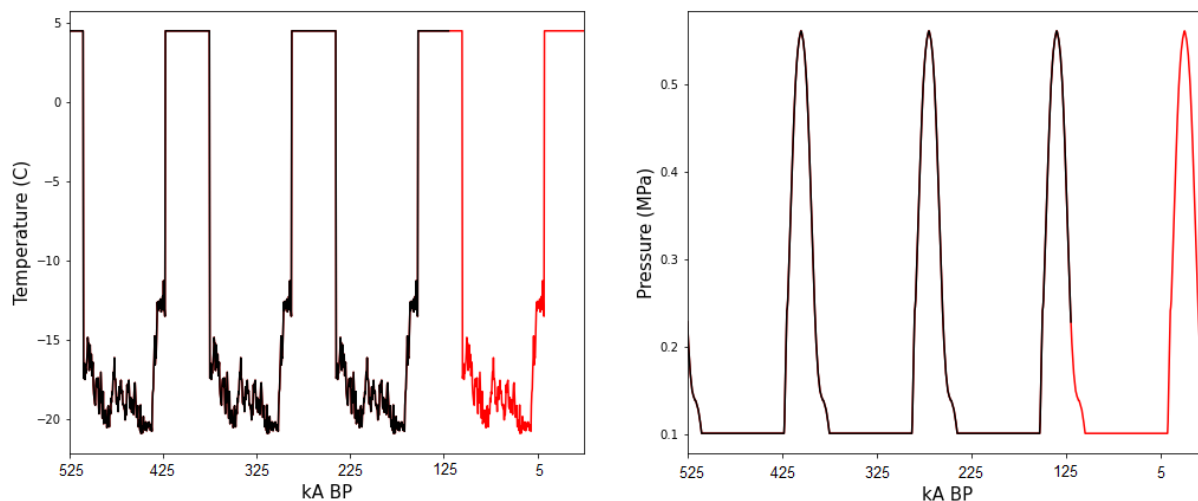


Figure 17. Top pressure (left) and pressure (right) boundary conditions for an example simulation, Profile 332811, located at 165.792°W and 66.4583°N, where present day water depth is 14 m. In black, the ice simulation portion of the simulation is shown, while in red the gas simulation portion is shown.

and enable probabilistic modeling results for methane gas flux. Additionally, both sedimentation and methanogenesis are turned on within the PFLOTTRAN simulation. This allows methane gas to form and migrate throughout the shallow offshore sediments. We note that the TOC has been multiplied by a factor of 50 in these demonstration simulations in order to increase the methanogenesis over the short time frame of simulation (150 kA). The reasoning behind this was purely for visualization purposes in the demonstration. In reality, such geologic systems have evolved over several millions of years, allowing gas and gas hydrate to form and deposit, yet we only simulate a fraction of this time. When we set the TOC to realistic values, very little gas accumulated within the short time span of the simulation. Rather than initialize the simulations with gas and methane hydrate, we simply increased the rate of methanogenesis so that we could demonstrate the capabilities of this workflow.

We note that the number of samples (25) should ideally be much larger to allow for proper statistical analysis of the probabilistic results. We limited our sampling to 25 due to computational resource limitations and the computation time available to us. With 25 samples, one should capture the average of a particular quantity in the ensemble results fairly well; however, in order to get meaningful statistics for the true possible range in model predictions (e.g., the tails of a distribution), at least 100 samples should be run.

An important value to note within the PFLOTTRAN input deck is GAS_RESIDUAL_SATURATION. This value indicates the minimum gas saturation required for free gas mobility throughout the sediment column. When gas saturation is lower than the residual gas saturation (set to 0.10 in these demonstration simulations), only dissolved gas is mobile and ebullition will not be possible. The residual gas saturation is a poorly constrained quantity in geophysical

modeling. Therefore, it is recommended that future modeling set-ups sample on this quantity, in addition to sedimentation rate and TOC. In general, we expect that smaller values of residual gas saturation will increase methane fluxes and may shift the onset of emissions to earlier times after permafrost thaw.

4.3. Modeling Results

We note, again, that the modeling workflow demonstration presented in this section is still a work in progress, but at the time of this writing, has been developed to an advanced stage capable of being exercised through its entire workflow – from model set-up to model prediction. As the capabilities of the PFLOTRAN HYDRATE Mode model are still being debugged and fine-tuned, the following results should be observed more as a proof of capabilities rather than as reliable, predictive science results. Given this expectation, we present the results of the ensemble simulations next.

4.3.1. Simulation Performance Statistics

A single ice simulation utilizing GPSM average parameter values was initialized for each of the 2,009 Alaskan domain locations (previously shown in **Figure 12**) and run for a total of 405 kA of simulation time. Of the 2,009 simulations, 1,973 were successful, while the remaining 36 PFLOTRAN HYDRATE Mode simulations failed. This equates to an ice simulation success rate of approximately 98%. Locations where the ice simulations failed to produce results are left blank on the figures.

Upon completion of successful ice simulations, gas simulations applying 25 Dakota-sampled GPSM parameters for combinations of sedimentation rate and TOC were initialized. They ran for 150 kA of simulation time. Of the 1,973 locations modeled to include methanogenesis, 1,493 were successful, while the remaining 480 PFLOTRAN HYDRATE Mode simulations failed to complete. This equates to a gas simulation success rate of approximately 75% - much lower than that of the ice simulations. In **Figure 18**, we present a map of all failed simulations for visual reference.

The main reason for failed simulations was numerical difficulty, especially when methanogenesis was included (e.g., the gas simulations). As described in **Section 2.3**, in the formulation developed here, liquid water is considered the only wetting phase. All other phases (free gas, gas hydrate, ice) are considered nonwetting phases. In order to apply a freezing point depression for ice phase or hydrate phase growth, the phase saturations are plugged into a capillary pressure function to obtain a phase capillary pressure and thus a freezing temperature offset. This is straightforward when transitioning from single phase liquid to two-phase (liquid-ice or liquid-hydrate). It is more complicated, however, when dealing with 3 or 4 phases coexisting in the pore space. In addition, PFLOTRAN's HYDRATE Mode uses primary variable switching to evaluate state properties of the system in any given grid cell at a given simulation time. There are 15 possible combinations for phase state, each corresponding to a set of primary variables. We have found that the numerical stability is most sensitive to the phase states and capillary pressure curve formulation. The majority of the failing simulations experienced trouble during phase transitions, such as when a freezing front passes through a region with mobile gas, for example. As seen in **Figure 18**, many of the failing gas simulations coincide with an area of anomalously low heat flow (**Figure 13C**), which would promote freezing. We note that PFLOTRAN HYDRATE Mode still requires development work to improve its robustness in this respect.

4.3.2. Submarine Permafrost Distribution

The average permafrost distribution for the LGM (18 ka BP), present day, and 10 kA in the future are presented in **Figure 19**, **Figure 20**, and **Figure 21**. These figures show the depth to the bottom of the submarine permafrost layer and were calculated as the deepest depth where ice saturation was

non-zero. How deep the frozen layer reaches at the LGM depends on several factors, namely the timing of exposure (which is primarily determined by the bathymetry), bottom water temperature prior to atmospheric exposure, and the heat flux (which determine the temperature profile in the sediment column at the time of atmospheric exposure). The depth of the submarine permafrost at the present day depends most heavily on the rate of thaw after submergence. The most important players are bottom water temperature, heat flux, and timing of submergence. Permafrost along the shelf edges has been submerged the longest and thus is expected to be the most degraded. However, regions where heat flux is high or bottom water temperatures are warm due to river discharge may create regions which support accelerated thaw unrelated to the timing of submergence. Regions where permafrost thaw is substantial is hypothesized to promote methanogenesis, due to an increased availability of organic matter, and increased sediment permeability for fluid flow.

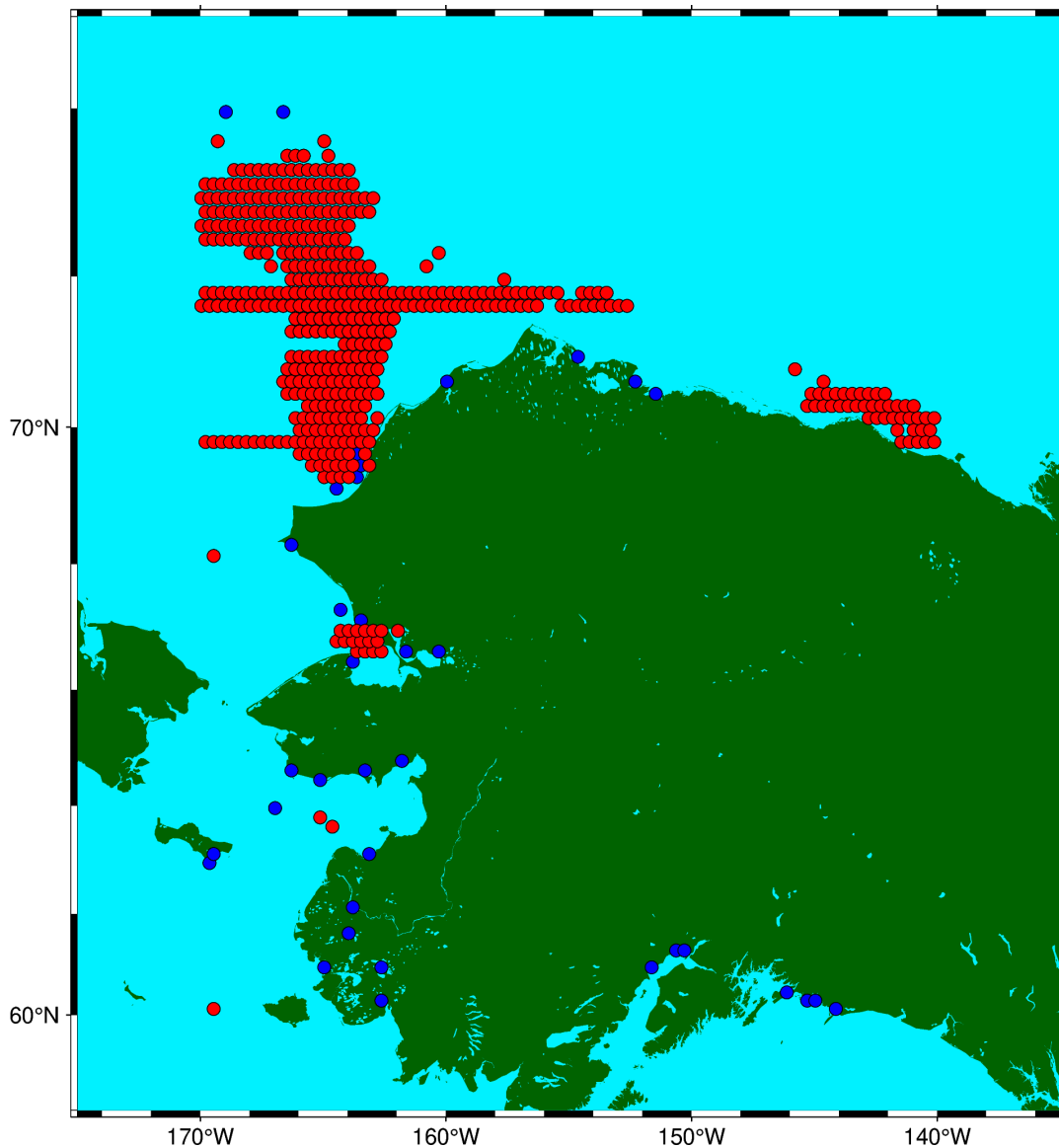


Figure 18. Map of failed PFLOTRAN HYDRATE Mode simulations where no data was available as a result. Blue show failed ice simulations (~ 2% of all simulations) and red show failed gas simulations (~ 25% of all simulations).

