



The Schwarz alternating method as a means for concurrent multi-scale coupling of conventional and data-driven models

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Outline

- 1. Alternating Schwarz Method for Coupling of Full Order Models (FOMs) in Solid Mechanics
 - Motivation & Background
 - Quasistatics
 - Extension to Dynamics
 - Summary & Next Steps
- 2. Alternating Schwarz Method for FOM-ROM* and ROM-ROM Coupling
 - Motivation & Background
 - Demonstration
 - Ongoing & Future Work







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Motivation for Concurrent Multiscale Coupling

- Large scale structural failure frequently originates from small scale phenomena such as defects, microcracks, inhomogeneities and more, which grow quickly in unstable manner.
- Failure occurs due to *tightly coupled interaction* between small scale (stress concentrations, material instabilities, cracks, etc.) and large scale (vibration, impact, high loads and other perturbations).

Concurrent multiscale methods are essential for understanding and prediction of behavior of engineering systems when a small scale failure determines the performance of the entire system.



Roof failure of Boeing 737 aircraft due to fatigue cracks. From *imechanica.org*





Previous Multiscale Coupling Work

Comput Mech (2014) 54:803-820 DOI 10.1007/s00466-014-1034-0

ORIGINAL PAPER

A multiscale overlapped coupling formulation for large-deformation strain localization

WaiChing Sun · Alejandro Mota

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Abstract We generalize the multiscale overlapped domain framework to couple multiple rate-independent standard dissipative material models in the finite deformation regime across different length scales. We show that a fully coupled multiscale incremental boundary-value problem can be recast as the stationary point that optimizes the partitioned incremental work of a three-field energy functional. We also establish inf-sup tests to examine the numerical stability issues that arise from enforcing weak compatibility in the three-field formulation. We also devise a new block solver for the domain coupling problem and demonstrate the performance of the formulation with one-dimensional numerical examples. These simulations indicate that it is sufficient to introduce a localization limiter in a confined region of interest to regularize the partial differential equation if loss of ellipticity occurs.

strain localization may lead to the eventual failure of materials, this phenomenon is of significant importance to modern engineering applications.

The objective of this work is to introduce concurrent coupling between sub-scale and macro-scale simulations for inelastic materials that are prone to strain localization. Since it is not feasible to conduct sub-scale simulations on macroscopic problems, we use the domain coupling method such that computational resources can be efficiently allocated to regions of interest [14,23,24,30]. To the best of our knowledge, this is the first work focusing on utilizing the domain coupling method to model strain localization in inelastic materials undergoing large deformation.

Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious meshdependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

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Three-field multiscale coupling formulation with compatibility enforced weakly using *Lagrange multipliers*.

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Method works well, but is difficult to implement into existing codes.

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Requirements for Multiscale Coupling Method

- Coupling is *concurrent* (two-way).
- *Ease of implementation* into existing massively-parallel HPC codes.
- *Scalable, fast, robust* (we target *real* engineering problems, e.g., analyses involving failure of bolted components!).
- *"Plug-and-play" framework*: simplifies task of meshing complex geometries
 - Ability to couple regions with *different non-conformal meshes*, *different element types* and *different levels of refinement*.
 - > Ability to use *different solvers/time-integrators* in different regions.
- Coupling does not introduce *nonphysical artifacts.*
- *Theoretical* convergence properties/guarantees.







 Ω_1

Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.

Basic Schwarz Algorithm

• Solve PDE by any method on Ω_1 w/ initial guess for transmission BCs on Γ_1 .

Iterate until convergence:

Initialize:

- Solve PDE by any method on Ω_2 w/ transmission BCs on Γ_2 based on values just obtained for Ω_1 .
- Solve PDE by any method on Ω_1 w/ transmission BCs on Γ_1 based on values just obtained for Ω_2 .







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Overlapping Schwarz: convergent with all-Dirichlet transmission BCs¹ if $\Omega_1 \cap \Omega_2 \neq \emptyset$.

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Non-overlapping Schwarz: convergent with Robin-Robin² or alternating Neumann-Dirichlet³ transmission BCs.



H. Schwarz

(1843 - 1921)





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Common use of Schwarz: preconditioner for Krylov iterative methods to solve linear systems.



H. Schwarz



How we use the Schwarz Alternating Method





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Quasistatic Solid Mechanics Formulation



• Energy functional defining weak form of the governing PDEs

$$\Phi[\boldsymbol{\varphi}] \coloneqq \int_{\Omega} A(\boldsymbol{F}, \boldsymbol{Z}) dV - \int_{\Omega} \rho \boldsymbol{B} \cdot \boldsymbol{\varphi} dV$$

- ➤ A(F, Z): Helmholtz free-energy density
- → $F := \nabla \phi$: deformation gradient
- Z: collection of internal variables (for plastic materials)
- $\blacktriangleright \rho$: density, **B**: body force
- Euler-Lagrange equations, obtained by minimizing $\Phi[\boldsymbol{\varphi}]$: $\begin{cases} \text{Div } \boldsymbol{P} + \rho \boldsymbol{B} = \boldsymbol{0} \text{, in } \Omega \\ \boldsymbol{\varphi} = \boldsymbol{\chi}, & \text{on } \partial \Omega \end{cases}$
- Quasistatics solves sequence of problems in which loading (body force) *B* is incremented quasistatically w.r.t. pseudo time t_i:

For i = 1, ..., nSolve Div $P + \rho B(t_i) = 0$ with appropriate boundary conditions (BCs) Increment pseudo time t_i to obtain t_{i+1}

Spatial Coupling via Alternating Schwarz



Overlapping Domain Decomposition

$$\begin{cases} \operatorname{Div} \boldsymbol{P}_{1}^{(n+1)} + \rho \boldsymbol{B}(t_{i}) = \boldsymbol{0} , \text{ in } \Omega_{1} \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\chi}, & \text{ on } \partial \Omega_{1} \backslash \Gamma_{1} \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\varphi}_{2}^{(n)} & \text{ on } \Gamma_{2} \end{cases}$$
$$\begin{cases} \operatorname{Div} \boldsymbol{P}_{2}^{(n+1)} + \rho \boldsymbol{B}(t_{i}) = \boldsymbol{0} , \text{ in } \Omega_{2} \\ \boldsymbol{\varphi}_{2}^{(n+1)} = \boldsymbol{\chi}, & \text{ on } \partial \Omega_{2} \backslash \Gamma_{2} \\ \boldsymbol{\varphi}_{2}^{(n+1)} = \boldsymbol{\varphi}_{1}^{(n+1)} & \text{ on } \Gamma_{2} \end{cases}$$



Model PDE: $\begin{cases} \text{Div } \boldsymbol{P} + \rho \boldsymbol{B} = \boldsymbol{0} \text{, in } \Omega \\ \boldsymbol{\varphi} = \boldsymbol{\chi}, & \text{on } \partial \Omega \end{cases}$

 Dirichlet-Dirichlet transmission BCs [Schwarz, 1870; Lions, 1988]

Non-overlapping Domain Decomposition

$$\begin{cases} \operatorname{Div} \boldsymbol{P}_{1}^{(n+1)} + \rho \boldsymbol{B}(t_{i}) = \boldsymbol{0} \text{, in } \Omega_{1} \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\chi}, & \operatorname{on } \partial \Omega_{1} \backslash \Gamma \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\lambda}_{n+1} & \operatorname{on } \Gamma \\ \end{cases} \\\begin{cases} \operatorname{Div} \boldsymbol{P}_{2}^{(n+1)} + \rho \boldsymbol{B}(t_{i}) = \boldsymbol{0} \text{, in } \Omega_{2} \\ \boldsymbol{\varphi}_{2}^{(n+1)} = \boldsymbol{\chi}, & \operatorname{on } \partial \Omega_{2} \backslash \Gamma \\ \boldsymbol{P}_{2}^{(n+1)} \boldsymbol{n} = \boldsymbol{P}_{2}^{(n+1)} \boldsymbol{n}, & \operatorname{on } \Gamma \\ \end{cases} \\ \boldsymbol{\lambda}_{n+1} = \theta \boldsymbol{\varphi}_{2}^{(n)} + (1 - \theta) \boldsymbol{\lambda}_{n}, \operatorname{on } \Gamma, \text{ for } n \geq 1 \end{cases}$$

- Relevant for multi-material and multi-physics coupling
- Alternating Dirichlet-Neumann transmission BCs [Zanolli *et al.,* 1987]
- Robin-Robin transmission BCs also lead to convergence [Lions, 1990]
- θ ∈ [0,1]: relaxation parameter (can help convergence)

Spatial Coupling via Alternating Schwarz



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Dirichlet-Dirichlet transmission

BCs [Schwarz, 1870; Lions, 1988]

Part 1 of talk

Non-overlapping Domain Decomposition

Part 2 of talk

$$\begin{cases} \operatorname{Div} \boldsymbol{P}_{1}^{(n+1)} + \rho \boldsymbol{B}(t_{i}) = \boldsymbol{0} \text{, in } \Omega_{1} \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\chi}, & \text{on } \partial \Omega_{1} \setminus \Gamma \\ \boldsymbol{\varphi}_{1}^{(n+1)} = \boldsymbol{\lambda}_{n+1} & \text{on } \Gamma \\ \end{cases}$$

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Overlapping Schwarz Coupling in Quasistatics



Advantages:

- Conceptually very *simple*.
- Allows the coupling of regions with *different non-conforming meshes*, *different element types*, and *different levels of refinement*.
- Information is exchanged among two or more regions, making coupling concurrent.
- *Different solvers* can be used for the different regions.
- *Different material models* can be coupled if they are compatible in the overlap region.
- Simplifies the task of *meshing complex geometries* for the different scales.

