

Applications of Machine Learning in Plasma Physics

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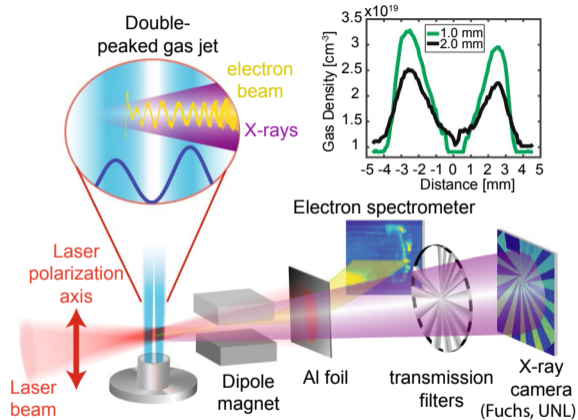
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Overview

- ▶ Primarily interested in basic plasma physics but with significant technological applications.
- ▶ Can't tolerate “black box” solutions. Need interpretability (mostly).
- ▶ Applications are all parameter estimation. Classification seems to have no use case in our work.
- ▶ Our training data is produced by simulations.
- ▶ Use ML to produce an approximation that can then be refined or validated by direct computations.
- ▶ Only makes sense if ML is computationally more efficient than the direct calculations.
- ▶ Big return from using ML to find reduced models.

Laser-Driven X-Ray Sources: Betatron Radiation

- ▶ Generate x-rays by causing energetic electrons to oscillate (violently) in a spatially varying electromagnetic field.
- ▶ Same mechanism as in a free-electron laser but incoherent, *i.e.*, without feedback.
- ▶ The electrons and fields are produced by a laser pulse ionizing a gas and subsequently driving large amplitude waves in the plasma.
- ▶ “Table top” version of a multi-kilometer facility.



Multitude of imaging applications ranging from medicine, biological and agricultural research to materials physics and national security.

Betatron Radiation

Calculating the Emitted Spectrum

- ▶ Comparing the radiation spectrum to theoretical predictions is a key element in validating our understanding of this system.
- ▶ Given an electron trajectory, compute the fields from the Liénard–Wiechert potentials:

$$\text{trajectories} \implies \mathbf{A}(\mathbf{r}, t) \implies \tilde{\mathbf{A}}(\mathbf{r}, \omega) \implies \frac{d^2 I}{d\omega d\Omega}.$$

Straightforward but expensive.

- ▶ Spectra from each particle are added incoherently.
- ▶ Neither electron trajectories nor the fields can be obtained from experiments.
- ▶ Obtaining the spectrum directly from the plasma physics simulations is not practical:
 - ▶ The radiation fields are much shorter wavelength than the mean plasma fields.
 - ▶ Using a grid to obtain the radiation from Maxwell's equations is not computationally reasonable.
 - ▶ To sample the particle trajectory with enough temporal resolution results in unmanageable amounts of data.

Betatron Radiation

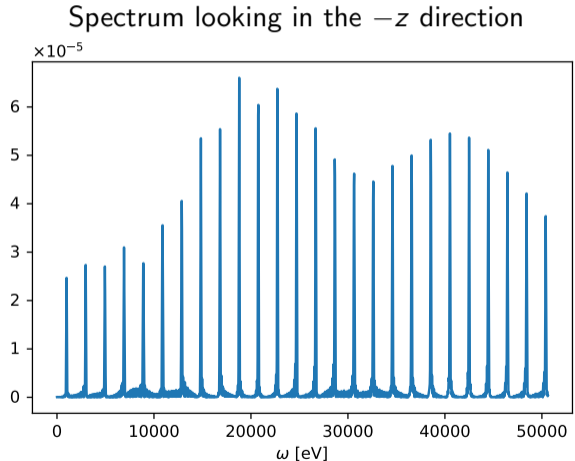
Calculating the Emitted Spectrum

- ▶ Since we don't have a good representation of the electron trajectories, we can live with some error in the spectrum.
- ▶ Sample the fields and particle phase space from the plasma physics simulations and use this as inputs to a MLP to estimate the single particle spectrum.
- ▶ We have a proof-of-concept implementation.