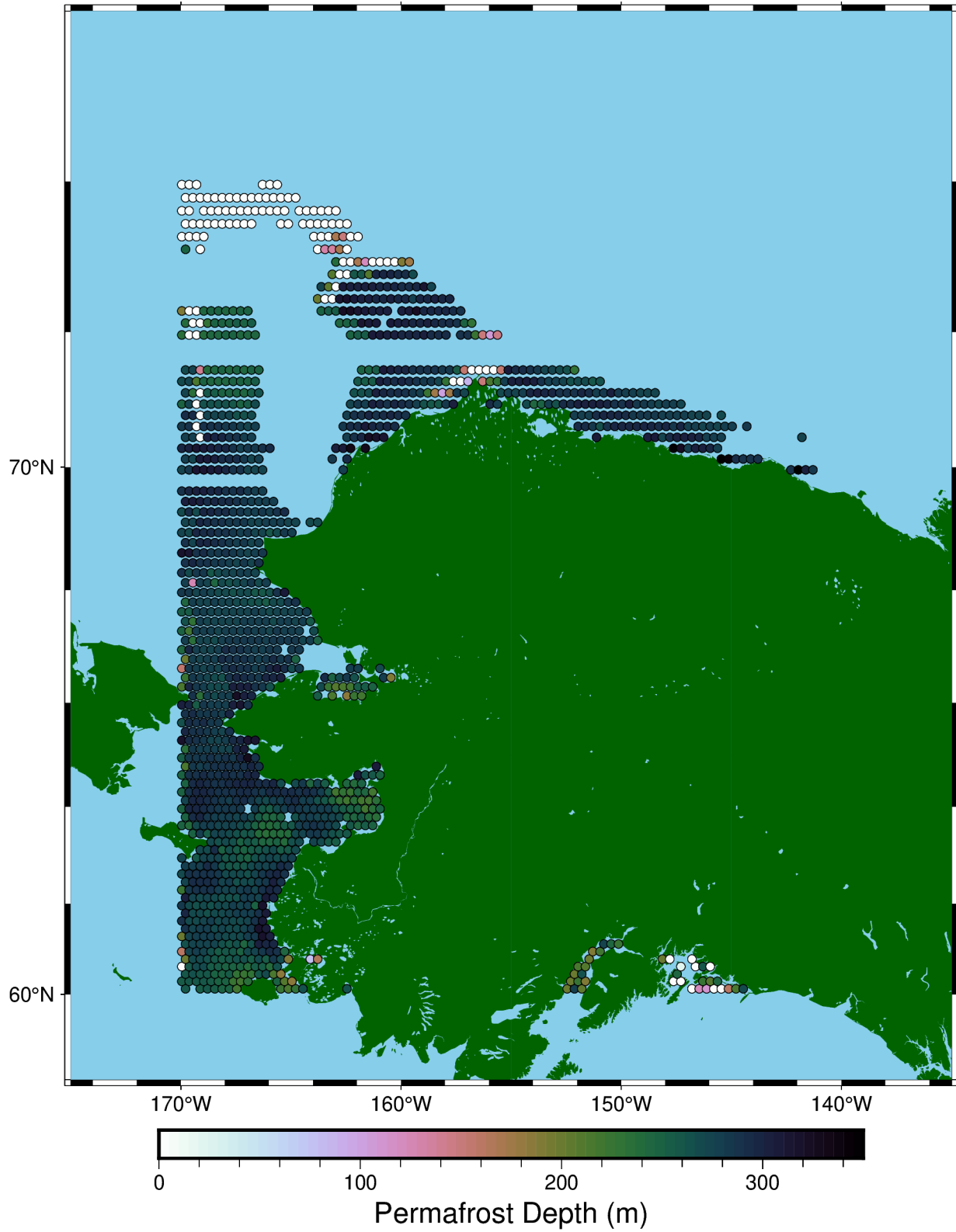


Figure 19. Permafrost depth (distance to the bottom of the permafrost layer) at the Last Glacial Maximum (LGM, 18 kA BP). Locations within the computational domain that are missing data did not complete (simulations failed).

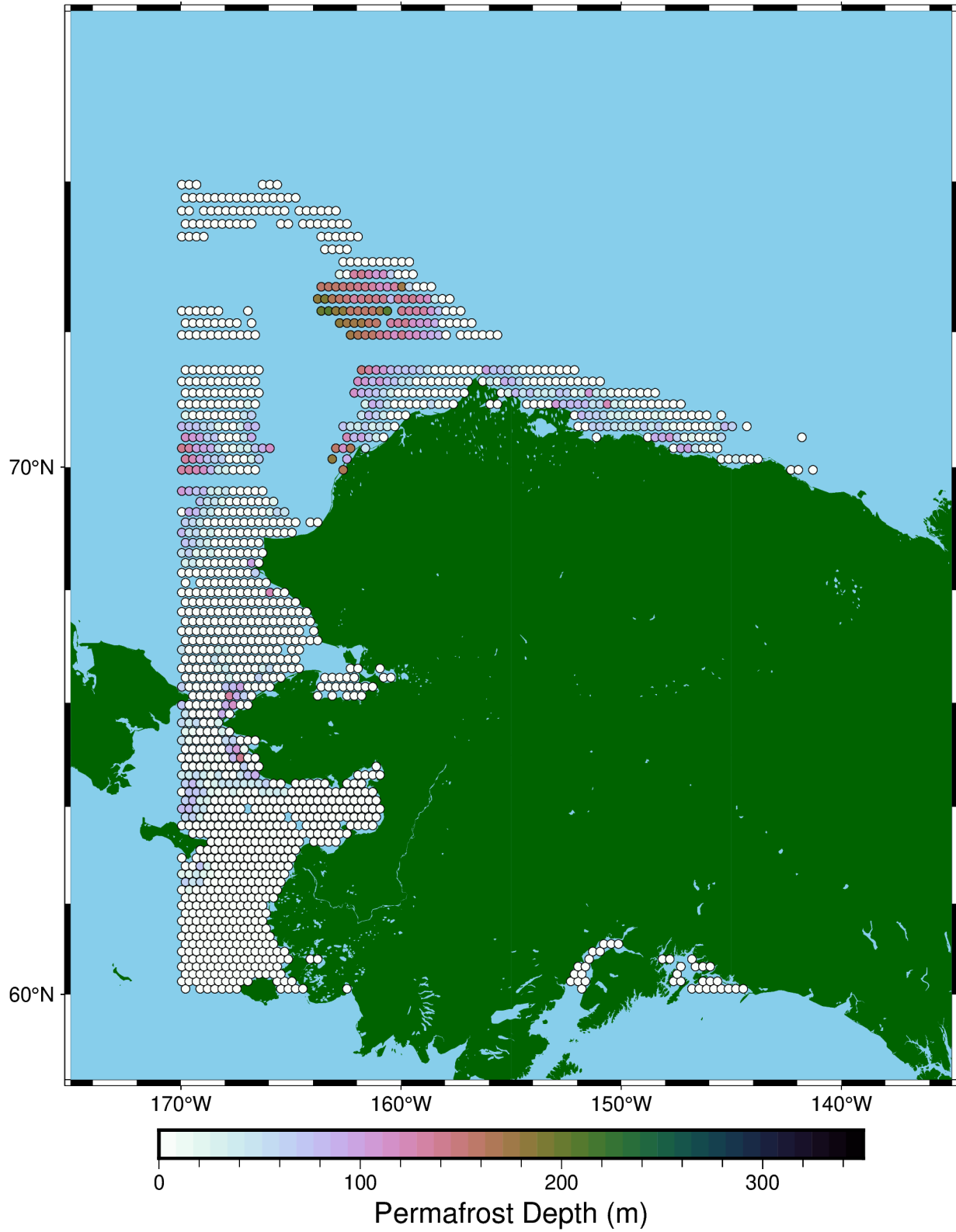


Figure 20. Permafrost depth (distance to the bottom of the permafrost layer) at the present day (0 kA BP). Locations within the computational domain that are missing data did not complete (simulations failed).

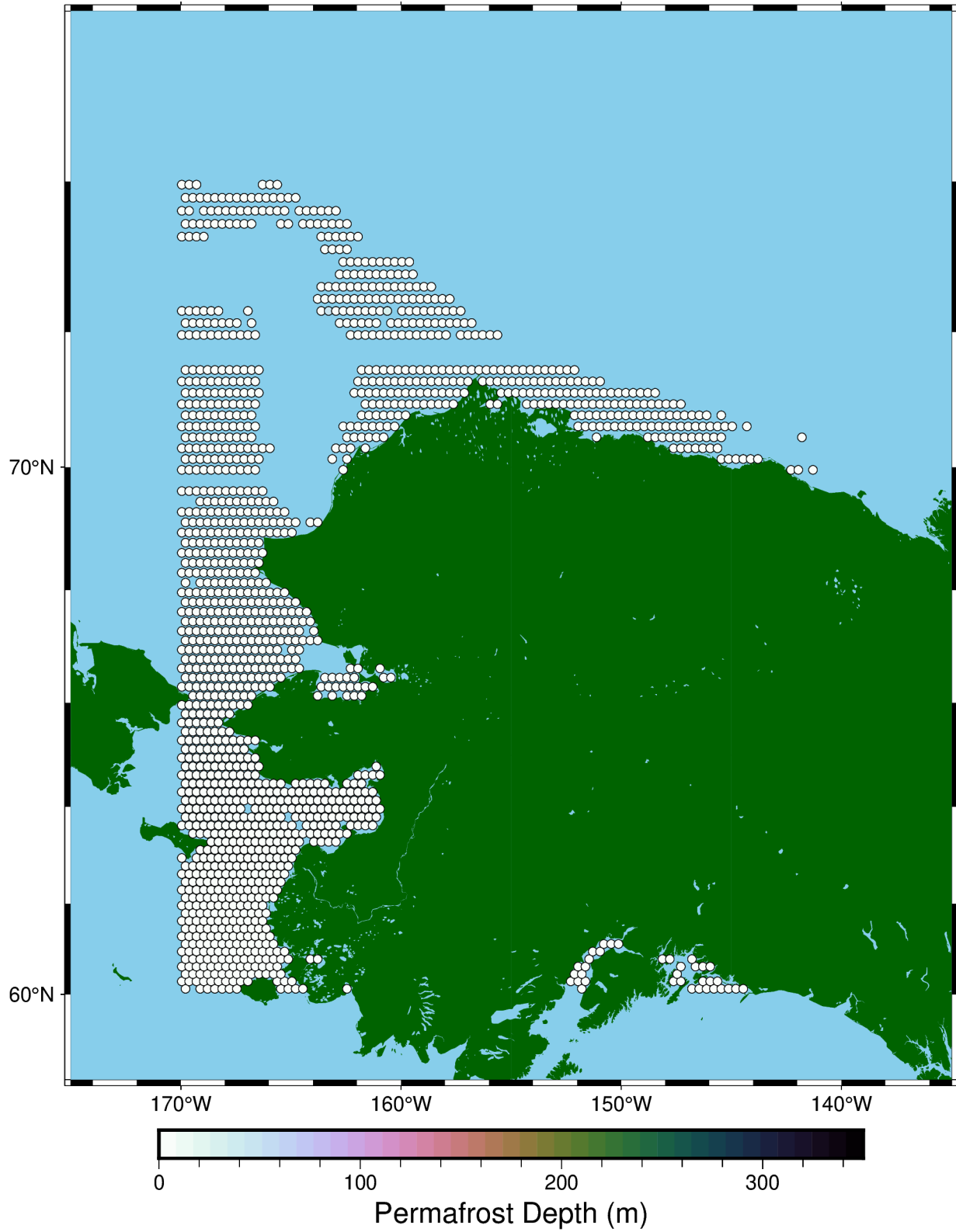


Figure 21. Permafrost depth (distance to the bottom of the permafrost layer) at 10 ka in the future. Locations within the computational domain that are missing data did not complete (simulations failed).

4.3.3. Methane Gas Distribution and Seabed-Water Flux

The average gas flux (kg yr^{-1}) at the 1 m^2 sediment-water column boundary, representing methane gas emissions from the sediments into the water column, for the LGM, present day, and 10 kA in the future are shown in **Figure 22**, **Figure 23**, and **Figure 24**. In these figures, bubble size and color are indicative of the magnitude of gas flux. We note that some simulation locations do not appear as their flux is either negligible, or the simulations failed to complete. A comparison against **Figure 18** is necessary to determine which of the two reasons for missing data applies, as presented. Once the numerical challenges are sorted out, failing simulations should no longer be an issue. For more detail, in **Figure 25**, **Figure 26**, and **Figure 27**, we also present the minimum and maximum values of gas flux at the LGM, present day, and 10 kA in the future.

While we must emphasize that these results are only notional and should not be interpreted as scientific evidence, certain trends can be identified. Little to no gas emissions from the seabed are predicted at the LGM, when the majority of the computational domain was exposed to the atmosphere. However, appreciable gas flux is seen in the northwest portion of the domain, where the seabed is deep enough that it is still submerged. At the present day, a significant increase in gas phase methane emissions is predicted. Emission rate magnitude is expected to increase with increasing TOC and sedimentation rate, as well as the state of permafrost degradation and the increase in sediment permeability associated with thaw. Interestingly, the model predicts a decrease in gas emissions 10 kA in the future relative to the present day. Although the methanogenesis parameters were held constant throughout the simulation period and gas mass increased over time, increasing sea level increases the pressure, and drives gaseous methane into solution, limiting gas phase saturation and thus mobility.