Theoretical Foundation

Using the Schwarz alternating as a *discretization method* for PDEs is natural idea with a sound *theoretical foundation*.

- <u>S.L. Sobolev (1936)</u>: posed Schwarz method for *linear elasticity* in variational form and *proved method's convergence* by proposing a convergent sequence of energy functionals.
- <u>S.G. Mikhlin (1951)</u>: proved convergence of Schwarz method for general linear elliptic PDEs.
- <u>P.-L. Lions (1988)</u>: studied convergence of Schwarz for *nonlinear monotone elliptic problems* using max principle.
- <u>A. Mota, I. Tezaur, C. Alleman (2017)</u>: proved *convergence* of the alternating Schwarz method for *finite deformation quasi-static nonlinear PDEs* (with energy functional Φ[φ]) with a *geometric convergence rate.*

$$\boldsymbol{\Phi}[\boldsymbol{\varphi}] = \int_{B} A(\boldsymbol{F}, \boldsymbol{Z}) \, dV - \int_{B} \boldsymbol{B} \cdot \boldsymbol{\varphi} \, dV$$
$$\nabla \cdot \boldsymbol{P} + \boldsymbol{B} = \boldsymbol{0}$$



S.L. Sobolev (1908 – 1989)



S.G. Mikhlin (1908 – 1990)



P.- L. Lions (1956-)



A. Mota, I. Tezaur, C. Alleman

Convergence Proof*



A. Moto, I. Tezane, C. Alleman Schwarz Alternating Method in Solid Mechanics	A. Mota, I. Tazaur, C. Allaman Schwarz Alternating Mathod in Solid Machanics	A. Hota, I. Trezave, C. Alfernau Scinware Alternating Method in Solid Mechanics	A. Mota, J. Teçane, C. Alloran Schwarz Alternating Method in Solid Mechanics
2 Formulation of the Schwarz Alternating Method		1. $\pi_{1}^{(1)} = \mathbf{X}_{1}^{(1)} \approx \Omega_{1}, \pi_{1}^{(1)} + \mathbf{\chi}(\mathbf{X}_{1}^{(1)}) \otimes \partial_{\theta} \Omega_{1},$ is initialized for Ω_{1} $\geq \pi_{B}^{(2)} = \mathbf{X}_{1}^{(2)} \approx \Omega_{1}, \pi_{1}^{(2)} + \mathbf{\chi}(\mathbf{X}_{1}^{(2)}) \otimes \partial_{\theta} \Omega_{1},$ is initialized for Ω_{1} .	Remark that [50] $\hat{S}_n = \hat{\omega}^{(n-1)} + \hat{V}_i \text{ for } \hat{\omega}^{(n-1)} \in \hat{S}_{n-1} \Rightarrow \hat{\omega}^{(n-1)} \in \hat{S}_n.$ (40)
we start by defining the standard inner detormation variational formulation to establish notation before presenting the formulation of the coupling method.	(m, r ₂ t)) r, m	$ \begin{array}{c} & \left(\sum_{j=1}^{N-1} \left(\sum_{j=1}^{N-1} \left(\sum_{j=1}^{N-1} \left(K_{Ajj}^{(1)} + K_{Ajj}^{(1)} H_{11} - K_{Ajj}^{(1)} H_{11} - K_{Ajj}^{(1)} H_{2j} \right) \right) \\ & \left(\sum_{j=1}^{N-1} \left(\sum_{j=1}^{N-1} \left(K_{Ajj}^{(2)} + K_{Ajj}^{(2)} H_{2j} - K_{Ajj}^{(2)} H_{2j} \right) \right) \\ & \left(\sum_{j=1}^{N-1} \left(K_{Ajj}^{(2)} + K_{Ajj}^{(2)} H_{2j} - K_{Ajj}^{(2)} H_{2j} \right) \right) \\ & \left(\sum_{j=1}^{N-1} \left(K_{Ajj}^{(2)} + K_{Ajj}^{(2)} + K_{Ajj}^{(2)} H_{2j} - K_{Ajj}^{(2)} + K_{Ajj}^{(2)} H_{2j} \right) \right) \\ & \left(\sum_{j=1}^{N-1} \left(K_{Ajj}^{(2)} + K_{Ajj}^{(2)} $	Theorem 1. Assume that the energy functional $\Phi(\varphi)$ ratisfies properties 1–5 above. Consider the Schwarz alternative webod of Section 2. defined for $(96.4)31$ and its conjunction form (39). Then
2.1 Variational Formulation on a Single Domain		$\begin{array}{cccccccccccccccccccccccccccccccccccc$	(a) $\Phi[\tilde{\varphi}^{(1)}] \ge \Phi[\tilde{\varphi}^{(1)}] \ge \cdots \ge \Phi[\tilde{\varphi}^{(n-1)}] \ge \Phi[\tilde{\varphi}^{(n)}] \ge \cdots \ge \Phi[\varphi]$, where φ is the minimizer of $\Phi[\varphi]$
Consider a basy as in a open set $U \subset V^*$ undergoing a motion described by the interpret $x = \varphi(X)$ if $M \to V^*$, $X \in \Omega$. Assume that the boundary of the body is $\partial \Omega = \partial_{\mu}\Omega \cup \partial_{T}\Omega$ with unit normal N , where $\partial_{\mu}\Omega$ is a displacement boundary, $\partial_{T}\Omega$ is a machino boundary, and $\partial_{\alpha}\Omega \cap \partial_{T}\Omega = 0$. The prescribed boundary	Figure 1: Two subdomains \mathbb{R}_2 and \mathbb{R}_2 and the corresponding boundaries Γ_1 and Γ_2 used by the Schwatz attenting method.	$2 \mod \left[\left(\left (\Delta x_R^{(1)} x _R^{(1)})^2 + \left(\left (\Delta x_R^{(1)} x_R^{(1)}) \right \right)^2 \right]^{1/2} \le \epsilon_{maxim}$ is tight tolerance.	(b) the sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in S.
displacements or Dirichlet boundary conditions are $\chi : \partial_{\varphi}\Omega \rightarrow \mathbb{R}^{3}$. The prescribed boundary tractions or Neumann boundary conditions are $\chi : \partial_{\varphi}\Omega \rightarrow \mathbb{R}^{3}$. Let $F := \text{Grad} \varphi$ be the deformation gradient. Let also $DI : \Omega \rightarrow \mathbb{R}^{3}$ be both force with B the most density in the afference configuration. Bethermore	that is $i = 1$ and $i = 2$ if a isorial and $i = 2$ and $i = 1$ if a isorem. Introduce the following definitions for	(35.34.4) Although use do not receively formal consumers records for the remaining variants of the	(c) the Schwarz minimum values Φ[φ ^(m)] converge monotonically to the minimum value Φ[φ] is S starting from any initial guess φ ⁽⁰⁾ .
introduce the energy functional	and hard and a set of the set of	Schwarz method, we offer some numerical results illustrating their convergence in Section 4. Consider the energy functional $\Phi[\phi]$ defined in (1). We will denote by $\langle , , \rangle$ the usual L^2 inner product	(d) if Φ [†] (φ) is Lipschitz continuous in a neighborhood of φ, then the sequence {φ ⁽ⁿ⁾ } converges geometrically to the minimizer φ ¹ .
$\Phi[\varphi] := \int_{\Omega} A(F, Z) dV - \int_{\Omega} RB \cdot \varphi dV - \int_{\partial \varphi \cdot \Omega} T \cdot \varphi dS,$ (1)	・ Coosine か := の/ wo ・ Dirichtet boundary: 奥田 := 奥田 臣.	over Ω , that is, $(\psi_1, \psi_2) := \int \psi_1 \cdot \psi_2 dV,$ (35)	Proof. See Appendix A.
in which $A(F, Z)$ is the Helmholtz free-energy density and Z is a collection of internal variables. The weak form of the problem is obtained by minimizing the energy functional $\Phi(\phi)$ over the Sobolev space $W_2^+(\Omega)$ that is compared of Ω for the problem is obtained by minimizing the energy functional form internal terms of the problem space.	Neumann boundary: @ ID := @ ID : ID	f_{f1} for $\psi_1, \psi_2 \in W_2^1(\Omega)$, with corresponding neural $ \cdot $. The proof of the convergence of the Schwarz alternating	Finally, while most of works cited above present their analysis for the specific case of two subdomains, extension to multiple subdomains is in general straightforward. The case of multiple subdomains is considered
that is comprised of an indication that are square-integrated and nave square-integrated integrated into activatives, beame	 Schwarz boundary: 1 / Se (86/1 K). 	memor requires that the functional w(p) satisfy the following properties over the space 5 denied in (2):	specifically in Liens [35], Badea [4], and Li-Shan and Evans [34].
$S := \{\varphi \in w_2(n) : \varphi = \chi \text{ on } b\varphi n\}$ (2) and	Note that were made deministers we guarantee that $(q; \phi) \in (q; \phi) = (; (q; \phi) + (; ; and (q; \phi) + (; ; ;))$ Now define the spaces	 Ψ(φ) is Fréchet differentiable, with Ψ'(φ) denoting its Fréchet derivative. 	4 Numerical Examples
$V := \{\xi \in W_2(\Omega) : \xi = 0 \text{ on } \partial_{\varphi}\Omega\}$ (3) where $\xi \in V$ is a tent function. The extension is minimized if and only if $M(z) \in M(z)$, of free 4	$S_i := \{ 2 W_2^1(m) : ' = \chi \text{ on } \otimes m_i ' = P_{2i_j / \Gamma_i} (m_j) \} \text{ on } \Gamma_i$, (7)	 Φ(φ) is strictly convex. 	In this section, we present numerical examples of the behavior of the Schwarz alternating method for two
where $\xi \in V$ is a rest function. The potential energy is minimized in and only if $\forall (\varphi) \subseteq \forall [\varphi] \forall e^-e_{\xi}$ for an $\xi \in V$ and $e \in \mathbb{R}$. It is straightforward to show that the minimum of $\Phi[\varphi]$ is the mapping $\varphi \in S$ that satisfies	and $V_i := \{-2 \ W_2^{(i)}(\infty) : i = 0 \text{ on } (0, \infty) \{\Gamma_i \},$ (8)	 Φ[φ] is tower semi-continuous. 	different implementations. First, we briefly describe the two implementations, one in MATLAB and the other in the open-source ALBANY finite element code [32]. Next, we discuss the error measures used throughout the open-source ALBANY finite element code [32].
$D\Phi[\varphi](\xi) = \int_{\Omega} \mathbf{P} : \text{Grad} \xi dV - \int_{\Omega} R\mathbf{B} \cdot \xi dV - \int_{\partial P \cdot \Omega} \mathbf{T} \cdot \xi dS = 0,$ (4)	where the symbol $P_{iij+l-1}$ () denotes the projection from the subdomain iij) onto the Schwarz boundary Γ_j . This projection operator plays a central role in the Schwarz alternating method. Its form and implementation	5. $\Phi'[\varphi]$ is uniformly continuous on K_B , where $F_{} = (x \in S, A(x) \in B, B \in B, K \in a)$ (26)	the numerican examples. Then, we commode with four examples that demonstrate different features of the Schwarz alternating method and our implementations. The first example, a one-dimensional singular bar, is most to demonstrate the hebraic of the four Schwarz variants of Schwarz alternation 2.4. The accord example, a subsid
where $P = \partial A/\partial F$ denotes the first Piola-Kirchhoff stress. The Euler-Lagrange equation corresponding to the variational statement (4) is	are discussed in subsequent sections. For the moment it is sufficient to assume that the operator is able to project a field 'from one subdomain to the Schwarz boundary of the other subdomain.	$m_R \rightarrow q \varphi \in \mathcal{O} : \pi(\varphi) \subset m, R \in \mathbb{P}, R \in \infty$ }. (38) It can be shown that the energy functional $\vartheta(\varphi)$ defined in (1) is strictly convex in S (property 3) provided	body of square body, aims to study the effect of the size of the overage region on the convergence of the method. The objective of the third example, a notched cylinder, is to analyze the numerical error in the results
	The curvatz anathening method surves a sequence of problems on log and log. The solution 1100 for the	that the Helmholtz free energy density $A(F, Z)$ is a must convex function of F [26]. Properties 1.2.4 and	and the demonstrate the ability of the method to couple different element topologies. The last example, a laser