Calculating the Emitted Spectrum

Single Particle Spectrum

- ▶ Magnetic field
 - ▶ $\mathbf{B} = B_0 \hat{\mathbf{y}}$.
 - ▶ $B_0 = 400$ T.
 - ▶ Period: 100 microns.
 - ▶ Extent: 1 mm.
- ▶ Electron
 - ▶ Moving in the z direction.
 - ▶ Energy: 307 MeV ($\gamma = 600$).
- ▶ Spectrum is a series of spikes with little other structure.
- ▶ Spikes depend on γ and B_0 .
- ▶ For an electron beam with a range of energies, the spikes are blurred out in the complete spectrum.

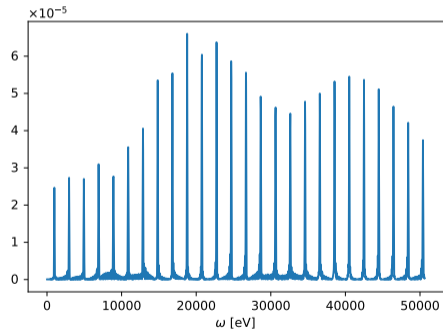


Calculating the Emitted Spectrum

ML Approach

- ▶ Extract ω_k and a_k for the first 25 spikes.
- ▶ Training data generated for $100 \leq \gamma \leq 1000$.
- ▶ Separate MLPs for estimating $\{\omega_k\}$ and $\{a_k\}$.
- ▶ Construct the spectrum from an electron beam by a weighted average of the single-particle spectra.
- ▶ Need to generalize field structure and initial electron direction.

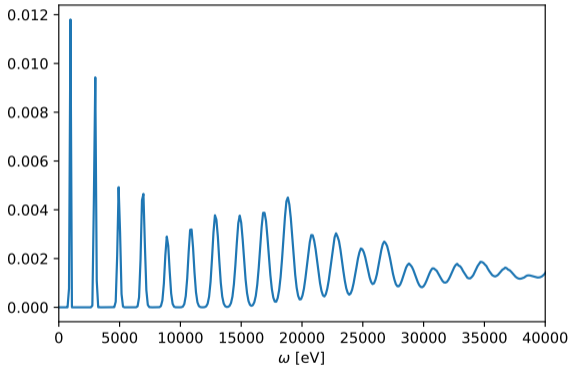
Spectrum looking in the $-z$ direction



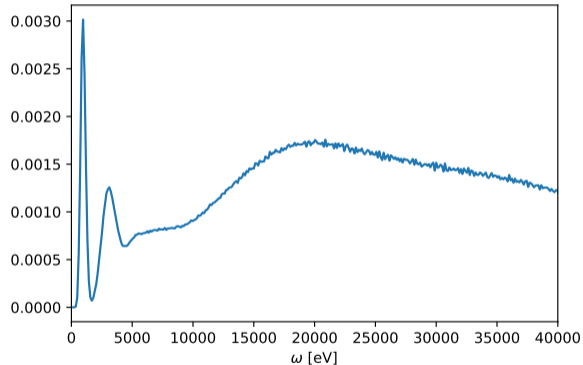
Calculating the Emitted Spectrum

ML Results

$\gamma = 600$, $\Delta\gamma/\gamma = 0.01$, 5000 particles



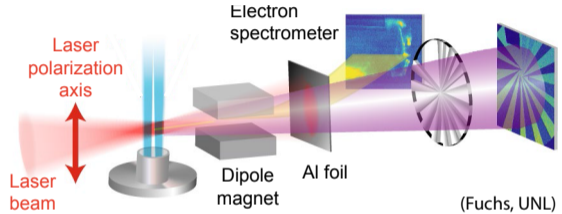
$\gamma = 600$, $\Delta\gamma/\gamma = 0.1$, 5000 particles



- ▶ Drop-off at large ω due to only keeping 25 harmonics.
- ▶ Several orders of magnitude faster than direct calculation from trajectories.
- ▶ More validation work needed.