In **Figure 28**, the methane gas flux at three profile locations is shown at several depths within the 1D sediment columns, highlighting the capability of the modeling workflow to track and quantify gas movement within the simulations. The locations for each profile are as follows - Profile 220481: 166°W & 64°N , Profile 324121: 169°W & 66°N , and Profile 592139: 155°W & 71°N . It is immediately noticeable that the majority of mobile gas exists near the sediment surface, and not at depth. This is a result of the short time span of the simulations when methanogenesis occurred. Because the source of gas is highest near the top of the sediment column and decays towards depth, the top of the sediment column is expected to form gas before the rest of it does. If these simulations were allowed to run for time spans more representative of the geologic record, gas formation would also be expected at deeper depths.

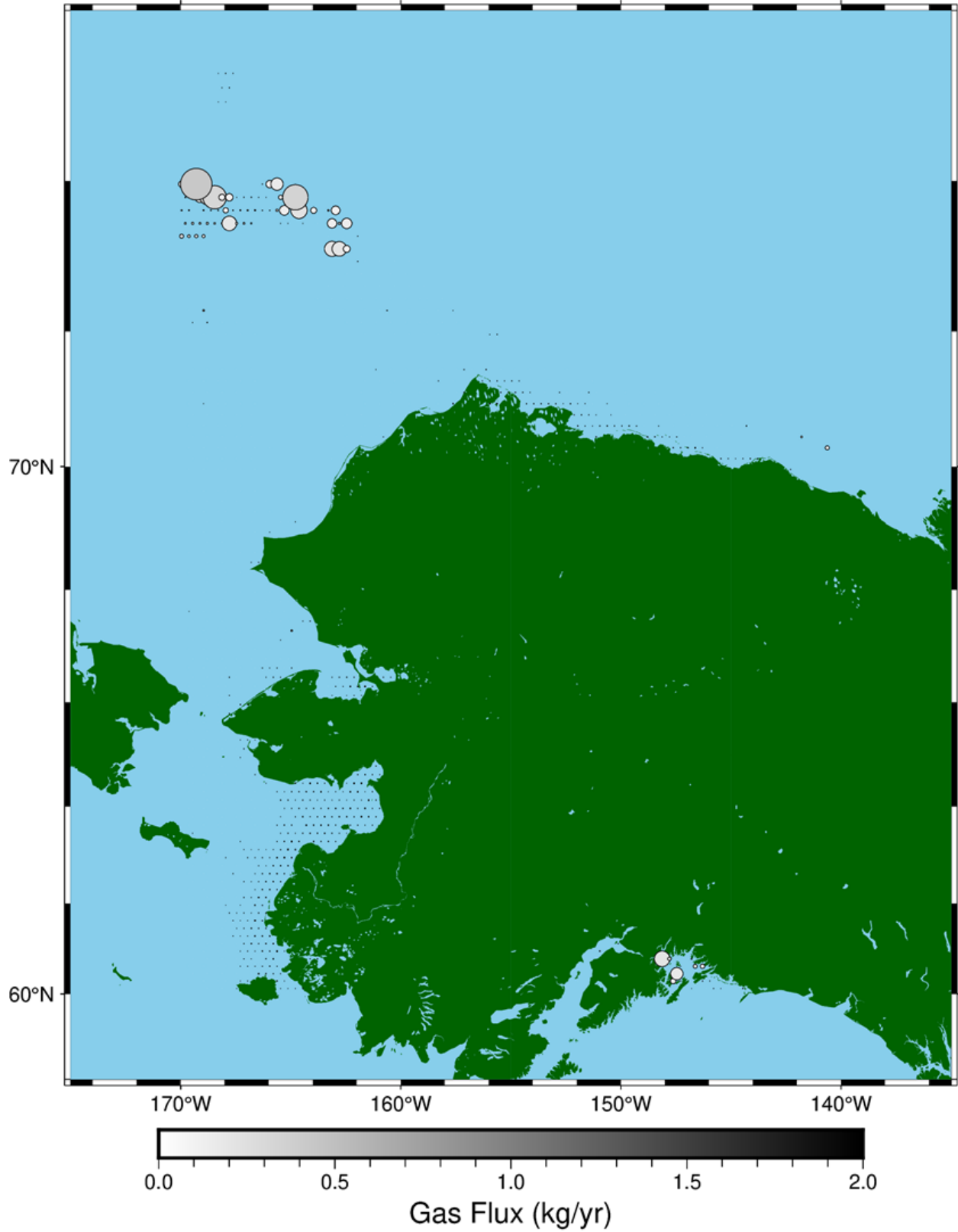


Figure 22. Average methane gas flux at the 1 m² sediment surface at the LGM (18 ka BP). Bubble size & color are indicative of magnitude. Some locations do not appear as their flux is negligible or the simulations failed.

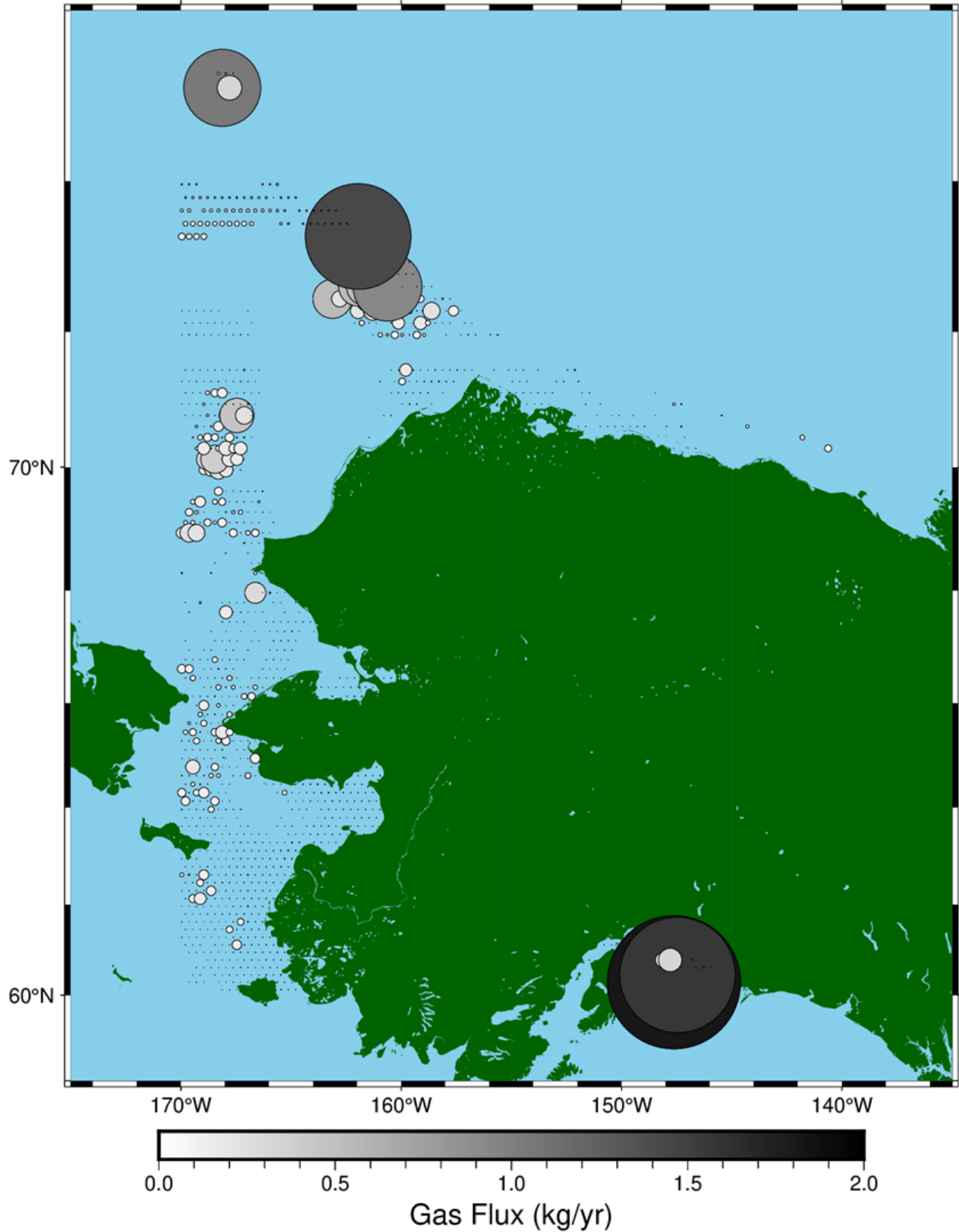


Figure 23. Average methane gas flux at the 1 m² sediment surface at the present day (0 kA BP). Bubble size & color are indicative of magnitude. Some locations do not appear as their flux is negligible or the simulations failed.

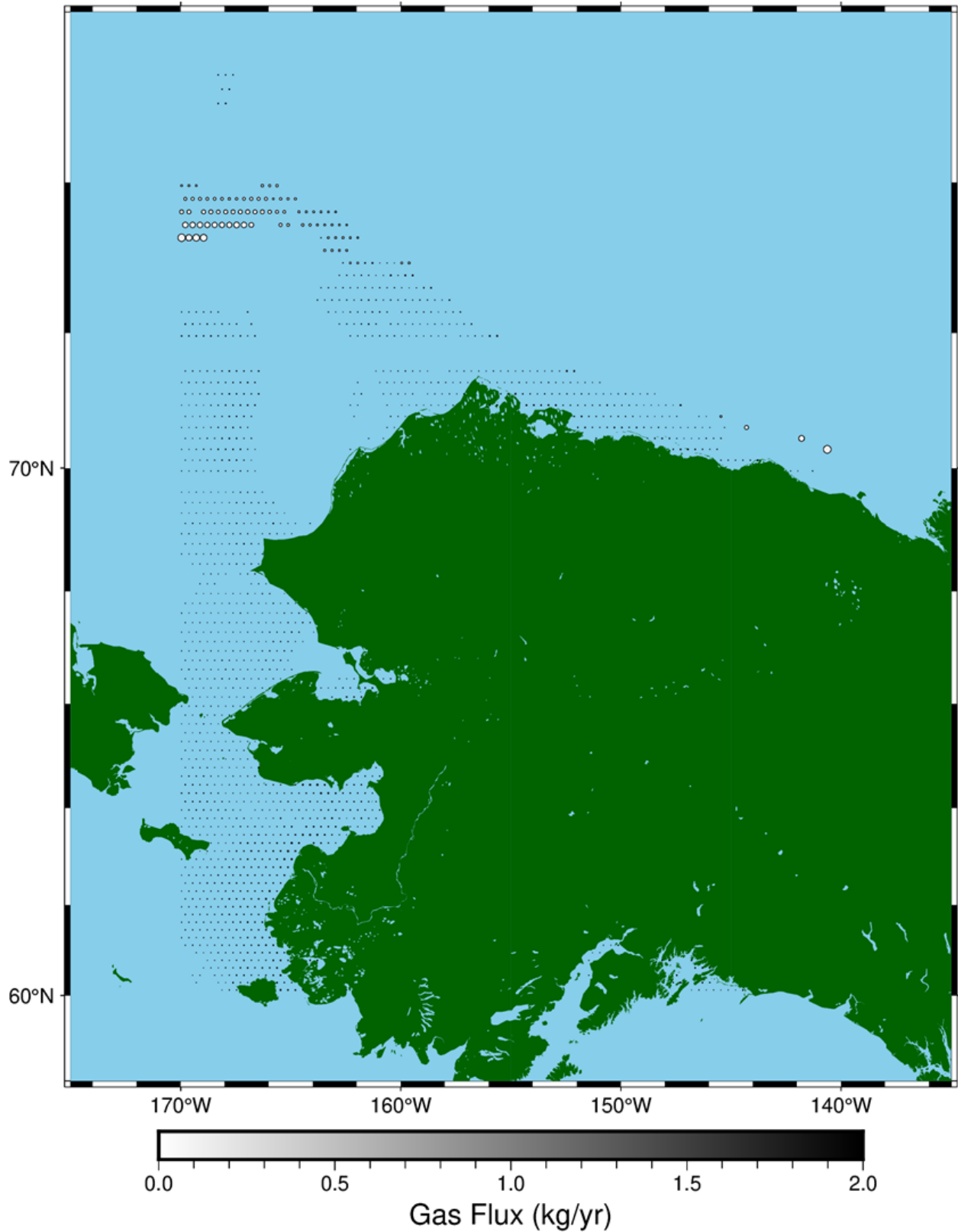


Figure 24. Average methane gas flux at the 1 m² sediment surface 10 kA in the future. Bubble size & color are indicative of magnitude. Some locations do not appear as their flux is negligible or the simulations failed.