Theorem 1. Assume that the energy functional $\Phi[\varphi]$ satisfies properties 1–5 above. Consider the Schwarz alternating method of Section 2 defined by (9)–(13) and its equivalent form (39). Then

- (a) $\Phi[\tilde{\varphi}^{(0)}] \ge \Phi[\tilde{\varphi}^{(1)}] \ge \cdots \ge \Phi[\tilde{\varphi}^{(n-1)}] \ge \Phi[\tilde{\varphi}^{(n)}] \ge \cdots \ge \Phi[\varphi]$, where φ is the minimizer of $\Phi[\varphi]$ over S.
- (b) The sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in S.
- (c) The Schwarz minimum values $\Phi[\tilde{\varphi}^{(n)}]$ converge monotonically to the minimum value $\Phi[\varphi]$ in S starting from any initial guess $\tilde{\varphi}^{(0)}$.

Remark 1 By the coercivity of $\Phi[\phi]$, it follows from the Lax-Milgram theorem that a unique minimizer to this functional maps S mints in the minimization of $\Delta[\phi]$ is well as well.	Avain using (57) and also (50) in (60) leads to	$\lim_{n\to\infty} \hat{\varphi}^{(n)} - \hat{\varphi}^{(n+1)} ^2 = 0, (69)$	$(\Phi'[\phi^{(n)}], \phi - \phi^{(n)}) \le (\Phi'[\phi^{(n)}], \phi - \phi^{(n)}) + \alpha_R \phi - \phi^{(n)} \le \Phi \phi - \Phi \phi^{(n)} $ (79)
Remark 2. By the Stampacchia theorem, the minimization of $\Phi(\varphi)$ in S is equivalent to finding $\varphi \in S$ such	$(\Phi^{\prime}[\hat{\phi}^{(n)}] - \Phi^{\prime}[\hat{\phi}^{(n-1)}], \hat{\zeta}_{2}) = (\Phi^{\prime}[\hat{\phi}^{(n)}], \xi) \le \Phi^{\prime}[\hat{\phi}^{(n)}] - \Phi^{\prime}[\hat{\phi}^{(n-1)}] \cdot \hat{\zeta}_{2} ,$ (61)	from which we can conclude that $\bar{\varphi}^{(n)} - \bar{\varphi}^{(n+1)} \rightarrow 0$ as $n \rightarrow \infty$. We must now show that $\bar{\varphi}^{(n)}$ converges to ω , the minimizer of $\Phi(\omega)$ on S . By (53) with $\psi_1 = \omega$ and	since $\alpha_B \ge 0$. Now, by the Cauchy-Schwarz inequality followed by the application of the Lipshitz continuity of $\Phi'[q]$ (66) we can write
mar $(\Phi'[\phi], \xi - \phi) \ge 0$ (51)	and substituting (56) into (61) we finally obtain that	$\psi_2 = \tilde{\varphi}^{(n)}$, we have	$(\Phi'[\tilde{\varphi}^{(n)}], \varphi - \tilde{\varphi}^{(n)}) \le \Phi'[\tilde{\varphi}^{(n)} \cdot \varphi - \tilde{\varphi}^{(n)} \le K \varphi - \tilde{\varphi}^{(n)} ^2.$ (80)
for all $\xi \in S$.	$(\Psi'[\tilde{\varphi}^{(n)}], \xi) \le C_0 \Psi'[\tilde{\varphi}^{(n)}] - \Psi'[\tilde{\varphi}^{(n-1)}] \cdot \xi ,$ (62)	$ \phi - \dot{\phi}^{(n)} ^2 \le \frac{1}{\alpha_R} \left\{ \Phi[\phi] - \Phi[\dot{\phi}^{(n)}] - \left(\Phi'[\dot{\phi}^{(n)}], \phi - \dot{\phi}^{(n)}\right) \right\}.$ (70)	Hence, from (79),
Remark 3 Recall that the strict convexity property of $\Phi[\varphi]$ can be written as	$\forall \xi \in S.$	Since φ is the minimum of $\Phi[\varphi]$, by (a) we have that $\Phi[\varphi] \le \Phi[\hat{\varphi}^{(n)}]$. It follows that	$\Phi[\phi^{(n)}] - \Phi[\phi] \le K \phi^{(n)} - \phi ^2.$ (81)
$\Phi[\psi_2] - \Phi[\psi_1] - (\Phi'[\psi_1], \psi_2 - \psi_1) \ge 0,$ (52)	Remark 8 For part (d) of Theorem 1, recall the definition of geometric convergence:	$\Phi[\phi] - \Phi[\bar{\phi}^{(n)}] - \left(\Phi'[\bar{\phi}^{(n)}], \phi - \bar{\phi}^{(n)}\right) \le - \left(\Phi'[\bar{\phi}^{(n)}], \phi - \bar{\phi}^{(n)}\right) = \left(\Phi'[\bar{\phi}^{(n)}], \bar{\phi}^{(n)} - \phi\right).$ (71)	Moreover, by (53) since $\Psi [\varphi] = 0$, $\Phi [\pi (\alpha)] = \Phi [\pi (\alpha)] = \Phi [\pi (\alpha)] = 0$ (22)
$\forall \psi_1, \psi_2 \in S$. From (36), remark that if $\Phi[\varphi]$ is strictly convex over $S \forall R \in \mathbb{R}$ such that $R < \infty$, we can find	$E_{n+1} \le CE_n$, (63)	Substituting (71) into (70) we have	$ \Psi \varphi^{\mu}, -\Psi \varphi \ge a_R \varphi^{\mu}, -\varphi $. (62)
an $\alpha_R > 0$ such that $\forall \psi_1, \psi_2 \in \Lambda_R$ we have	$\forall n \in \{0, 1, 2,\}$ for some $C > 0$, where	$ _{i\sigma} = \hat{a}^{(n)}_{i\sigma} ^{2} \le \frac{1}{2} \left(g^{i}_{i\sigma} _{i\sigma}^{(n)} - \hat{a}^{(n)}_{i\sigma} - \alpha^{i}_{i\sigma} \right)$ (75)	Using (k) and (k) we obtain (ad-(ad), ad-)) (ad-(ad-), ad-)) can a(a) (ad-(ad-)) (ad-
$\Psi[\psi_2] - \Psi[\psi_1] - (\Psi[\psi_1], \psi_2 - \psi_1) \ge a_R[\psi_2 - \psi_1]^*.$ (53)	$E_n := \phi^{(n+1)} - \phi^{(n)} .$ (64)	$ar - r - a_R (r - r - r), \qquad (ar - r)$	$\langle \Psi \varphi^{(n)} - \Psi \varphi \rangle = \langle \Psi \varphi^{(n+n)} - \Psi \varphi \rangle \le K \varphi^{(n)} - \varphi ^n - \alpha_R \varphi^{(n+n)} - \varphi ^n.$ (83)
Remark 4 By property 5, the uniform continuity of $\Phi'[\varphi]$, there exists a modulus of continuity $\omega > 0$, with $\omega \in U_n \rightarrow U_n$, such that	Bemark 9 Recall from the definition of continuity that if $\Phi^{(i)}(a)$ is Linsbitz continuous at $a^{(a)}$ near a , then	Now by (c_) (Remark 7),	Combining (83) and (78) leads to
$ \Psi'(\psi_1) - \Psi'(\psi_2) \le \omega(\psi_1 - \psi_2),$ (54)	there exists a constant $K \ge 0$ such that	$(\Psi'[\hat{\varphi}^{(n')}], \hat{\varphi}^{(n)} - \varphi) \le C_0 \Psi'[\hat{\varphi}^{(n')}] - \Psi'[\hat{\varphi}^{(n'-1)}] \cdot \hat{\varphi}^{(n)} - \varphi .$ (73)	$\frac{a_R}{C_1} \phi^{(n)} - \phi ^2 \le \left(\Phi[\phi^{(n)}] - \Phi[\phi]\right) - \left(\Phi[\phi^{(n+1)}] - \Phi[\phi]\right) \le K \phi^{(n)} - \phi ^2 - a_R \phi^{(n+1)} - \phi ^2.$
$\forall \psi_1, \psi_2 \in K_R$. By definition, $\omega(e) \rightarrow 0$ as $e \rightarrow 0$.	$\frac{\ \Phi'[\phi^{(\alpha)}] - \Phi'[\phi]\ }{\ \phi\ _{1}} \le K.$ (65)	Substituting (73) into (72) leads to	0r (84)
Remark 5 It was shown in [35] that in the case $\Omega_1 \cap \Omega_2 \neq \emptyset$, $\forall \varphi \in S$, there exist $\zeta_1 \in S_1$ and $\zeta_2 \in S_2$ such that	$\ \varphi^{(n)} - \varphi\ $	$ \hat{\varphi}^{(n)} - \varphi \le \frac{\zeta_0}{\alpha_R} \Psi'[\hat{\varphi}^{(n)}] - \Psi'[\hat{\varphi}^{(n-1)}] .$ (74)	$ \phi^{(n+1)} - \phi \le B \phi^{(n)} - \phi $ (85)