Inverse Problem: Electron Beam Phase Space from the Energy Spectrum

- ▶ Knowledge of the electron beam phase space is critical to interpreting many experiments.
- ▶ Direct measurement of the beam phase space is not practical.
- ▶ Measure total charge.
- ▶ Measure energy with a magnetic spectrometer.
- ▶ Imaging of the beam in one spatial direction.
- ▶ All measurements are integrated over the beam head-to-tail.
- ▶ Determine macroscopic beam parameters (6-D correlation matrix, *i.e.*, beam emittance).
- ▶ Ill-posed inverse problem. Regularize using physics insight.
- ▶ Use ML parameter estimation to produce a starting point for a nonlinear optimization.



Vlasov–Maxwell System

The Vlasov–Maxwell equation is a phase space transport equation that well-describes many low collisionality plasmas.

- ▶ 1-particle phase-space distribution function $f(\mathbf{r}, \mathbf{p}, t)$.
- ▶ The evolution of f is given by

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} + q \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right) \cdot \frac{\partial f}{\partial \mathbf{p}} = 0,$$

where

$$\mathbf{v} = \mathbf{p}/\gamma m \quad \text{and} \quad \gamma = \sqrt{1 + p^2/m^2 c^2}.$$

- ▶ \mathbf{E} and \mathbf{B} are determined from the charge and current densities

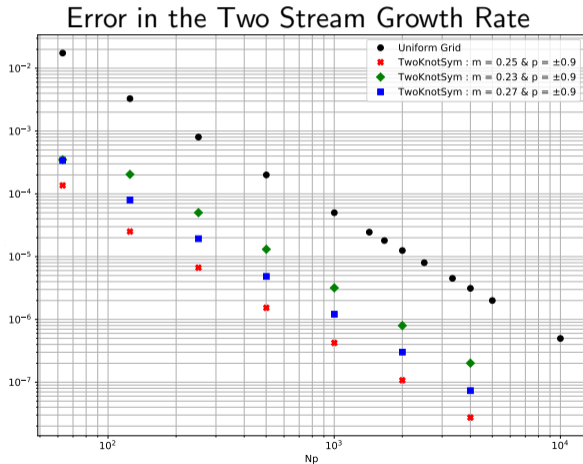
$$\rho(\mathbf{r}, t) = q \int d^3 \mathbf{p} f(\mathbf{r}, \mathbf{p}, t) \quad \text{and} \quad \mathbf{j}(\mathbf{r}, t) = q \int d^3 \mathbf{p} \mathbf{v} f(\mathbf{r}, \mathbf{p}, t).$$

- ▶ Computationally demanding in two or more spatial dimensions.
- ▶ Not all of phase space contributes equally to the fields; some regions are more important than others.

Vlasov–Poisson System

Improving Computational Performance with a Non-Uniform Momentum Grid.

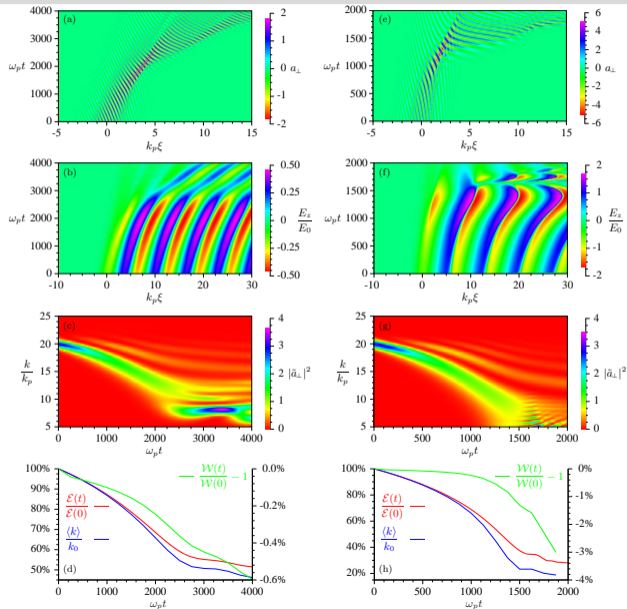
- ▶ Specialize to one spatial dimension, non-relativistic motion, and consider only electrostatic fields.
- ▶ Non-uniform grid in momentum space leads to significant computational savings.
- ▶ Use ML to predict grid parameters for a class of problems.
- ▶ How much training data is needed?
- ▶ How close can the predictions get to the optimal grid?
- ▶ This work is just beginning.