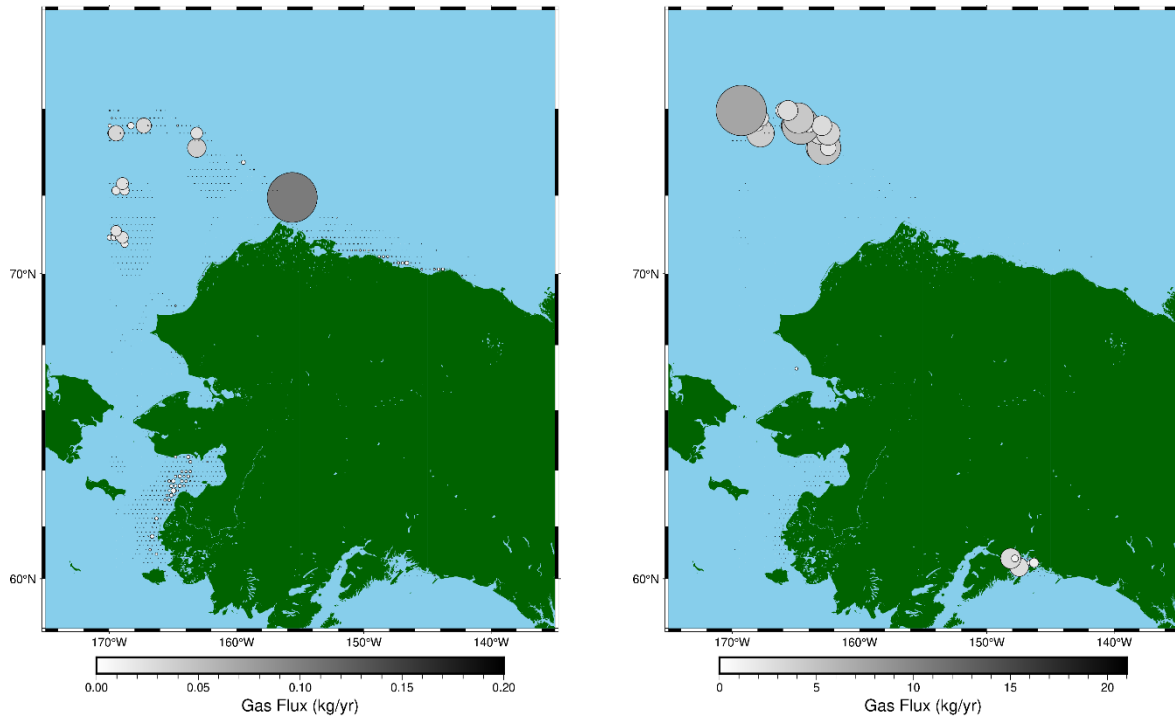


Figure 25. Minimum (left) and maximum (right) methane gas flux at the 1 m² sediment surface at the LGM (18 ka BP). Bubble size & color are indicative of magnitude. Some locations do not appear as their flux is negligible or the simulations failed.

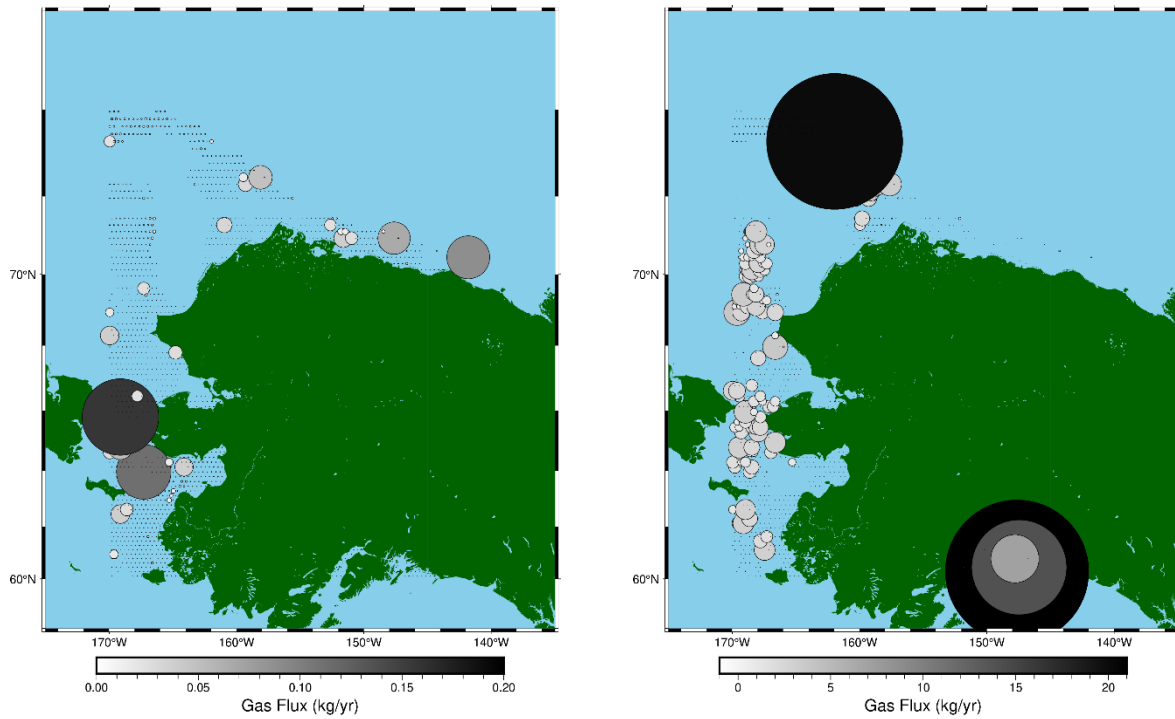
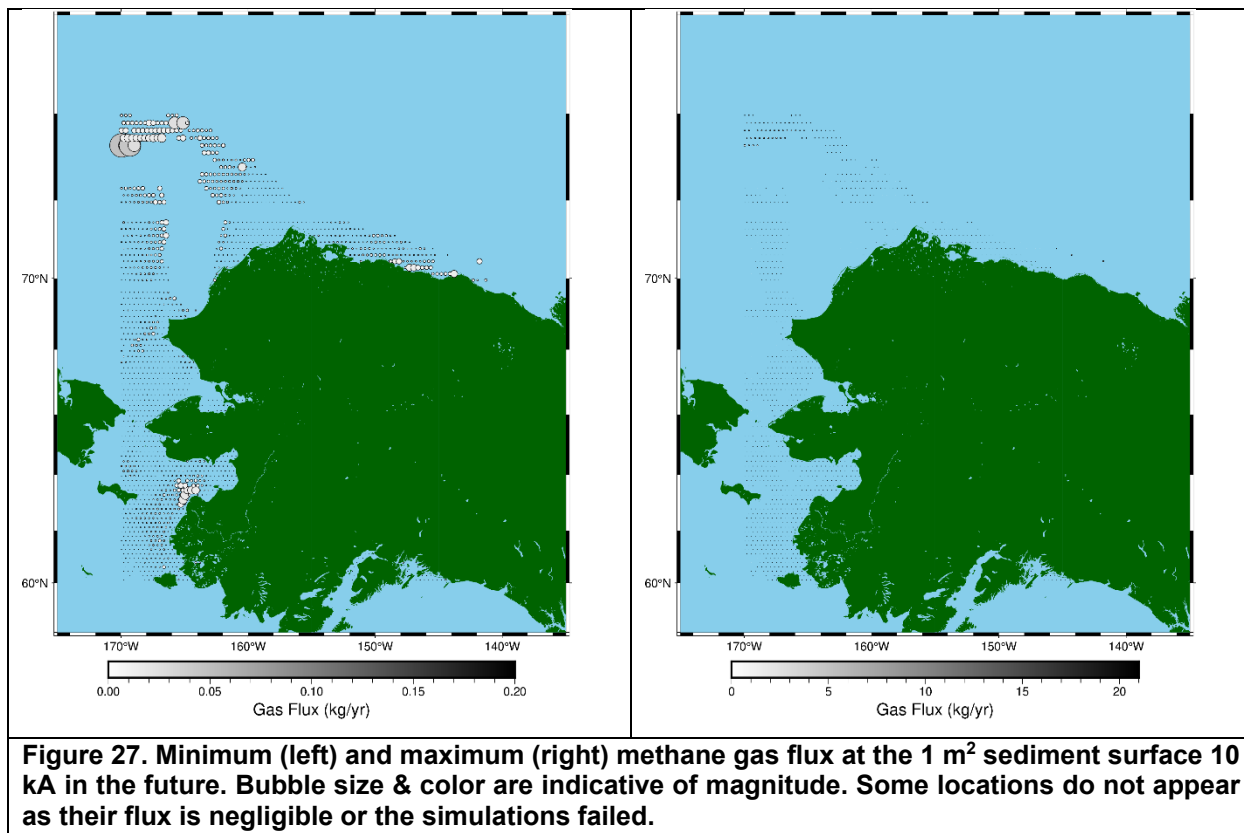
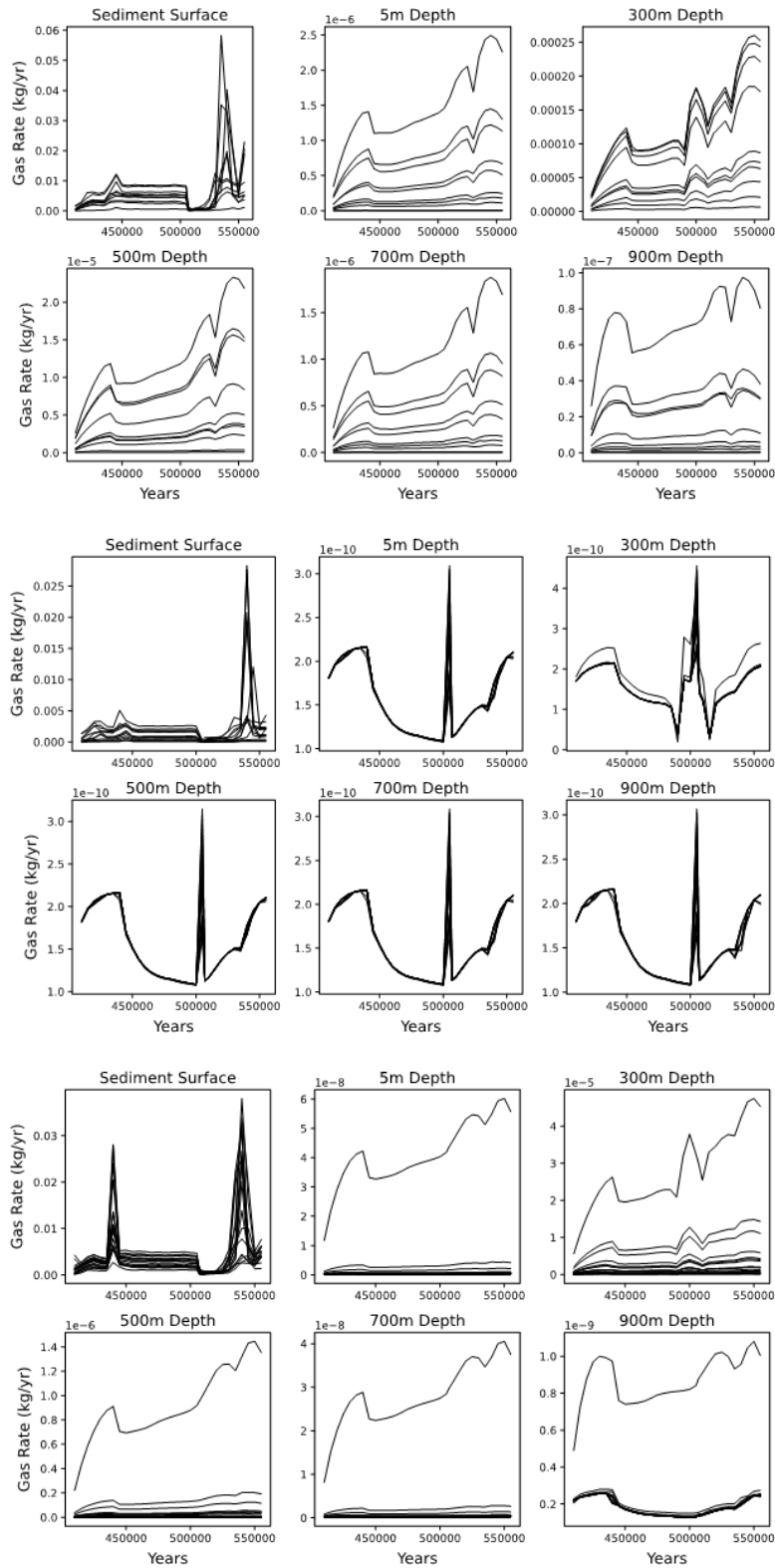


Figure 26. Minimum (left) and maximum (right) methane gas flux at the 1 m² sediment surface at the present day (0 ka BP). Bubble size & color are indicative of magnitude. Some locations do not appear as their flux is negligible or the simulations failed.