$\varphi = \zeta_1 + \zeta_2$, (55)	Considering that $\Phi'[\varphi] = 0$ since φ is the minimizer of $\Phi(\varphi)$, (65) is equivalent to	Applying the uniform continuity assumption (54), we obtain	with $B := \sqrt{\frac{K-1}{K-1}}$ (96)
and $\max(f_1 , f_2) \le C_0 x $, (56)	$ \Psi'[\phi^{(n)}] \le K \phi^{(n)} - \phi .$ (66)	$ \hat{\varphi}^{(n)} - \varphi \le \frac{C_0}{\omega} \left(\hat{\varphi}^{(n)} - \hat{\varphi}^{(n-1)} \right).$ (75)	$D := \sqrt{\alpha_R} - \overline{C_1}$, (60) and $B \in \mathbb{R}$ are show $C \to -V'$. Furthermore, $(\lambda(0))$ are set of the top
for some $C_0 > 0$ independent of φ .	Proof of Theorem 1	Og v 7 Be (00, 0.50) - 500-501 - 0 are	the minimizer φ of $\Phi(\varphi)$ by (b) and (c), it follows that $B \in (0, 1)$. Define $C := 1 - B \in (0, 1)$, then (85)
Remark 6 Note that (39) can be written as	Proof of (a). Let $\hat{\varphi}^{(1)} = \arg \min_{\varphi \in \hat{S}_1} \Phi(\varphi)$. By (40), $\hat{\varphi}^{(1)} \in \hat{S}_2$. Let $\hat{\varphi}^*$ be the minimizer of $\Phi(\varphi)$ over \hat{S}_2 and $\Phi(\varphi) = \Phi(\varphi)$. The set of the	By (0) , $\ \psi^{-1} - \psi^{-1}\ \to 0$ as $n \to \infty$. From any we obtain the result, namely that $\psi^{-1} \to \psi$ is $n \to \infty$.	can be recast as $ \phi^{(n+1)} - \phi^{(n)} \le C \phi^{(n)} - \phi^{(n-1)} $ (87)
$(\Phi'[\phi^{(n)}], \xi^{(i)}) = 0$, for $\phi^{(n)} \in \hat{S}_n, \forall \xi^{(i)} \in S_i$, (57)	that $\Phi[\tilde{\varphi}^{(1)}] < \Phi[\tilde{\varphi}^{(2)}]$, where $\tilde{\varphi}^{(2)} = \arg\min_{\varphi \in \tilde{\varphi}^0} \varphi $. It follows by induction that	Proof of (c). This follows immediately from (a) and (b).	whereupon the claim is proven.
for $i \in \{1, 2\}$ and $n \in \{0, 1, 2,\}$ (recall from (6) the relation between i and n). This is due to the uniqueness	$\Phi(\bar{\omega}^{(n)}) \le \Phi(\bar{\omega}^{(n-1)})$ (67)	Proof of (d). By (b), for large enough n, there exists some $C_1 > 0$ independent of n such that	D. Annalasta Calastica for Directo Pilantis Channelso Dec
of the solution to each minimization problem over \hat{S}_n and the definition of $\hat{\varphi}^{(n)}$ as the minimizer of $\Phi[\varphi]$ over \hat{S}_n	for $n \in \{1, 2, 3,\}$. Now let ω be the minimizer of $\Phi[\omega]$ over S . Since the problem is well-need ω is	$ \phi^{(n)} - \phi ^2 \le C_1 \phi^{(n+1)} - \phi^{(n)} ^2.$ (76)	B Analytic Solution for Linear-Elastic Singular Bar
Remark 7 Let $\hat{\varphi}^{(n)} \in \hat{S}_n$, and let $\xi \in S$. By Remark 5, there exist $\zeta_1 \in S_1$ and $\zeta_2 \in S_2$ such that	unique. Hence $\Phi(\varphi) \le \Phi[\varphi^{(n)}]$ for all $n \in \{1, 2, 3,\}$.	Let us choose C_1 such that $C_1 > \alpha_R/K$, where K is the Lipshitz continuity constant in (66). Combining (68) with (76) leads to	As reference, herein we provide the solution of the singular bar of Section 4.3 for linear elasticity. The equilibrium equation is
$(\Phi^{i}[\hat{\varphi}^{(n)}], \xi) = (\Phi^{i}[\hat{\varphi}^{(n)}], \xi_{1} + \xi_{2}).$ (58)		$\frac{1}{\alpha_N} \left(\Phi[\tilde{\varphi}^{(n)}] - \Phi[\tilde{\varphi}^{(n+1)}] \right) \ge \tilde{\varphi}^{(n+1)} - \tilde{\varphi}^{(n)} ^2 \ge \frac{1}{C_1} \tilde{\varphi}^{(n)} - \varphi ^2. (77)$	$P = \sigma(X)A(X) = \text{const.}, \sigma(X) = E\epsilon(X), e(X) := a'(X), A(X) = A_0\left(\frac{X}{L}\right)^{\frac{1}{2}},$ (88)
34	35	36	37

Implementation in Albany LCM Code

The proposed *quasistatic alternating Schwarz method* is implemented within the *Albany LCM* open-source parallel, C++, multi-physics, finite element code.

- **Component-based** design for rapid development.
- Contains a wide variety of *constitutive models*.
- Extensive use of libraries from the open-source *Trilinos* project.
 - Use of the *Phalanx* package to decompose complex problem into simpler problems with managed dependencies.
 - Use of the *Sacado* package for *automatic differentiation*.
 - Use of *Teko* package for block preconditioning for Modified and Monolithic Schwarz variants.
- Parallel implementation of Schwarz alternating method uses the Data Transfer Kit (DTK).
- All software available on *GitHub*.

Albany

https://github.com/SNL Computation/LCM



https://github.com/trilinos/trilinos



https://github.com/ORNL-CEES/DataTransferKit

Quasistatic Example #1: Cuboid Problem



- Coupling of *two cuboids* with square base (above).
- *Neohookean*-type material model.



Schwarz Iteration

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Cuboid Problem: Convergence with Overlap & Refinement



Cuboid Problem: Schwarz Error





Subdomain	u_3 relative error	σ_{33} relative error	
$\Omega_1 \ \Omega_2$	1.24×10^{-14} 7.30×10^{-15}	2.31×10^{-13} 3.06×10^{-13}	L





- *Notched cylinder* that is stretched along its axial direction.
- Domain decomposed into *two subdomains*.
- Neohookean-type material model.

Notched Cylinder: TET - HEX Coupling



- The Schwarz alternating method is capable of coupling *different mesh topologies*.
- The notched region, where stress concentrations are expected, is *finely* meshed with *tetrahedral* elements.
- The top and bottom regions, presumably of less interest, are meshed with *coarser hexahedral* elements.



Notched Cylinder: TET - HEX Coupling





Notched Cylinder: TET - HEX Coupling



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Notched Cylinder: Coupling Different Materials

The Schwarz method is capable of coupling regions with *different material models*.

- Notched cylinder subjected to tensile load with an *elastic* and *J2 elasto-plastic* regions.
- *Coarse* region is *elastic* and *fine* region is *elasto-plastic*.
- The *overlap region* in the first mesh is nearer the notch, where plastic behavior is expected.

Coupled regions

Coarse, elastic region

Fine, elasto-plastic region



Notched Cylinder: Coupling Different Materials

Need to be careful to do domain decomposition so that material models are *consistent* in overlap region.