Reduced Models: Wakefield Excitation and Laser Pulse Evolution

- ▶ The laser pulse deposits energy into the plasma by exciting waves.
- ▶ The energy lost by the laser pulse manifests as a redshifting of the pulse.
- ▶ Not obvious; seems that the pulse amplitude should just decrease.
- ▶ An adiabatic invariant, wave action, forces the laser wave number to decrease as energy is lost:

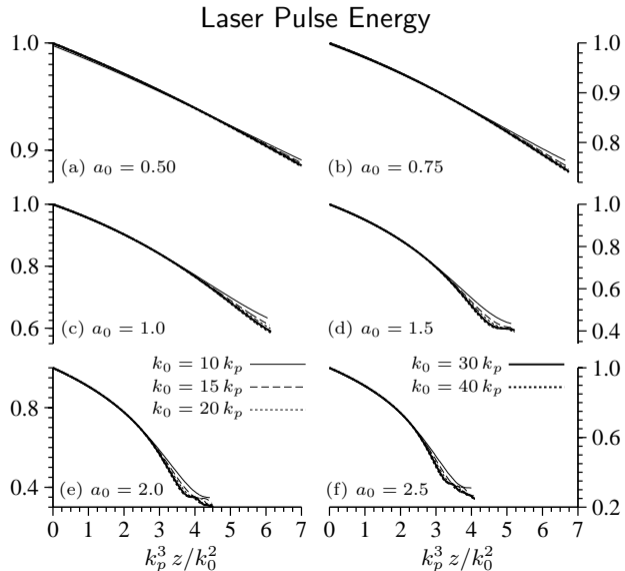
$$\frac{\langle k \rangle}{k_0} \mathcal{W} = \frac{k_0}{k_p} \mathcal{E}.$$



Reduced Models

- ▶ There is a set of universal scaling laws.
- ▶ Except for small k_0/k_p , the depletion is independent of k_0 .
- ▶ The spatial scale length for energy depletion is

$$\frac{k_p^3 z}{k_0^2}.$$



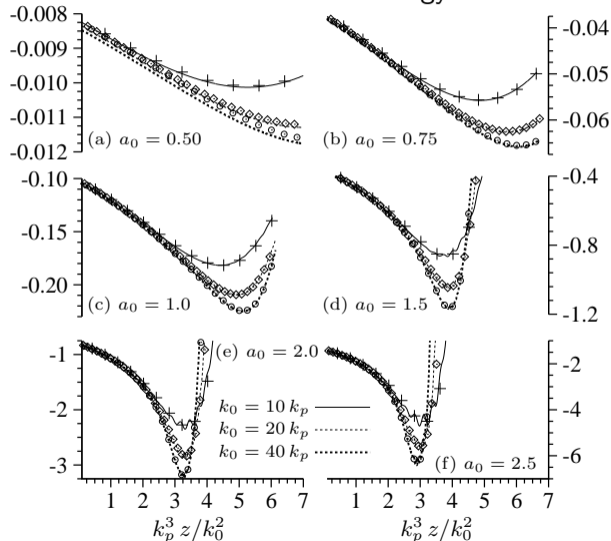
Reduced Models

- ▶ The rate of energy loss of the laser pulse is given by

$$\frac{d\mathcal{E}}{d\omega_p t} = - \left(\frac{k_p}{k_0} \frac{E_{\max}}{E_0} \right)^2.$$

- ▶ We have evidence for additional “hidden” scaling laws but these have not yet been fully explored or understood.
- ▶ Could this have been found by ML using latent-space methods?

Rate of Laser Pulse Energy Loss



Conclusions

- ▶ Numerous interesting applications of ML in basic plasma physics.
- ▶ Proof-of-concept implementation for computing x-ray spectra in laser-driven light sources.
- ▶ Finding scaling laws and reduced models is a high-payoff application. How?
- ▶ Possible applications in real-time guiding of experiments?

Acknowledgments

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