In this section, we have presented the capability that this hybrid workflow can provide for ensemble modeling of the thermodynamics and methanogenesis on the Arctic continental shelf, resulting in probabilistic predictions for methane gas emissions from the seabed to the sediment column. We must emphasize, again, that the demonstration presented here should not be interpreted as scientific evidence, but rather, a proof of concept of the workflow capability. Much more improvements are necessary in the PFLOTTRAN HYDRATE Mode process model, and the selection of model input parameters, before results can be used for scientific purposes.



Profile 220481

Profile 324121

Profile 592139

Figure 28. Methane gas flux (per 1 m²) over simulation time at different depths for Profiles 220481, 324121, and 592139. The LGM corresponds to 507 kA.

5. APPLICATION TO GLOBAL CLIMATE MODELING

Global climate models allow us to better understand how the Earth's climate will evolve as an interconnected, complex system. Models such as the DOE's E³SM, an advanced Earth system model, do this through the coupled simulation of land, ocean, and atmospheric physics, and require the world's most powerful supercomputers to run. One of the stated goals of this work is to use the developed hybrid modeling workflow to provide a reliable estimate of a time-varying source term for marine GHG flux to global oceanographic circulation and Earth system models. In this section, we briefly explore the history and current status of DOE's E³SM, and what it may require before E³SM can utilize such a source term.

E³SM Version 1 (E³SMv1) was released to the broader scientific community in April 2018, and E³SM Version 2 (E³SMv2) was tagged in September 2021. At the time of this writing, E³SM is still in its second version. E³SM Science is organized around three science drivers with simulation campaigns focused on water cycle, biogeochemistry, and cryosphere research. Each campaign employs unique model configurations optimized to address the associated research challenges by varying their choice of biogeochemical species. Marine GHG flux most closely fits into E³SM's biogeochemistry science campaign.

In the current E3SMv2, the model prescribes a global methane mole fraction in the atmosphere. For historical runs the methane input data are from years 0 to 2014 and are obtained from [103], which provides consolidated data sets of historical atmospheric (volume) 31 mixing ratios of 43 greenhouse gases specifically for the purpose of climate model runs (see **Figure 29**). The future runs (2014-2500) use data from the same group, but the methane data are modeled with MAGICC7.0 (see **Figure 30**). MAGICC stands for 'Model for the Assessment of Greenhouse Gas Induced Climate Change'. It is a prime reduced-complexity model, often used by IPCC, for key scientific publications and by a number of Integrated Assessment Models, according to the MAGICC website.

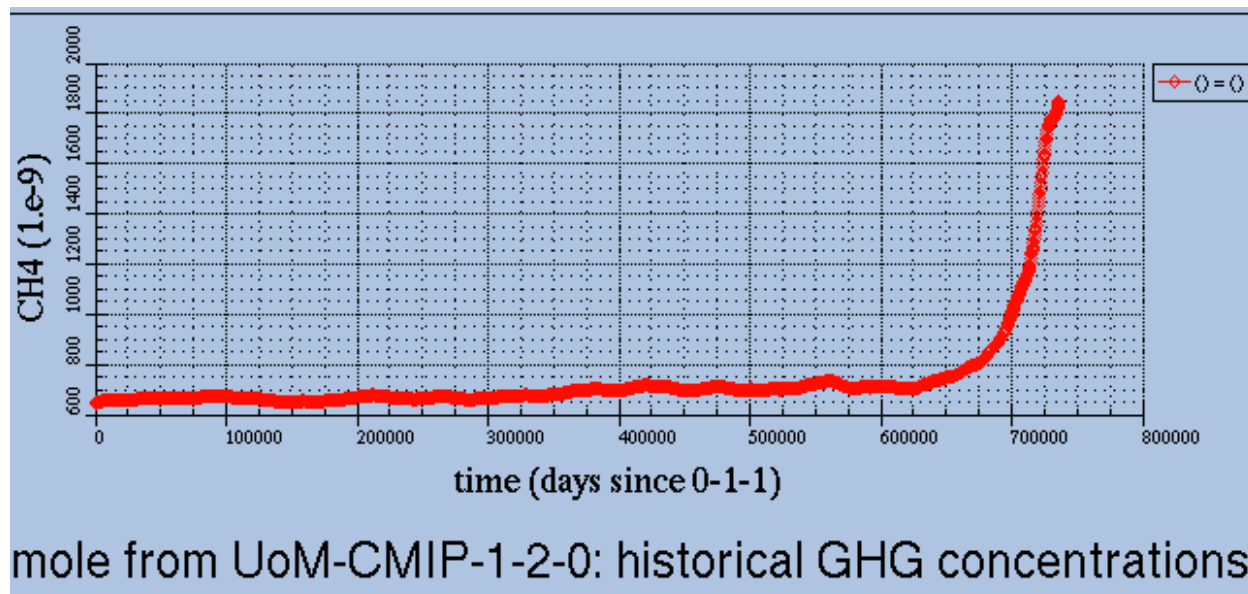


Figure 29. Atmospheric methane mole fraction prescribed for historical runs in E³SMv2.

Of particular emphasis is the current version of E³SM *prescribes* a methane mole fraction in the atmosphere, rather than allowing interaction between the atmosphere and land or marine sources of methane. Furthermore, the methane mole fraction in the atmosphere is applied homogenously, meaning the amount of methane does not vary in space, only in time.

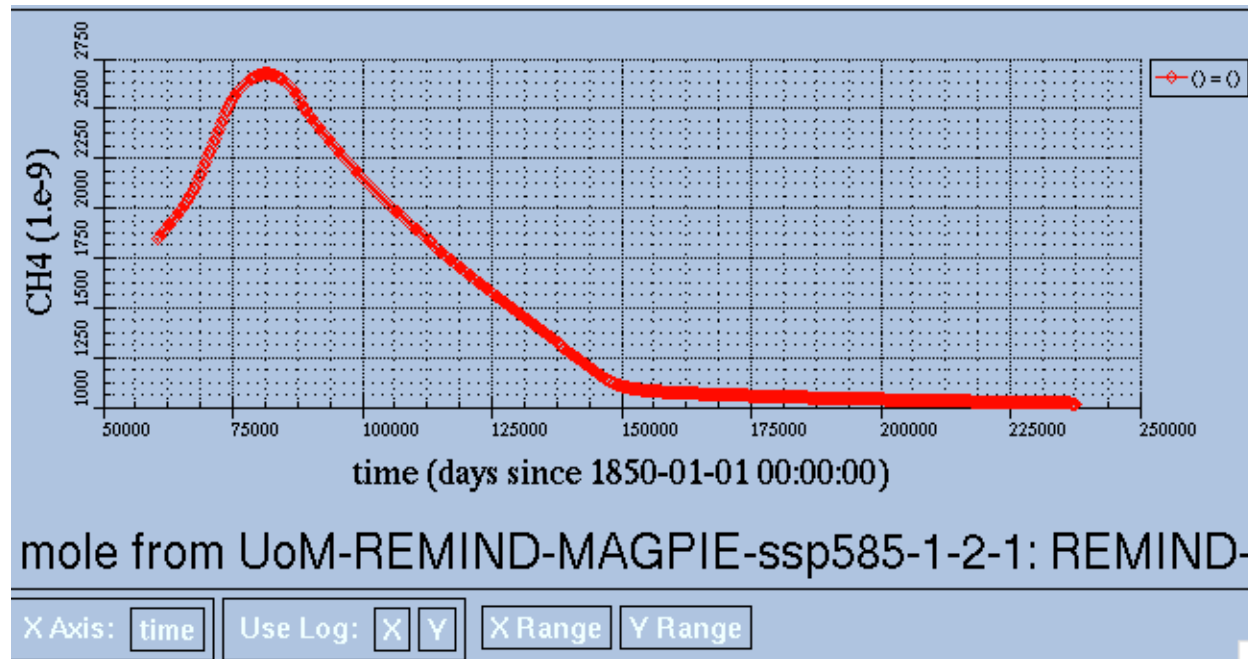


Figure 30. Atmospheric methane mole fraction prescribed for future runs in E3SMv2.

Until now, biogeochemistry configurations of E³SM have only included CO₂ as an interactively modeled GHG in the atmosphere. However, development is underway on E³SM Version 3 (E³SMv3), in which researchers have proposed to work on several critical improvements for the next phase, focusing on methane. These include the capability to: (i) simulate methane based on surface emissions and climate feedbacks, and (ii) simulate methane gas in the stratosphere, which is critical to simulating the climate impact of its evolving emissions. This work is projected to occur under Phase III of the project, which will start in FY23 and go through FY25. Only after this work is completed, will E³SM be configured such that a spatially and temporally varying source of marine GHG (particularly methane) can be of use to climate projects.

6. CONCLUSIONS AND FUTURE WORK

This report summarizes our group's effort at developing a numerical modeling framework designed to produce a first-of-its-kind estimate of Arctic methane gas releases from the marine sediments to the water column, and potentially to the atmosphere, where positive climate feedback may occur.

Newly developed modeling capability supported by the Laboratory Directed Research and Development (LDRD) program at Sandia National Laboratories now gives us the ability to probabilistically map gas distribution and quantity in the seabed by using a hybrid approach of geospatial machine learning, and predictive numerical thermodynamic ensemble modeling. The novelty in this approach is its ability to produce maps of useful data in regions that are only sparsely sampled, a common challenge in the Arctic, and a major obstacle to progress in the past. We have summarized the workflow developed in this study, and provided details on each software component: GPSM, Dakota, and PFLOTRAN HYDRATE Mode.

By applying this model to the circum-Arctic continental shelves and integrating the flux of free gas from *in situ* methanogenesis and dissociating gas hydrates from the sediment column under climate forcing, we can provide the most reliable estimate of a spatially and temporally varying source term for greenhouse gas flux that can be used by global oceanographic circulation and Earth system models (such as DOE's E³SM). We have presented a demonstration of this workflow for a site which considers Alaska's continental shelf. In this demonstration, we've made several simplifications to the input parameters, highlighting where improvements should be made in future work. We also note that debugging and fine-tuning PFLOTRAN HYDRATE Mode performance is ongoing. Nevertheless, we have demonstrated the capabilities of the hybrid modeling workflow, including its ability to probabilistically map submarine permafrost, methane hydrate, and methane gas distribution and flux over historical time periods through the future.

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APPENDIX A. PFLOTRAN HYDRATE MODE INPUT DECKS

A.1. hydrate_thermal.in (Section 2.3.5.1)

Note this input deck is displayed in two-column format for space efficiency.