- When the *overlap* region is *far from the notch*, no plastic deformation exists in it: the coarse and fine regions predict the *same behavior*.
- When the *overlap* region is *near the notch*, plastic deformation spills onto it and the two models predict different behavior, affecting convergence *adversely*.



Overlap far from notch.

Overlap near notch.

Quasistatic Example #3: Laser Weld





Laser weld specimen

- Problem of *practical scale*.
- Isotropic elasticity and J₂ plasticity with linear isotropic hardening.
- *Identical parameters* for weld and base materials for proof of concept, to become independent models.



Single domain discretization

Laser Weld: Strong Scalability of Parallel **1** Schwarz with DTK



• Near-ideal linear speedup (64-1024 cores).





Data Transfer Kit (DTK)

Sandia National

Quasistatic Example #4: Tensile Bar



The alternating Schwarz method can be used as part of a *homogenization* (upscaling) process to bridge gap b/w *microscopic* and *macroscopic* regions

Goal: study strain localization in microstructure. Microstructure embedded in ASTM tensile geometry (right). Fix microstructure, investigate macroensemble of uniaxial loads. scale Fit flow curves with a *macroscale* J_2 plasticity model (below). 350 micro-300 Cauchy Stress 1 structure true stress (MPa) 005 007 150 macro-LO CP ensembles scale 12 fit Work by C. Alleman, J. Foulk, 100 0.005 0.010 0.015 0.020 0.025 0.030 0.035 0.040 equivalent plastic strain(mm/mm) D. Littlewood, G. Bergel

Outline

- 1. Alternating Schwarz Method for Coupling of Full Order Models (FOMs) in Solid Mechanics
 - Motivation & Background
 - Quasistatics
 - Extension to Dynamics
 - Summary & Next Steps
- 2. Alternating Schwarz Method for FOM-ROM* and ROM-ROM Coupling
 - Motivation & Background
 - Demonstration
 - Ongoing & Future Work







Dynamic Solid Mechanics Formulation

• Kinetic energy:
$$T(\dot{\boldsymbol{\varphi}}) \coloneqq \frac{1}{2} \int_{\Omega} \rho \dot{\boldsymbol{\varphi}} \cdot \dot{\boldsymbol{\varphi}} \, dV$$

- Potential Energy: $V(\boldsymbol{\varphi}) \coloneqq \int_{\Omega} A(\boldsymbol{F}, \boldsymbol{Z}) dV \int_{\Omega} \rho \boldsymbol{B} \cdot \boldsymbol{\varphi} dV$
- Lagrangian: $L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) \coloneqq T(\dot{\boldsymbol{\varphi}}) V(\boldsymbol{\varphi})$
- Action functional: $S[\boldsymbol{\varphi}] \coloneqq \int_{I} L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) dt$
- Euler-Lagrange equations: $\begin{cases}
 \text{Div } \boldsymbol{P} + \rho \boldsymbol{B} = \rho \boldsymbol{\ddot{\varphi}}, & \text{in } \Omega \times I \\
 \boldsymbol{\varphi}(\boldsymbol{X}, t_0) = \boldsymbol{x}_0, & \text{in } \Omega \\
 \boldsymbol{\dot{\varphi}}(\boldsymbol{X}, t_0) = \boldsymbol{v}_0, & \text{in } \Omega \\
 \boldsymbol{\varphi}(\boldsymbol{X}, t) = \boldsymbol{v}_0, & \text{on } \partial \Omega \times I
 \end{cases}$
- Semi-discrete problem following FEM discretization in space:

$$M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}}$$



Schwarz Alternating Method for Dynamics

 In the literature the Schwarz method is applied to dynamics by using *spacetime discretizations*.



Overlapping non-matching meshes and time steps in dynamics.


Schwarz Alternating Method for Dynamics

 In the literature the Schwarz method is applied to dynamics by using *spacetime discretizations*.

Pro ☺: Can use *non-matching* meshes and time-steps (see right figure).

Con ②: *Unfeasible* given the design of our current codes and size of simulations.



Overlapping non-matching meshes and time steps in dynamics.





<u>Step 0</u>: Initialize i = 0 (controller time index).

at which subdomains are synchronized

Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Schwarz Alternating Method for Dynamic Multiscale Coupling



Controller time stepper: defines global ΔTs at which subdomains are synchronized

Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize i = 0 (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.



at which subdomains are synchronized

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Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.



at which subdomains are synchronized

Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

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Step 3: Check for convergence at time T_{i+1} .



at which subdomains are synchronized

Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize i = 0 (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

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Step 3: Check for convergence at time T_{i+1} .

If unconverged, return to Step 1.

Schwarz Alternating Method for Dynamic ^{Sandia} Multiscale Coupling Controller time stepper: defines global ΔTs at which subdomains are synchronized



 $\begin{array}{c|c} T_1 & T_2 \\ \hline \\ Integrate using \Delta t_1 \\ \hline \\ \Omega_2 \text{ to } \Gamma_1 & & \\ \end{array} \\ \end{array}$

<u>Step 0</u>: Initialize i = 0 (controller time index).

<u>Step 1</u>: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

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<u>Step 3</u>: Check for convergence at time T_{i+1} .

- ➢ If unconverged, return to Step 1.
- ▶ If converged, set i = i + 1 and return to Step 1.

Schwarz Alternating Method for Dynamic ^{Sandia} Multiscale Coupling Controller time stepper: defines global ΔTs at which subdomains are synchronized





<u>Step 0</u>: Initialize i = 0 (controller time index).

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<u>Step 3</u>: Check for convergence at time T_{i+1} .

- ➢ If unconverged, return to Step 1.
- ▶ If converged, set i = i + 1 and return to Step 1.

Can use *different integrators* with *different time steps* within each domain!

Schwarz Alternating Method for Dynamic Multiscale Coupling: Theory

- Like for quasistatics, dynamic alternating Schwarz method converges provided each single-domain problem is *well-posed* and *overlap region* is *non-empty*, under some *conditions* on Δt .
- Well-posedness for the dynamic problem requires that action functional $S[\boldsymbol{\varphi}] \coloneqq \int_{I} \int_{\Omega} L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) dV dt$ be strictly convex or strictly concave, where $L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) \coloneqq T(\dot{\boldsymbol{\varphi}}) + V(\boldsymbol{\varphi})$ is the Lagrangian.
 - > This is studied by looking at its second variation $\delta^2 S[\boldsymbol{\varphi}_h]$
- We can show assuming a *Newmark* time-integration scheme that for the *fully-discrete* problem:

$$\delta^2 S[\boldsymbol{\varphi}_h] = \boldsymbol{x}^T \left[\frac{\gamma^2}{(\beta \Delta t)^2} \boldsymbol{M} - \boldsymbol{K} \right] \boldsymbol{x}$$

 $\succ \delta^2 S[\boldsymbol{\varphi}_h]$ can always be made positive by choosing a *sufficiently small* Δt

> Numerical experiments reveal that Δt requirements for *stability/accuracy* typically lead to automatic satisfaction of this bound.

Implementation in Albany LCM & Sierra/SM

Numerical results shown here for *dynamic* Schwarz are from two codes: *Albany LCM* and *Sierra/Solid Mechanics (Sierra/SM*)

- *Albany LCM:* Trilinos-based open-source* parallel, C++, multi-physics, FE code.
- Sierra Mechanics Framework: Sandia Lagrangian 3D code for FEA of solids & structures.



Sierra/SM: quasistatics, implicit/explicit dynamics + loose coupling via *Arpeggio*



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> Schwarz alternating method was **"implemented"** in *Sierra/SM* using *Arpeggio* loose coupling framework

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Sierra/SM: quasistatics, implicit/explicit dynamics + loose coupling via *Arpeggio*



Schwarz alternating method was **"implemented"** in *Sierra/SM* using *Arpeggio* loose coupling framework

We did not have to write any code in Sierra/SM to implement Schwarz!

Dynamic Example #1: Elastic Wave Propagation

- Linear elastic *clamped beam* with Gaussian initial condition for the *z*-displacement.
- Simple problem with analytical exact solution but very stringent test for discretization methods.
- Test Schwarz with **2** subdomains: $\Omega_0 = (0,0.001) \times (0,0.001) \times (0,0.75), \Omega_1 = (0,0.001) \times (0,0.001) \times (0.25,1).$



Elastic Wave: Diff. Integrators, Same Δts

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Table 1: Averaged (over times + domains) relative errors in **z-displacement** (blue) and **z-velocity** (green) for several different Schwarz couplings, 50% overlap volume fraction

	Implicit-Implicit		Explicit(CM)-Implicit		Explicit(LM)-Implicit	
Conformal HEX - HEX	2.79e-3	7.32e-3	3.53e-3	8.70e-3	4.72e-3	1.19e-2
Nonconformal HEX - HEX	2.90e-3	7.10e-3	2.82e-3	7.29e-3	2.84e-3	7.33e-3
TET - HEX	2.79e-3	7.58e-3	3.52e-3	8.92e-3	4.72e-3	1.19e-2

Elastic Wave: Diff. Integrators, Diff. Δt s





Figures above: Plots of displacement, velocity and acceleration for the elastic wave propagation problem using different time integrators (implicit and explicit) and different time steps (1e-2s and 2e-7s) for each subdomain, superimposed over the analytic single domain solution.

The analytic solution is *indistinguishable* from Schwarz solutions (hidden behind the solutions for Ω_0 (red) and Ω_1 (green))!