<pre> #Description: Gas hydrate example: 1D Radial Thermal Dissociation SIMULATION SIMULATION_TYPE SUBSURFACE PROCESS_MODELS SUBSURFACE_FLOW flow MODE HYDRATE OPTIONS FIX_UPWIND_DIRECTION / / / END SUBSURFACE HYDRATE HYDRATE_PHASE_BOUNDARY MORIDIS_SIMPLE EFFECTIVE_SAT_SCALING ADJUST_SOLUBILITY_WITHIN_GHSZ GT_3PHASE WITH_GIBBS_THOMSON # NO_PC END NUMERICAL_METHODS flow NEWTON_SOLVER USE_INFINITY_NORM_CONVERGENCE CENTRAL_DIFFERENCE_JACOBIAN END LINEAR_SOLVER FLOW SOLVER DIRECT END END #===== discretization ===== GRID TYPE STRUCTURED CYLINDRICAL NXYZ 1500 1 1 DXYZ 1000@0.02d0 \ 0.020000000000000000\ 0.0202600212511425\ 0.0205234230548373\ 0.0207902493618595\ 0.0210605446943913\ 0.0213343541534502\ 0.0216117234264150\ 0.0218926987946491\ 0.0221773271412226\ 0.0224656559587354\ 0.0227577333572418\ 0.0230536080722776\ 0.0233533294729928\ </pre>	<pre> 0.0236569475703883\ 0.0239645130256615\ 0.0242760771586591\ 0.0245916919564404\ 0.0249114100819516\ 0.0252352848828133\ 0.0255633704002216\ 0.0258957213779658\ 0.0262323932715626\ 0.0265734422575092\ 0.0269189252426573\ 0.0272688998737076\ 0.0276234245468296\ 0.0279825584174049\ 0.0283463614098980\ 0.0287148942278549\ 0.0290882183640325\ 0.0294663961106586\ 0.0298494905698262\ 0.0302375656640228\ 0.0306306861467960\ 0.0310289176135581\ 0.0314323265125318\ 0.0318409801558372\ 0.0322549467307234\ 0.0326742953109462\ 0.0330990958682938\ 0.0335294192842617\ 0.0339653373618804\ 0.0344069228379691\ 0.0348542493939071\ 0.0353073916706589\ 0.0357664252784980\ 0.0362314268109885\ 0.0367024738574920\ 0.0371796450161145\ 0.0376630199068207\ 0.0381526791847194\ 0.0386487045535219\ 0.0391511787791740\ 0.0396601857036672\ 0.0401758102590278\ 0.0406981384814885\ 0.0412272575258448\ 0.0417632556799970\ 0.0423062223796818\ 0.0428562482233957\ 0.0434134249875117\ 0.0439778456415934\ 0.0445496043639073\ 0.0451287965571376\ 0.0457155188643046\ 0.0463098691848909\ 0.0469119466911759\ 0.0475218518447843\ 0.0481396864134487\ 0.0487655534879903\ </pre>
--	--

<p> 0.0493995574995205\ 0.0500418042368660\ 0.0506924008642208\ 0.0513514559390274\ 0.0520190794300901\ 0.0526953827359247\ 0.0533804787033461\ 0.0540744816462975\ 0.0547775073649251\ 0.0554896731648998\ 0.0562110978769910\ 0.0569419018768944\ 0.0576822071053175\ 0.0584321370883268\ 0.0591918169579586\ 0.0599613734730988\ 0.0607409350406337\ 0.0615306317368752\ 0.0623305953292656\ 0.0631409592983642\ 0.0639618588601191\ 0.0647934309884295\ 0.0656358144380007\ 0.0664891497674970\ 0.0673535793629942\ 0.0682292474617387\ 0.0691163001762143\ 0.0700148855185222\ 0.0709251534250784\ 0.0718472557816315\ 0.0727813464486062\ 0.0737275812867762\ 0.0746861181832711\ 0.0756571170779206\ 0.0766407399899423\ 0.0776371510449758\ 0.0786465165024684\ 0.0796690047834169\ 0.0807047864984699\ 0.0817540344763959\ 0.0828169237929209\ 0.0838936317999411\ 0.0849843381551166\ 0.0860892248518470\ 0.0872084762496402\ 0.0883422791048732\ 0.0894908226019546\ 0.0906542983848912\ 0.0918329005892653\ 0.0930268258746285\ 0.0942362734573153\ 0.0954614451436841\ 0.0967025453637906\ 0.0979597812054984\ 0.0992333624490332\ 0.100523501601987\ 0.101830413934776\ 0.103154317516559\ 0.104495433251630\ 0.105853984916268\ 0.107230199196085\ 0.108624305723846\ 0.110036537117786\ 0.111467129020424 </p>

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12.6000000000000

1@1.d0
1@1.d0
/
END

#===== fluid properties
=====
FLUID_PROPERTY
PHASE LIQUID
DIFFUSION_COEFFICIENT 1.d-9
END

FLUID_PROPERTY
PHASE GAS
DIFFUSION_COEFFICIENT 2.d-5
END

```

```

#===== more fluid properties
=====
EOS WATER
DENSITY IF97
ENTHALPY IF97
STEAM_DENSITY IF97
STEAM_ENTHALPY IF97
END

EOS GAS
DENSITY IDEAL # PR_METHANE
ENTHALPY IDEAL_METHANE
VISCOSITY DEFAULT #CONSTANT 1.1d-5 Pa-s #DEFAULT
END

#===== material properties
=====
MATERIAL_PROPERTY soil1
ID 1000
CHARACTERISTIC_CURVES default
POROSITY 0.3
TORTUOSITY 1.0
ROCK_DENSITY 2600.d0
THERMAL_CONDUCTIVITY_DRY 2.d0 #W/m-C
THERMAL_CONDUCTIVITY_WET 2.18d0 #W/m-C
HEAT_CAPACITY 1000 J/kg-C
PERMEABILITY
PERM_ISO 1.d-12
/
#SOIL_REFERENCE_PRESSURE INITIAL_PRESSURE
#POROSITY_COMPRESSIBILITY 1.d-9
END

#===== characteristic curves
=====
CHARACTERISTIC_CURVES default
SATURATION_FUNCTION IGHCC2_COMP #VAN_GENUCHTEN
ALPHA 8.d-5
M 0.45
LIQUID_RESIDUAL_SATURATION 0.12d0
MAX_CAPILLARY_PRESSURE 1.d9 #0.d0
#CALCULATE_INTERFACIAL_TENSION
/
PERMEABILITY_FUNCTION IGHCC2_COMP_LIQ
#BURDINE_BC_LIQ
PHASE LIQUID
LAMBDA 3
LIQUID_RESIDUAL_SATURATION 0.12
/
PERMEABILITY_FUNCTION IGHCC2_COMP_GAS
#BURDINE_BC_GAS
PHASE GAS
LAMBDA 3
LIQUID_RESIDUAL_SATURATION 0.12
GAS_RESIDUAL_SATURATION 0.02
/
END

#===== output options
=====
OUTPUT

```

```

SNAPSHOT_FILE
TIMES d 2 5 10 20 30 60
FORMAT TECPLOT POINT
/
UNFILTER_NON_STATE_VARIABLES

VARIABLES
TEMPERATURE
LIQUID_PRESSURE
GAS_PRESSURE
LIQUID_SATURATION
GAS_SATURATION
HYDRATE_SATURATION
ICE_SATURATION
LIQUID_MASS_FRACTIONS
LIQUID_MOLE_FRACTIONS
GAS_MASS_FRACTIONS
LIQUID_DENSITY
LIQUID_DENSITY MOLAR
GAS_DENSITY
GAS_DENSITY MOLAR
PERMEABILITY
LIQUID_RELATIVE_PERMEABILITY
GAS_RELATIVE_PERMEABILITY
POROSITY
#EFFECTIVE_POROSITY
/
END

#===== times
=====
TIME
FINAL_TIME 60 d
INITIAL_TIMESTEP_SIZE 1.d-6 h
#MAXIMUM_TIMESTEP_SIZE 8.d-3 d #8.d-3 d #1.d-2 d #5.d-3
d
MAXIMUM_TIMESTEP_SIZE 1.d-1 d
END

#===== regions
=====
REGION all
COORDINATES
0.d0 0.d0 0.d0
1.d3 1.d0 1.d0
/
END

REGION edge
FACE EAST
COORDINATES
1.0d3 0.d0 0.d0
1.0d3 1.d0 1.d0
/
END

REGION center
COORDINATES
0.d0 0.d0 0.d0
0.d0 1.d0 1.d0
/
END

```

```

#===== flow conditions
=====

FLOW_CONDITION initial
TYPE
GAS_PRESSURE DIRICHLET
HYDRATE_SATURATION DIRICHLET
TEMPERATURE DIRICHLET
/
GAS_PRESSURE 4.6d6
HYDRATE_SATURATION 5.d-1
TEMPERATURE 3.d0
END

FLOW_CONDITION edge
TYPE
GAS_PRESSURE DIRICHLET
HYDRATE_SATURATION DIRICHLET
TEMPERATURE DIRICHLET
/
GAS_PRESSURE 4.6d6
HYDRATE_SATURATION 5.d-1
TEMPERATURE 3.d0
END

FLOW_CONDITION center
SYNC_TIMESTEP_WITH_UPDATE
TYPE
RATE MASS_RATE
/
RATE 0.d0 0.d0 150.d0 g/s g/s W
END

#===== condition couplers
=====
# initial condition
INITIAL_CONDITION all
FLOW_CONDITION initial
REGION all
END

SOURCE_SINK center
FLOW_CONDITION center
REGION center
END

BOUNDARY_CONDITION edge
FLOW_CONDITION edge
REGION edge
END

#===== stratigraphy couplers
=====
STRATA
REGION all
MATERIAL soil1
END

#===== convergence criteria
=====
END_SUBSURFACE

```


A.2. hydrate_ice.in (Section 2.3.5.2)

```
#Description: Gas hydrate example
SIMULATION
  SIMULATION_TYPE SUBSURFACE
  PROCESS_MODELS
    SUBSURFACE_FLOW flow
    MODE HYDRATE
  /
  /
END

SUBSURFACE

NUMERICAL_METHODS flow
  NEWTON_SOLVER
  USE_INFINITY_NORM_CONVERGENCE
  PRES_ABS_UPDATE_INF_TOL 1.d-1
  PRES_REL_UPDATE_INF_TOL 1.d-4
  SAT_ABS_UPDATE_INF_TOL 1.d-6
  SAT_REL_UPDATE_INF_TOL 1.d-4
  RESIDUAL_ABS_INF_TOL 1.d-6
  RESIDUAL_SCALED_INF_TOL 1.d-7
  XMOL_ABS_UPDATE_INF_TOL 1.d-10
  XMOL_REL_UPDATE_INF_TOL 1.d-5
  CENTRAL_DIFFERENCE_JACOBIAN
  /
END

HYDRATE
  HYDRATE_PHASE_BOUNDARY MORIDIS
  WITH_GIBBS_THOMSON
  GT_3PHASE
  SCALE_PERM_BY_HYD_SAT
  EFFECTIVE_SAT_SCALING
  ADJUST_SOLUBILITY_WITHIN_GHSZ
  BC_REFERENCE_PRESSURE 1.d20
END

#===== discretization =====
GRID
  TYPE STRUCTURED
  NXYZ 1 1 1000
  DXYZ
    1@1.d1
    1@1.d1
    1000@1.d0
  /
END

#===== fluid properties =====
FLUID_PROPERTY
  PHASE LIQUID
  DIFFUSION_COEFFICIENT 1.d-9
```

END

FLUID_PROPERTY
PHASE GAS
DIFFUSION_COEFFICIENT 1.d-5
END

#===== more fluid properties =====

EOS WATER
DENSITY IF97
ENTHALPY IF97
STEAM_DENSITY IF97
STEAM_ENTHALPY IF97
END

EOS GAS
DENSITY DEFAULT
VISCOSITY DEFAULT
HENRYS_CONSTANT DEFAULT #CONSTANT 1.d11
END

#===== material properties =====

MATERIAL_PROPERTY marine_sediments
ID 10
CHARACTERISTIC_CURVES marine_sediments
POROSITY 0.4
TORTUOSITY_FUNCTION_OF_POROSITY 1.4 #Van Loon and Mibus 2015
ROCK_DENSITY 2700.d0
THERMAL_CONDUCTIVITY_DRY 4.0 W/m-C
THERMAL_CONDUCTIVITY_WET 1.0 W/m-C
HEAT_CAPACITY 830 J/kg-C
PERMEABILITY
PERM_ISO 1.d-15
/
END

#===== characteristic curves =====

CHARACTERISTIC_CURVES marine_sediments
SATURATION_FUNCTION VAN_GENUCHTEN
LOOP_INVARIANT
UNSATURATED_EXTENSION LCPC
ALPHA 1.d-4
M 0.2
LIQUID_RESIDUAL_SATURATION 0.1d0
MAX_CAPILLARY_PRESSURE 1.d8
/
PERMEABILITY_FUNCTION MUALEM_VG_LIQ
PHASE LIQUID
M 0.2
LIQUID_RESIDUAL_SATURATION 0.1d0
/
PERMEABILITY_FUNCTION MUALEM_VG_GAS
PHASE GAS

```
M 0.2
LIQUID_RESIDUAL_SATURATION 0.1d0
GAS_RESIDUAL_SATURATION 0.1d0
/
TEST
END
```

```
#===== output options =====
```

```
OUTPUT
SNAPSHOT_FILE
TIMES y 1 10 100 1000 10000
FORMAT TECPLOT POINT
/
UNFILTER_NON_STATE_VARIABLES
```

```
VARIABLES
TEMPERATURE
LIQUID_PRESSURE
GAS_PRESSURE
LIQUID_SATURATION
GAS_SATURATION
HYDRATE_SATURATION
ICE_SATURATION
LIQUID_MOLE_FRACTIONS
/
END
```

```
#===== times =====
```

```
TIME
FINAL_TIME 10000 y
INITIAL_TIMESTEP_SIZE 1.d0 h
MAXIMUM_TIMESTEP_SIZE 1.d2 y
```

```
END
```

```
#===== regions =====
```

```
REGION all
COORDINATES
0.d0 0.d0 0.d0
1.d1 1.d1 1000.d0
/
END
```

```
REGION top
FACE TOP
COORDINATES
0.0d0 0.d0 1000.d0
1.0d1 1.d1 1000.d0
/
END
```

```
REGION bottom
FACE BOTTOM
COORDINATES
0.d0 0.d0 0.d0
1.d1 1.d1 0.d0
```

```
/  
END
```

```
REGION methanogen  
COORDINATES  
0.d0 0.d0 700.d0  
1.d1 1.d1 800.d0
```

```
/  
END
```

```
##### flow conditions #####
```

```
FLOW_CONDITION initial  
TYPE  
LIQUID_PRESSURE HYDROSTATIC  
MOLE_FRACTION DIRICHLET  
TEMPERATURE DIRICHLET
```

```
/  
DATUM 0.d0 0.d0 0.d0  
LIQUID_PRESSURE 20.d6 #Pa  
MOLE_FRACTION 1.d-4  
TEMPERATURE 2.d0  
END
```

```
FLOW_CONDITION bcB  
TYPE  
LIQUID_PRESSURE DIRICHLET  
MOLE_FRACTION DIRICHLET  
TEMPERATURE DIRICHLET
```

```
/  
LIQUID_PRESSURE 20.d6 #Pa  
MOLE_FRACTION 1.d-4  
TEMPERATURE 2.d0  
END
```

```
FLOW_CONDITION bcT  
TYPE  
GAS_PRESSURE DIRICHLET_SEEPAGE  
RELATIVE_HUMIDITY DIRICHLET  
TEMPERATURE DIRICHLET
```

```
/  
GAS_PRESSURE 10.d6 #Pa  
RELATIVE_HUMIDITY 10  
TEMPERATURE -10.d0  
END
```

```
##### condition couplers #####
```

```
# initial condition  
INITIAL_CONDITION all  
FLOW_CONDITION initial  
REGION all  
END
```

```
BOUNDARY_CONDITION top  
FLOW_CONDITION bcT  
REGION top
```

END

BOUNDARY_CONDITION bottom

FLOW_CONDITION bcB

REGION bottom

END

#===== stratigraphy couplers =====

STRATA

REGION all

MATERIAL marine_sediments

END

END_SUBSURFACE

A.3. hydrate_ice_with_gas.in (Section 2.3.5.3)

```
#Description: Gas hydrate example
SIMULATION
  SIMULATION_TYPE SUBSURFACE
  PROCESS_MODELS
    SUBSURFACE_FLOW flow
    MODE HYDRATE
  /
  /
END

SUBSURFACE

NUMERICAL_METHODS flow
  NEWTON_SOLVER
  USE_INFINITY_NORM_CONVERGENCE
  PRES_ABS_UPDATE_INF_TOL 1.d-1
  PRES_REL_UPDATE_INF_TOL 1.d-4
  SAT_ABS_UPDATE_INF_TOL 1.d-6
  SAT_REL_UPDATE_INF_TOL 1.d-4
  RESIDUAL_ABS_INF_TOL 1.d-6
  RESIDUAL_SCALED_INF_TOL 1.d-7
  XMOL_ABS_UPDATE_INF_TOL 1.d-10
  XMOL_REL_UPDATE_INF_TOL 1.d-5
  CENTRAL_DIFFERENCE_JACOBIAN
  /
END

HYDRATE
  HYDRATE_PHASE_BOUNDARY MORIDIS
  WITH_GIBBS_THOMSON
  GT_3PHASE
  SCALE_PERM_BY_HYD_SAT
  EFFECTIVE_SAT_SCALING
  ADJUST_SOLUBILITY_WITHIN_GHSZ
END

#===== discretization =====
GRID
  TYPE STRUCTURED
  NXYZ 1 1 1000
  DXYZ
    1@1.d1
    1@1.d1
    1000@1.d0
  /
END

#===== fluid properties =====
FLUID_PROPERTY
  PHASE LIQUID
  DIFFUSION_COEFFICIENT 1.d-9
END

FLUID_PROPERTY
```

```
PHASE GAS
DIFFUSION_COEFFICIENT 1.d-5
END
```

```
#===== more fluid properties =====
```

```
EOS WATER
DENSITY IF97
ENTHALPY IF97
STEAM_DENSITY IF97
STEAM_ENTHALPY IF97
END
```

```
EOS GAS
DENSITY DEFAULT
VISCOSITY DEFAULT
HENRYS_CONSTANT DEFAULT #CONSTANT 1.d11
END
```

```
#===== material properties =====
```

```
MATERIAL_PROPERTY marine_sediments
ID 10
CHARACTERISTIC_CURVES marine_sediments
POROSITY 0.4
TORTUOSITY_FUNCTION_OF_POROSITY 1.4 #Van Loon and Mibus 2015
ROCK_DENSITY 2700.d0
THERMAL_CONDUCTIVITY_DRY 4.0 W/m-C
THERMAL_CONDUCTIVITY_WET 1.0 W/m-C
HEAT_CAPACITY 830 J/kg-C
PERMEABILITY
  PERM_ISO 1.d-15
/
END
```

```
#===== characteristic curves =====
```

```
CHARACTERISTIC_CURVES marine_sediments
SATURATION_FUNCTION VAN_GENUCHTEN
LOOP_INVARIANT
UNSATURATED_EXTENSION LCPC
ALPHA 1.d-1
M .05
LIQUID_RESIDUAL_SATURATION 0.1d0
MAX_CAPILLARY_PRESSURE 1.d8
/
PERMEABILITY_FUNCTION MUALEM_VG_LIQ
PHASE LIQUID
M 0.05
LIQUID_RESIDUAL_SATURATION 0.1d0
/
PERMEABILITY_FUNCTION MUALEM_VG_GAS
PHASE GAS
M 0.05
LIQUID_RESIDUAL_SATURATION 0.1d0
GAS_RESIDUAL_SATURATION 0.1d0
```

```
/
TEST
END
```

```
#===== output options =====
```

```
OUTPUT
SNAPSHOT_FILE
  TIMES y 1 10 100 1000 10000
  FORMAT TECPLOT POINT
/
UNFILTER_NON_STATE_VARIABLES
```

```
VARIABLES
TEMPERATURE
LIQUID_PRESSURE
GAS_PRESSURE
LIQUID_SATURATION
GAS_SATURATION
HYDRATE_SATURATION
ICE_SATURATION
LIQUID_MOLE_FRACTIONS
/
END
```

```
#===== times =====
```

```
TIME
FINAL_TIME 10000 y
INITIAL_TIMESTEP_SIZE 1.d0 h
MAXIMUM_TIMESTEP_SIZE 1.d2 y
END
```

```
#===== regions =====
```

```
REGION all
COORDINATES
  0.d0 0.d0 0.d0
  1.d1 1.d1 1000.d0
/
END
```

```
REGION top
FACE TOP
COORDINATES
  0.0d0 0.d0 1000.d0
  1.0d1 1.d1 1000.d0
/
END
```

```
REGION bottom
FACE BOTTOM
COORDINATES
  0.d0 0.d0 0.d0
  1.d1 1.d1 0.d0
/
END
```

```
REGION methanogen
```



```
COORDINATES
  0.d0 0.d0 700.d0
  1.d1 1.d1 800.d0
/
END
```

```
#===== flow conditions =====
```

```
FLOW_CONDITION initial
TYPE
  LIQUID_PRESSURE HYDROSTATIC
  MOLE_FRACTION DIRICHLET
  TEMPERATURE DIRICHLET
/
DATUM 0.d0 0.d0 0.d0
LIQUID_PRESSURE 20.d6
MOLE_FRACTION 5.d-4
TEMPERATURE 2.d0
END
```

```
FLOW_CONDITION bcB
TYPE
  LIQUID_PRESSURE DIRICHLET
  MOLE_FRACTION DIRICHLET
  TEMPERATURE DIRICHLET
/
LIQUID_PRESSURE 20.d6 #Pa
MOLE_FRACTION 5.d-4
TEMPERATURE 2.d0
END
```

```
FLOW_CONDITION bcT
TYPE
  GAS_PRESSURE DIRICHLET
  RELATIVE_HUMIDITY DIRICHLET
  TEMPERATURE DIRICHLET
/
GAS_PRESSURE 10.d6 #Pa
RELATIVE_HUMIDITY 10
TEMPERATURE -10.d0
END
```

```
FLOW_CONDITION methanogen
SYNC_TIMESTEP_WITH_UPDATE
TYPE
  RATE MASS_RATE
/
#RATE 0.d0 1.d-4 0.d0 g/s g/s MW
INTERPOLATION LINEAR
RATE FILE ./methanogen.gen
END
```

```
#===== condition couplers =====
```

```
# initial condition
INITIAL_CONDITION all
```

```

FLOW_CONDITION initial
REGION all
END

BOUNDARY_CONDITION top
FLOW_CONDITION bcT
REGION top
END

SOURCE_SINK methanogen
FLOW_CONDITION methanogen
REGION methanogen
END

BOUNDARY_CONDITION bottom
FLOW_CONDITION bcB
REGION bottom
END

#===== stratigraphy couplers =====
STRATA
REGION all
MATERIAL marine_sediments
END

END_SUBSURFACE

```

methanogen.gen

```

TIME_UNITS y
DATA_UNITS g/s g/s MW
0.000E+00 0.d0 1.d-4 0.d0
1.E+03 0.d0 1.d-4 0.d0
1.0E+03 0.d0 0.d0 0.d0
1.0E+05 0.d0 0.d0 0.d0

```

A.4. ice_run.in

```

WALLCLOCK_STOP 0.5 min
/
END

SIMULATION
SIMULATION_TYPE SUBSURFACE
PROCESS_MODELS
SUBSURFACE_FLOW flow
MODE HYDRATE
OPTIONS
/
/
CHECKPOINT

```

```

TIMES y 135000 270000 405000 510000 645000
FORMAT HDF5
TIME_UNITS y
/
END

SUBSURFACE
  NUMERICAL_METHODS flow
  NEWTON_SOLVER
  USE_INFINITY_NORM_CONVERGENCE
  SAT_ABS_UPDATE_INF_TOL 1.d-7
  SAT_REL_UPDATE_INF_TOL 1.d-8
/
  TIMESTEPPER
  PRESSURE_CHANGE_GOVERNOR 1.0d5
  TEMPERATURE_CHANGE_GOVERNOR 1.0d0
/
END

GRID
  TYPE STRUCTURED CARTESIAN
  NXYZ 1 1 1000
  DXYZ
    1@1.d0
    1@1.d0
    1000@1.d0
/
  GRAVITY 0.d0 0.d0 -9.8d0
END

DATASET porosity
  FILENAME ../phi_k_dataset.h5
  HDF5_DATASET_NAME porosity
END

REGION all
  COORDINATES
    0.d0 0.d0 0.d0
    1.d0 1.d0 1000.d0
/
END

REGION top
  FACE TOP
  COORDINATES
    0.0d0 0.d0 1000.d0
    1.0d0 1.d0 1000.d0
/
END

REGION bottom
  FACE BOTTOM
  COORDINATES
    0.d0 0.d0 0.d0
    1.d0 1.d0 0.d0
/
END

END

HYDRATE
  HYDRATE_PHASE_BOUNDARY MORIDIS_SIMPLE
  WITH_GIBBS_THOMSON
  GT_3PHASE
  BC_REFERENCE_PRESSURE 1.d20
  SCALE_PERM_BY_HYD_SAT
  EFFECTIVE_SAT_SCALING
  ADJUST_SOLUBILITY_WITHIN_GHSZ

```

```

#WITH_SEDIMENTATION
SALINITY 0.035
#METHANOGENESIS
#NAME ss_methanogenesis
#ALPHA 0.0
#K_ALPHA 2241
#LAMBDA 1e-14
#V_SED 0.0
#SMT_DEPTH 15
#/
#PERM_SCALING_FUNCTION DAI_AND_SEOL
END