Dynamic Example #2: Tension Specimen

- Uniaxial aluminum cylindrical tensile specimen with *inelastic J₂ material model*.
- Domain decomposition into *two* subdomains (right): Ω_0 = ends, Ω_1 = gauge.
- Nonconformal HEX + composite TET10 coupling via Schwarz.
- Implicit Newmark time-integration with adaptive time-stepping algorithm employed in both subdomains.
- Slight *imperfection* introduced at center of gauge to force *necking* upon pulling in vertical direction.







Tension Specimen: Expected Result



Tension Specimen: Disp. & EQPS





*EQPS = Equivalent Plastic Strain

Tension Specimen: Domain Decomposition







55

Tension Specimen: EQPS





Tension Specimen: Cauchy Stress





Dynamic Example #3: Bolted Joint Problem

Problem of *practical scale*.



• Schwarz solution compared to single-domain solution on composite TET10 mesh.



- $\Omega_1 = \text{bolts}$ (Composite TET10), $\Omega_2 = \text{parts}$ (HEX).
- Inelastic J₂ material model in both subdomains.
 - Ω_1 : steel
 - Ω_2 : steel component, aluminum (bottom) plate



- BC: x-disp = 0.02 at T = 1.0e-3 on top of parts.
- Run until T = 5.0e-4 w/ dt = 1e-5 + implicit Newmark with analytic mass matrix for composite tet 10s.



Bolted Joint Problem: Displacement lime: 0.000000 ΔY AY Schwarz Single $\boldsymbol{\Omega}$

Bolted Joint Problem: EQPS







Single Ω

Schwarz

Bolted Joint Problem: Convergence Rate





Figure above: Convergence behavior of the dynamic Schwarz algorithm for the bolted joint problem



	CPU times (64 procs*)	Avg # Schwarz iters	Max # Schwarz iters	
Single Domain	3h 34m	_	_	
Schwarz	2h 42m	3.22	4	
Single Domain (finer)	17h 00m	_	Sie	and
Schwarz (finer mesh of bolts)	29h 29m	3.28	4	



	CPU times (64 procs*)	Avg # Schwarz iters	Max # Schwarz iters	
Single Domain	3h 34m	—	—	
Schwarz	2h 42m	3.22	4	
Single Domain (finer)	17h 00m	_	Sie	and
Schwarz (finer mesh of bolts)	29h 29m	3.28	4	

 Despite its iterative nature, Schwarz can actually be *faster* than single domain run for discretizations having comparable # of elements in the bolts.



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Single Domain	3h 34m	—	—	
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Schwarz (finer mesh of bolts)	29h 29m	3.28	4	and the

- Despite its iterative nature, Schwarz can actually be *faster* than single domain run for discretizations having comparable # of elements in the bolts.
 - Even if the method is more computationally expensive for some resolutions, it may be preferred for its ability to *rapidly change* and *evaluate* a *variety* of *engineering designs* (our typical use case).



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- Despite its iterative nature, Schwarz can actually be *faster* than single domain run for discretizations having comparable # of elements in the bolts.
 - Even if the method is more computationally expensive for some resolutions, it may be preferred for its ability to *rapidly change* and *evaluate* a *variety* of *engineering designs* (our typical use case).
- Dynamic Schwarz converges in between 2-4 Schwarz iterations per time-step despite the overlap region being very small for this problem.

overlap region

Outline

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Summary



The Schwarz alternating method has been developed/implemented for concurrent multiscale quasistatic & dynamic modeling in Sandia's *Albany LCM* and *Sierra/SM* codes.

- Coupling is *concurrent* (two-way).
- ③ *Ease of implementation* into existing massively-parallel HPC codes.
- Scalable, fast, robust (we target real engineering problems, e.g., analyses involving failure of bolted components!).
- ③ *"Plug-and-play" framework*: simplifies task of meshing complex geometries!
 - Oblight to couple regions with *different non-conformal meshes*, *different element types* and *different levels of refinement*.
 - ③ Ability to use *different solvers/time-integrators* in different regions.
- © Coupling does not introduce *nonphysical artifacts.*
- Theoretical convergence properties/guarantees.

Ongoing Work and Next Steps



- Continuing to *apply* the Schwarz alternating method to problems of interest to *production* using *Sierra/SM*.
- Advancing the Schwarz alternating method to enable *coupling* of *structural elements* to *continuum elements*.
- Developing a *Schwarz-like algorithm* for simulating *contact* to address wellknown challenges (contact constraint enforcement, multiple scales).
 - Done by introducing a *combination* of *Dirichlet* and *Neumann* boundary conditions into different subdomains in a *non-overlapping* alternating fashion.
 - See the following *reference* for more info:

J. Hoy, **I. Tezaur**, A. Mota. "The Schwarz alternating method for multiscale contact mechanics". in *Computer Science Research Institute Summer Proceedings 2021*, J.D. Smith and E. Galvan, eds., Technical Report SAND2021-0653R, Sandia National Labs, 360-378, 2021. (<u>https://www.sandia.gov/ccr/csri-summer-programs/2021-proceedings/</u>)

> Journal article with D. Koliesnikova in preparation.

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Motivation

The past decades have seen tremendous investment in **simulation frameworks** for **coupled multi-scale** and **multi-physics** problems.

- Frameworks rely on **established mathematical theories** to couple physics components.
- Most existing coupling frameworks are based on traditional discretization methods.





Complex System Model

- PDEs, ODEs
- Nonlocal integral
- Classical DFT
- Atomistic, ...

Traditional Methods

- Mesh-based (FE, FV, FD)
- Meshless (SPH, MLS)
- Implicit, explicit
- Eulerian, Lagrangian, ...

Coupled Numerical Model

 N_4

- Monolithic (Lagrange multipliers)
- Partitioned (loose) coupling
- Iterative (Schwarz, optimization)

ical Model Traditional + Data-Driven Methods

- PINNs
- Neural ODEs
- Projection-based ROMs, ...

aboratories

(EAM)

Land Ice (MALI) Ocean (MPAS-

Land (ALM)

Sea Ice

Motivation

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N_3







Coupled Numerical Model Traditional + Data-Driven Methods PINNs ٠

- Neural ODEs
- Projection-based ROMs, ...

- **Complex System Model**
- PDEs. ODEs
- Nonlocal integral
- Classical DFT
- Atomistic, ...

Traditional Methods

- Mesh-based (FE, FV, FD)
- Meshless (SPH, MLS)
- Implicit, explicit
- Eulerian, Lagrangian, ... •

- Monolithic (Lagrange multipliers)
- Partitioned (loose) coupling
- Iterative (Schwarz, optimization)
- There is currently a big push to integrate data-driven methods into modeling & simulation toolchains.

Unfortunately, existing algorithmic and software infrastructures are **ill-equipped** to handle plug-and-play integration of **non-traditional**, data-driven models!
Flexible Heterogeneous Numerical Methods (fHNM) Project



Principal research objective:



This talk.

 $\int \mathcal{M}^2 dt$

- **Discover mathematical principles** guiding the assembly of **standard** and **data-driven** numerical models in stable, accurate and physically consistent ways.
- Principal research challenges: we lack mathematical and algorithmic understanding of how to
- "Mix-and-match" standard and data-driven models from three-classes

Class A: projection-based reduced order models (ROMs)

- Class B: machine-learned models, i.e., Physics-Informed Neural Networks (PINNs)
- Class C: flow map approximation models, i.e., dynamic model decomposition (DMD) models
- Ensure **well-posedness & physical consistency** of resulting **heterogeneous models**.
- Solve such heterogeneous models efficiently.

Three coupling methods:

- This talk.
- Alternating Schwarz-based coupling
- Optimization-based coupling
- Coupling via generalized mortar methods



Projection-Based Model Order Reduction via the POD/Galerkin Method







ROM = Projection-Based Reduced Order Model. HROM = Hyper-Reduced ROM

Schwarz Extensions to FOM-ROM and ROM-ROM Couplings



Enforcement of Dirichlet boundary conditions (DBCs) in ROM at indices $i_{\rm Dir}$

• Method I in [Gunzburger *et al.* 2007] is employed

 $\boldsymbol{u}(t) \approx \overline{\boldsymbol{u}} + \boldsymbol{\Phi} \widehat{\boldsymbol{u}}(t), \quad \boldsymbol{v}(t) \approx \overline{\boldsymbol{v}} + \boldsymbol{\Phi} \widehat{\boldsymbol{v}}(t), \quad \boldsymbol{a}(t) \approx \overline{\boldsymbol{a}} + \boldsymbol{\Phi} \widehat{\boldsymbol{a}}(t)$

> POD modes made to satisfy homogeneous DBCs: $\Phi(i_{\text{Dir}},:) = 0$

 \succ BCs imposed by modifying \overline{u} , \overline{v} , \overline{a} : $\overline{u}(i_{\text{Dir}}) \leftarrow \chi_u$, $\overline{v}(i_{\text{Dir}}) \leftarrow \chi_v$, $\overline{a}(i_{\text{Dir}}) \leftarrow \chi_a$

Choice of domain decomposition

• Error-based indicators that help decide in what region of the domain a ROM can be viable should drive domain decomposition (future work) [Bergmann *et al.* 2018]

Snapshot collection and reduced basis construction

- Ideally, generate snapshots/reduced bases separately in each subdomain Ω_i [Hoang *et al.* 2021, Smetana *et al.*, 2022]
- POD results presented herein use snapshots obtained via FOM-FOM coupling on $\Omega = \bigcup_i \ \Omega_i$

For nonlinear problems, hyper-reduction needs to preserve Hamiltonian structure

- We employ Energy-Conserving Sampling & Weighting Method (ECSW) [Farhat et al. 2015]
- Investigating structure-preserving coupling is near-term future work.

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Numerical Example: 1D Dynamic Wave Propagation Problem

- **1D beam** geometry $\Omega = (0,1)$, clamped at both ends, with Gaussian initial condition discretized using FEM + Newmark- β
- Simple problem but very stringent test for discretization/ coupling methods.
- Two constitutive models considered:
 - Linear elastic (problem has exact analytical solution)
 - Nonlinear hyperelastic Henky
 This talk.



ROMs results are *reproductive* and *predictive*, and are based on the *POD/Galerkin* method, with POD calculated from FOM-FOM coupled model.

> 50 POD modes capture ~100% snapshot energy for linear variant of this problem.

- ➢ 536 POD modes capture ~100% snapshot energy for Henky variant of this problem.
- Hyper-reduced ROMs (HROMs) perform *hyper-reduction* using ECSW [Farhat *et al.,* 2015]
 Ensures that *Lagrangian structure* of problem is preserved in HROM.
- **Couplings tested:** overlapping, non-overlapping, FOM-FOM, FOM-ROM, ROM-ROM, FOM-HROM, HROM-HROM, implicit-explicit, implicit-implicit, explicit-explicit. **This talk.**



Numerical Example: 1D Dynamic Wave **Propagation Problem**

- **Two variants** of problem, with different initial conditions (ICs):
 - Symmetric Gaussian IC (top right)
 - Rounded Square IC (bottom right)
- Non-overlapping domain decomposition (DD) of $\Omega = \Omega_1 \cup$ $Ω_2$, where $Ω_1$ = [0, 0.6] and $Ω_2$ = [0.6, 1.0]
 - DD based on heuristics: during time-interval considered $(0 \le t \le 1 \times 10^3)$, sharper gradient forms in Ω_2 , suggesting FOM or larger ROM is needed there.
- Reproductive problem:
 - Displacement snapshots collected using FOM-FOM nonoverlapping coupling with Symmetric Gaussian IC
 - ➢ FOM-ROM, FOM-HROM, ROM-ROM and HROM-HROM run with Symmetric Gaussian IC
- Predictive problem:
 - Snapshots same as for reproductive problem.
 - FOM-ROM, FOM-HROM, ROM-ROM and HROM-HROM run with Rounded Square IC



0.1





Numerical Example: Reproductive Results in Sandia Automatical Example: Reproductive Results in Sandia



Model	M_1/M_2	$N_{e,1}/N_{e,2}$	CPU time (s)	$rac{\mathcal{E}_{ ext{MSE}}(ilde{m{u}}_1)}{\mathcal{E}_{ ext{MSE}}(ilde{m{u}}_2)}$	$rac{\mathcal{E}_{ ext{MSE}}(ilde{m{v}}_1)/}{\mathcal{E}_{ ext{MSE}}(ilde{m{v}}_2)}$	$egin{array}{c} \mathcal{E}_{ ext{MSE}}(ilde{m{a}}_1)/\ \mathcal{E}_{ ext{MSE}}(ilde{m{a}}_2) \end{array}$	N_S	
FOM	-/-	-/-	1.871×10^{3}	-/-	-/-	-/-	-	
ROM	60/-	-/-	1.398×10^{3}	$1.659 \times 10^{-2}/-$	1.037×10^{-1}	$4.681 \times 10^{-1}/-$	-	
HROM	60/-	155/-	$5.878 imes 10^2$	$1.730 \times 10^{-2}/-$	$1.063 \times 10^{-1}/-$	$4.741 \times 10^{-1}/-$	—	
ROM	200/-	-/-	1.448×10^{3}	$2.287 \times 10^{-4}/-$	$4.038 \times 10^{-3}/-$	$4.542 \times 10^{-2}/-$	-	
HROM	200/-	428/-	9.229×10^{2}	$8.396 \times 10^{-4}/-$	$8.947 \times 10^{-3}/-$	$7.462 \times 10^{-2}/-$	-	Green
FOM-FOM	-/-	-/-	2.345×10^{3}	_	_	_	24,630	
FOM-ROM	-/80	-/-	2.341×10^{3}	2.171×10^{-6}	3.884×10^{-5}	$2.982 \times 10^{-4}/$	25,227	shading
				1.253×10^{-5}	2.401×10^{-4}	2.805×10^{-3}		highlights
FOM-HROM	-/80	-/130	2.085×10^{3}	2.022×10^{-4}	$1.723e \times 10^{-3}/$	$7.421 \times 10^{-3}/$	$29,\!678$	inginging
				5.734×10^{-4}	5.776×10^{-3}	3.791×10^{-2}		most
FOM-BOM	-/200	-/-	2.449×10^{3}	4.754×10^{-12}	1.835×10^{-10}	5.550×10^{-9}	24,630	compotitivo
1 0111 100111	,200	/	2.110 × 10	7.357×10^{-11}	4.027×10^{-9}	1.401×10^{-7}	21,000	
FOM-HROM	-/200	-/252	2.352×10^{3}	1.421×10^{-5}	1.724×10^{-4}	9.567×10^{-4}	27.156	coupled
	,	,		4.563×10^{-4}	2.243×10^{-3}	1.364×10^{-2}		coupled
ROM-ROM	200/80	-/-	2.778×10^{3}	4.861×10^{-5}	1.219×10^{-3}	1.586×10^{-2}	27.810	models
		,		3.093×10^{-3}	4.177×10^{-4}	3.936×10^{-3}	.,	
HROM-HROM	200/80	315/130	1.769×10^{3}	3.410×10^{-3}	4.110×10^{-2}	2.485×10^{-1}	29,860	
				6.662×10^{-4}	6.432×10^{-3}	4.307×10^{-2}		
ROM-ROM	300/80	-/-	2.646×10^{3}	2.580×10^{-6}	6.226×10^{-5}	9.470×10^{-4}	25.059	
	,	,		1.292×10^{-3}	2.483×10^{-4}	2.906×10^{-3}		
HROM-HROM	300/80	405/130	1.938×10^{3}	6.960×10^{-3}	6.328×10^{-2}	3.137×10^{-1}	29,896	
				7.230×10^{-4}	7.403×10^{-3}	$ 4.960 \times 10^{-2}$		

- All coupled models evaluated converged on average in **<3** Schwarz iterations per time-step
- FOM-ROM coupling has same total # Schwarz iters ($N_{\rm S}$) as FOM-FOM coupling
- Other couplings require more Schwarz iters than FOM-FOM coupling to converge
 - More Schwarz iters required when coupling less accurate models
 - Larger 300/80 mode ROM-ROM takes less time than smaller 200/80 mode ROM-ROM
- FOM-HROM & HROM-HROM couplings outperform FOM-FOM coupling in CPU time by 12.5-32.6%
- All couplings involving ROMs/HROMs are **at least as accurate** as single-domain ROMs/HROMs

Numerical Example: Reproductive Results 🔂





- Single-domain ROM and • **HROM** are most efficient
- Couplings involving **ROMs and HROMs** enable one to achieve smaller errors
- Benefits of hyperreduction are limited on 1D problem

Numerical Example: Reproductive Results 🔂



Numerical Example: Predictive Results





- Predictive accuracy/robustness can be improved by coupling ROM or HROM to FOM
 FOM-ROM coupling is remarkably accurate, achieving displacement error O(1 × 10⁻⁸)

 FOM-HROM and ROM-ROM couplings are more accurate than single-domain ROMs
 HROM-HROM on par with single-domain HROM in terms of accuracy
- Wall-clock times of coupled models can be improved
 - FOM-HROM, ROM-ROM and HROM-HROM models are slower than FOM-FOM model as more Schwarz iterations required to achieve convergence
 - > Hyper-reduction samples ~60% of total mesh points
 - Greater gains from hyper-reduction anticipated for 2D/3D problems

Numerical Example: Predictive Results





- Predictive **single-domain ROM** solution exhibits **spurious oscillations** in velocity and acceleration
- Predictive FOM-HROM solution is smooth and oscillation-free
 > Highlights coupling method's ability to improve ROM predictive accuracy

Numerical Example: Predictive Results





 Single-domain FOM solution 	– Solution in Ω_1	– Solution in Ω_2
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Ongoing & Future Work

Ongoing Work

- Extension to multi-material problems (method works; results not shown herein).
- Development of additive Schwarz variant, which has more potential for parallelization and speed-ups (preliminary FOM-FOM results show promise).
- Extension of method to multi-D through newly-developed Python code, pressio-demoapps¹
 library and/or 3D solid mechanics Julia code².

> Multi-D implementation requires transfer operators for transmission BCs

- Journal publication on results presented here + 2D/3D results.
- Alternating Schwarz-based coupling of Physics-Informed Neural Networks (PINNs)

Next Steps

- Development of smart DD approaches based on error indicators, to determine optimal placement of ROM and FOM (including on-the-fly ROM-FOM switching).
- Analysis of method's convergence properties for non-overlapping and ROM/HROM coupling cases.
- Extension of coupling approach to POD modes built from snapshots on independentlysimulated subdomains.
- Application to other problems, including multi-physics problems.
- Structure-preservation within ROMs and couplings involving ROMs







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Careers at Sandia National Labs



<u>Students</u>: please consider Sandia and other national labs as a potential employer for summer internships and when you graduate!



- Sandia is a multidisciplinary national lab and Federally Funded Research & Development Center (FFRDC).
- Contractor for U.S. DOE's National Nuclear Security Administration (NNSA).
- **Two main sites**: Albuquerque, NM and Livermore, CA

Careers at Sandia National Labs



<u>Students</u>: please consider Sandia and other national labs as a potential employer for summer internships and when you graduate!