FLUID_PROPERTY
PHASE LIQUID
DIFFUSION_COEFFICIENT 1.d-9
END

FLUID_PROPERTY
PHASE GAS
DIFFUSION_COEFFICIENT 2.d-5
END

EOS WATER
DENSITY IF97
ENTHALPY IF97
STEAM_DENSITY IF97
STEAM_ENTHALPY IF97
END

EOS GAS
DENSITY PR_METHANE
ENTHALPY IDEAL_METHANE
VISCOSITY DEFAULT #CONSTANT 1.1d-5 Pa-s #DEFAULT
END

MATERIAL_PROPERTY marine_sediments
ID 1
CHARACTERISTIC_CURVES marine_sediments
POROSITY DATASET porosity
TORTUOSITY_FUNCTION_OF_POROSITY 1.4 #Van Loon and Mibus 2015
ROCK_DENSITY 2700.d0
THERMAL_CONDUCTIVITY_DRY 1.0 W/m-C
THERMAL_CONDUCTIVITY_WET 1.0 W/m-C
HEAT_CAPACITY 830 J/kg-C
PERMEABILITY
  PERM_ISO 1.d-15
/
END

CHARACTERISTIC_CURVES marine_sediments
SATURATION_FUNCTION VAN_GENUCHTEN
  LOOP_INVARIANT
  UNSATURATED_EXTENSION ECPC
  ALPHA 5.8d-4
  M 0.189
  LIQUID_RESIDUAL_SATURATION 0.1d0
  MAX_CAPILLARY_PRESSURE 1.d8 #roughly read from Jacinto et al. 2009
/
PERMEABILITY_FUNCTION MUALEM_VG_LIQ
  PHASE LIQUID
  M 0.189
  LIQUID_RESIDUAL_SATURATION 0.1d0
/
PERMEABILITY_FUNCTION MUALEM_VG_GAS
  PHASE GAS
  M 0.189
  LIQUID_RESIDUAL_SATURATION 0.1d0

```

```

GAS_RESIDUAL_SATURATION 0.15d0
/
END

STRATA
  REGION all
  MATERIAL marine_sediments
END

TIME
  FINAL_TIME 645000 y
  INITIAL_Timestep_SIZE 1.d-3 y
  MAXIMUM_Timestep_SIZE 250.0 y
END

OUTPUT
  SNAPSHOT_FILE
  PERIODIC_TIME 15000 y
  FORMAT TECPLOT POINT
  #NO_PRINT_INITIAL
/
#UNFILTER_NON_STATE_VARIABLES

VARIABLES
  GAS_SATURATION
  HYDRATE_SATURATION
  GAS_DENSITY
  POROSITY
  ICE_SATURATION
  TEMPERATURE
  THERMODYNAMIC_STATE
  LIQUID_SATURATION
  GAS_PRESSURE
  LIQUID_PRESSURE
  PERMEABILITY
/
END

EXTERNAL_FILE flowConditions.txt

# initial condition
INITIAL_CONDITION all
  FLOW_CONDITION initial
  REGION all
END

BOUNDARY_CONDITION top
  FLOW_CONDITION top
  REGION top
END

BOUNDARY_CONDITION bottom
  FLOW_CONDITION bottom
  REGION bottom
END

END_SUBSURFACE

```

A.5. methane_run.in

```
SIMULATION
```

```

SIMULATION_TYPE SUBSURFACE
PROCESS_MODELS
  SUBSURFACE_FLOW flow
  MODE HYDRATE
  OPTIONS
    RESTRICT_STATE_CHANGE
  /
  /
  /
RESTART
  FILENAME ../ice_run1/ice_run1-405000.0000y.h5
  /
END

SUBSURFACE
NUMERICAL_METHODS flow
NEWTON_SOLVER
  USE_INFINITY_NORM_CONVERGENCE

  SAT_ABS_UPDATE_INF_TOL 1.d-2
  SAT_REL_UPDATE_INF_TOL 1.d-2

  GAS_PRES_ABS_UPDATE_INF_TOL 1.d+4
  GAS_PRES_REL_UPDATE_INF_TOL 1.d-2

  LIQUID_PRES_ABS_UPDATE_INF_TOL 1.d+4
  LIQUID_PRES_REL_UPDATE_INF_TOL 1.d-2

  GAS_RESIDUAL_ABS_INF_TOL 1.d-8
  GAS_RESIDUAL_SCALED_INF_TOL 1.d-5

  LIQUID_RESIDUAL_ABS_INF_TOL 1.d-7
  LIQUID_RESIDUAL_SCALED_INF_TOL 1.d-4
  /
TIMESTEPPER
  PRESSURE_CHANGE_GOVERNOR 1.0d6
  SATURATION_CHANGE_GOVERNOR 0.05
  #TEMPERATURE_CHANGE_GOVERNOR 0.5d0
  /
END

WALLCLOCK_STOP 120 min

GRID
  TYPE STRUCTURED CARTESIAN
  NXYZ 1 1 1000
  DXYZ
    1@1.d0
    1@1.d0
    1000@1.d0
  /
  GRAVITY 0.d0 0.d0 -9.8d0
END

DATASET porosity
  FILENAME ../phi_k_dataset.h5
  HDF5_DATASET_NAME porosity
END

REGION all
  COORDINATES
    0.d0 0.d0 0.d0
    1.d0 1.d0 1000.d0
  /
END

REGION top

```

```

FACE TOP
COORDINATES
  0.0d0 0.0d0 1000.d0
  1.0d0 1.0d0 1000.d0
/
END

REGION bottom
FACE BOTTOM
COORDINATES
  0.d0 0.d0 0.d0
  1.d0 1.d0 0.d0
/

END

# positive flux is up
INTEGRAL_FLUX sediment_surface
FLUXES_OPTION POSITIVE_FLUXES_ONLY
PLANE
  0.d0 0.d0 1000.d0
  1.d0 1.d0 1000.d0
  0.d0 1.d0 1000.d0
/
END

# positive flux is up
#INTEGRAL_FLUX ice_surface
# CELL_IDS
# 327 328
# /
#END

# This will be dynamic, find ice > 0 and make CELL_IDS
INTEGRAL_FLUX 5m_depth
CELL_IDS
  322 323
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 300m
CELL_IDS
  700 701
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 325m
CELL_IDS
  675 676
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 350m
CELL_IDS
  650 651
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 375m
CELL_IDS

```

```
625 626
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 400m
CELL_IDS
600 601
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 425m
CELL_IDS
575 576
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 450m
CELL_IDS
550 551
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 475m
CELL_IDS
525 526
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 500m
CELL_IDS
500 501
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 525m
CELL_IDS
475 476
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 550m
CELL_IDS
450 451
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 575m
CELL_IDS
425 426
/
END
```



```
# positive flux is up
INTEGRAL_FLUX
  NAME 600m
  CELL_IDS
    400 401
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 625m
  CELL_IDS
    375 376
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 650m
  CELL_IDS
    350 351
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 675m
  CELL_IDS
    325 326
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 700m
  CELL_IDS
    300 301
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 725m
  CELL_IDS
    275 276
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 750m
  CELL_IDS
    250 251
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 775m
  CELL_IDS
    225 226
  /
END
```

```
# positive flux is up
INTEGRAL_FLUX
  NAME 800m
  CELL_IDS
```

```

200 201
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 825m
CELL_IDS
175 176
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 850m
CELL_IDS
150 151
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 875m
CELL_IDS
125 126
/
END

# positive flux is up
INTEGRAL_FLUX
NAME 900m
CELL_IDS
100 101
/
END

HYDRATE
HYDRATE_PHASE_BOUNDARY MORIDIS
WITH_GIBBS_THOMSON
GT_3PHASE
BC_REFERENCE_PRESSURE 1.d20
SCALE_PERM_BY_HYD_SAT
EFFECTIVE_SAT_SCALING
ADJUST_SOLUBILITY_WITHIN_GHSZ
WITH_SEDIMENTATION
SALINITY 0.035
METHANOGENESIS
NAME ss_methanogenesis
ALPHA 0.1597591273811524
K_ALPHA 2241
LAMBDA 1e-14
V_SED 3.303563591360281e-13
SMT_DEPTH 5.0
/
#PERM_SCALING_FUNCTION DAI_AND_SEOL
END

FLUID_PROPERTY
PHASE LIQUID
DIFFUSION_COEFFICIENT 1.d-9
END

FLUID_PROPERTY
PHASE GAS
DIFFUSION_COEFFICIENT 2.d-5
END

EOS WATER

```

```

DENSITY IF97
ENTHALPY IF97
STEAM_DENSITY IF97
STEAM_ENTHALPY IF97
END

EOS GAS
DENSITY PR METHANE
ENTHALPY IDEAL METHANE
VISCOSITY DEFAULT #CONSTANT 1.1d-5 Pa-s #DEFAULT
END

MATERIAL_PROPERTY marine_sediments
ID 1
CHARACTERISTIC_CURVES marine_sediments
POROSITY DATASET porosity
TORTUOSITY_FUNCTION_OF_POROSITY 1.4 #Van Loon and Mibus 2015
ROCK_DENSITY 2700.d0
THERMAL_CONDUCTIVITY_DRY 1.0 W/m-C
THERMAL_CONDUCTIVITY_WET 1.0 W/m-C
HEAT_CAPACITY 830 J/kg-C
PERMEABILITY
  PERM_ISO 5.d-14
/
END

CHARACTERISTIC_CURVES marine_sediments
PERMEABILITY_FUNCTION BURDINE_LINEAR_LIQ
  PHASE LIQUID
  LIQUID_RESIDUAL_SATURATION 0.2d0
/
PERMEABILITY_FUNCTION BURDINE_LINEAR_GAS
  PHASE GAS
  LIQUID_RESIDUAL_SATURATION 0.2d0
  GAS_RESIDUAL_SATURATION 0.10d0
/
SATURATION_FUNCTION VAN_GENUCHTEN
  LOOP_INVARIANT
  UNSATURATED_EXTENSION ECPC
  ALPHA 1.0e-5
  M 0.2
  LIQUID_RESIDUAL_SATURATION 0.1d0
  MAX_CAPILLARY_PRESSURE 5.d10
/
END

STRATA
  REGION all
  MATERIAL marine_sediments
END

TIME
  FINAL_TIME 555000.0 y
  INITIAL_TIMESTEP_SIZE 1.d-3 y
  MAXIMUM_TIMESTEP_SIZE 100.0 y

END

OUTPUT
  SNAPSHOT_FILE
    TIMES y 507000.0
    PERIODIC TIME 5000 y
    FORMAT TECPLOT POINT
    #NO_PRINT_INITIAL
  /
  OBSERVATION_FILE
    TIMES y 507000.0
    PERIODIC TIME 5000 y

```

```

/
#UNFILTER_NON_STATE_VARIABLES

VARIABLES
GAS_SATURATION
HYDRATE_SATURATION
GAS_DENSITY
POROSITY
ICE_SATURATION
TEMPERATURE
THERMODYNAMIC_STATE
LIQUID_SATURATION
GAS_PRESSURE
LIQUID_PRESSURE
PERMEABILITY
/
END

EXTERNAL_FILE flowConditions.txt

# initial condition
INITIAL_CONDITION all
FLOW_CONDITION initial
REGION all
END

BOUNDARY_CONDITION top
FLOW_CONDITION top
REGION top
END

BOUNDARY_CONDITION bottom
FLOW_CONDITION bottom
REGION bottom
END

END_SUBSURFACE

```

A.6. flowConditions.txt

Note this input deck is displayed in two-column format for space efficiency.

```

FLOW_CONDITION initial
TYPE
LIQUID_PRESSURE HYDROSTATIC
MOLE_FRACTION DIRICHLET
TEMPERATURE DIRICHLET
/
DATUM 0.d0 0.d0 1000.d0
GRADIENT
TEMPERATURE 0.d0 0.d0 -0.051790238209626256
/
LIQUID_PRESSURE 228791.4427924983
MOLE_FRACTION 1.d-10
TEMPERATURE 4.485964300000311
END

FLOW_CONDITION top
TYPE
GAS_PRESSURE DIRICHLET_SEEPAGE
RELATIVE_HUMIDITY DIRICHLET

```

```

TEMPERATURE DIRICHLET
/
GAS_PRESSURE LIST
TIME_UNITS year
DATA_UNITS Pa
INTERPOLATION LINEAR
0.d0 228791.4427924983d0
500.d0 216079.41141697354d0
1000.d0 204254.24913385956d0
1500.d0 195348.21356150892d0
2000.d0 186439.81370538234d0
2500.d0 178040.14871894888d0
3000.d0 170205.6557750324d0
3500.d0 164602.3892874983d0
4000.d0 159895.85107054017d0
4500.d0 154968.39927724394d0
5000.d0 149877.65678848544d0
5500.d0 147933.47740728944d0
6000.d0 145340.1116006376d0
6500.d0 143319.08471963252d0

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9000.d0 137756.4155298107d0
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 LIQUID_PRESSURE HYDROSTATIC
 MOLE_FRACTION DIRICHLET
 TEMPERATURE DIRICHLET

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DATA_UNITS Pa
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