- Sandia is a *great* place to work!
 - Lots of *interesting* problems that require *fundamental research* in applied math/computational science and impact *mission-critical applications*.
 - Great work/life balance.
- **Opportunities** at/with Sandia:
 - Interns (summer, year-round)
 - Post docs
 - Several prestigious post doctoral fellowships (von Neumann, Truman, Hruby, Data Science)
 - ➤ Staff

Please see: <u>www.sandia.gov/careers</u> for info about current opportunities.





Careers at Sandia National Labs



Intern- Scientific Machine Learning Technical Graduate Summer

Job ID 686827 Location Albuquerque, NM Department Scientific Machine Learning Job Family Student Posted Date 11/30/2022

Intern - Computer Science Research Institute (CSRI) - R&D Graduate Summer

Job ID 685380 Location Multiple Department Scalable Algorithms Job Family Student



Posted Date 12/14/2022

Intern - Computer Science Research Institute (CSRI) - R&D Undergraduate Summer

Job ID 685381

Location Multiple

Department Scalable Algorithms

Job Family Student

Posted Date 12/01/2022

•

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Left: current postings from <u>www.sandia.gov/careers</u> for Summer 2023. *Summer interns are hired by early March.*

I am hiring interns to help with ROM-ROM/ROM-FOM coupling work!

If interested, email me (<u>ikalash@sandia.gov</u>) and apply to the CSRI Internship Posting (left).



Start of Backup Slides

Appendix. Four Variants* of Schwarz





Full Schwarz



$1: \boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)} \text{ in } \Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_1, \boldsymbol{x}_\beta^{(1)} \leftarrow \boldsymbol{X}_\beta^{(1)} \text{ on } \Gamma_1$	\triangleright initialize for Ω_1
2: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{X}_B^{(2)}$ in $\Omega_2, \boldsymbol{x}_b^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_2, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_2 3: repeat	⊳ initialize for Ω ₂ ⊳ Newton-Schwarz loop
4: $\boldsymbol{x}_{eta}^{(1)} \leftarrow \boldsymbol{P}_{12} \boldsymbol{x}_{B}^{(2)} + \boldsymbol{Q}_{12} \boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12} \boldsymbol{x}_{eta}^{(2)}$	\triangleright project from Ω_2 to Γ_1
5: $\Delta \boldsymbol{x}_B^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_B^{(1)}; \boldsymbol{x}_b^{(1)}; \boldsymbol{x}_\beta^{(1)}) \setminus \boldsymbol{R}_A^{(1)}(\boldsymbol{x}_B^{(1)}; \boldsymbol{x}_b^{(1)}; \boldsymbol{x}_\beta^{(1)})$	⊳ linear system
6: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{x}_B^{(1)} + \Delta \boldsymbol{x}_B^{(1)}$	
7: $\pmb{x}_{eta}^{(2)} \leftarrow \bar{\pmb{P}_{21}} \pmb{x}_{B}^{(1)} + \bar{\pmb{Q}}_{21} \pmb{x}_{b}^{(1)} + \pmb{G}_{21} \pmb{x}_{eta}^{(1)}$	\triangleright project from Ω_1 to Γ_2
$ 8: \qquad \bigtriangleup \bm{x}_B^{(2)} \leftarrow -\bm{K}_{AB}^{(2)}(\bm{x}_B^{(2)}; \bm{x}_b^{(2)}; \bm{x}_\beta^{(2)}) \backslash \bm{R}_A^{(2)}(\bm{x}_B^{(2)}; \bm{x}_b^{(2)}; \bm{x}_\beta^{(2)}) $	⊳ linear system
9: $oldsymbol{x}_B^{(2)} \leftarrow oldsymbol{x}_B^{(2)} + riangle oldsymbol{x}_B^{(2)}$	
$10: \text{ until } \left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \right)^2 \right]^{1/2} \le \epsilon_{\text{machine}}$	⊳ tight tolerance

Modified Schwarz



Monolithic Schwarz

Inexact Schwarz

*A. Mota, I. Tezaur, C. Alleman. "The Schwarz Alternating Method in Solid Mechanics", CMAME 319 (2017), 19-51. 92

Appendix. Four Variants* of Schwarz



Modified and *Monolithic* Schwarz variants create and solve a *block system*.



Full Schwarz



1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_1, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_1 \triangleright initialize for Ω_1 2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} \triangleright initialize for Ω_2 3: repeat ▷ Newton-Schwarz loop 4: $\boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{P}_{12}\boldsymbol{x}_{\beta}^{(2)} + \boldsymbol{Q}_{12}\boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12}\boldsymbol{x}_{\beta}^{(2)}$ \triangleright project from Ω_2 to Γ_1 5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$ ▷ linear system $\begin{array}{l} 6: \quad \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{x}_{B}^{(1)} + \Delta \boldsymbol{x}_{B}^{(1)} \\ 7: \quad \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{P}_{21} \boldsymbol{x}_{B}^{(1)} + \boldsymbol{Q}_{21} \boldsymbol{x}_{b}^{(1)} + \boldsymbol{G}_{21} \boldsymbol{x}_{\beta}^{(1)} \end{array}$ \triangleright project from Ω_1 to Γ_2 8: $\Delta \boldsymbol{x}_{B}^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{B}^{(2)}) \setminus \boldsymbol{R}_{A}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{b}^{(2)})$ ⊳ linear system 9: $\boldsymbol{x}_{P}^{(2)} \leftarrow \boldsymbol{x}_{P}^{(2)} + \Delta \boldsymbol{x}_{P}^{(2)}$ $10: \text{ until } \left[\left(|| \triangle \boldsymbol{x}_B^{(1)} || / || \boldsymbol{x}_B^{(1)} || \right)^2 + \left(|| \triangle \boldsymbol{x}_B^{(2)} || / || \boldsymbol{x}_B^{(2)} || \right)^2 \right]^{1/2} \le \epsilon_{\text{machine}}$ ▷ tight tolerance

Modified Schwarz



Monolithic Schwarz

Inexact Schwarz

*A. Mota, I. Tezaur, C. Alleman. "The Schwarz Alternating Method in Solid Mechanics", CMAME 319 (2017), 19-51.

Appendix. Four Variants* of Schwarz



Most performant method: monotonic convergence, theoretical convergence guarantee.



Full Schwarz

1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_1, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_1	\triangleright initialize for Ω_1
2: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{X}_B^{(2)}$ in $\Omega_2, \boldsymbol{x}_b^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_2, \boldsymbol{x}_{\boldsymbol{\beta}}^{(2)} \leftarrow \boldsymbol{X}_{\boldsymbol{\beta}}^{(2)}$ on Γ_2	\triangleright initialize for Ω_2
3: repeat	▷ Schwarz loop
4: $\boldsymbol{y}^{(1)} \leftarrow \boldsymbol{x}^{(1)}_B$	⊳ for convergence check
5: $\boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{P}_{12} \boldsymbol{x}_{\beta}^{(2)} + \boldsymbol{Q}_{12} \boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12} \boldsymbol{x}_{\beta}^{(2)}$	\triangleright project from Ω_2 to Γ_1
6: repeat	\triangleright Newton loop for Ω_1
7: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$	⊳ linear system
8: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{x}_B^{(1)} + riangle \boldsymbol{x}_B^{(1)}$	
9: until $ \triangle \boldsymbol{x}_{B}^{(1)} / \boldsymbol{x}_{B}^{(1)} \le \epsilon$	\triangleright loose tolerance, e.g. $\epsilon \in [10^{-4}, 10^{-1}]$
10: $\boldsymbol{y}^{(2)} \leftarrow \boldsymbol{x}^{(2)}_B$	⊳ for convergence check
11: $\boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{P}_{21}\boldsymbol{x}_{\beta}^{(1)} + \boldsymbol{Q}_{21}\boldsymbol{x}_{\beta}^{(1)} + \boldsymbol{G}_{21}\boldsymbol{x}_{\beta}^{(1)}$	\triangleright project from Ω_1 to Γ_2
12: repeat	\triangleright Newton loop for Ω_2
13: $ \Delta \boldsymbol{x}_{B}^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{\beta}^{(2)}) \setminus \boldsymbol{R}_{A}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{\beta}^{(2)}) $	⊳ solve linear system
14: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{x}_B^{(2)} + \bigtriangleup \boldsymbol{x}_B^{(2)}$	
15: until $ \triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \le \epsilon$	\triangleright loose tolerance, e.g. $\epsilon \in [10^{-4}, 10^{-1}]$
$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	⊳ tight tolerance

Inexact Schwarz

$1: \boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)} \text{ in } \Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_1, \boldsymbol{x}_\beta^{(1)} \leftarrow \boldsymbol{X}_\beta^{(1)} \text{ on } \Gamma_1$	\triangleright initialize for Ω_1
2: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{X}_B^{(2)}$ in $\Omega_2, \boldsymbol{x}_b^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_2, \boldsymbol{x}_\beta^{(2)} \leftarrow \boldsymbol{X}_\beta^{(2)}$ on Γ_2	\triangleright initialize for Ω_2
3: repeat	⊳ Newton-Schwarz loop
$\begin{vmatrix} 4: & \bm{x}_{\beta}^{(1)} \leftarrow \bm{P}_{12}\bm{x}_{B}^{(2)} + \bm{Q}_{12}\bm{x}_{b}^{(2)} + \bm{G}_{12}\bm{x}_{\beta}^{(2)} \end{vmatrix}$	\triangleright project from Ω_2 to Γ_1
5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)})$	⊳ linear system
6: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{x}_B^{(1)} + riangle \boldsymbol{x}_B^{(1)}$	
7: $m{x}_eta^{(2)} \leftarrow m{P}_{21}m{x}_B^{(1)} + m{Q}_{21}m{x}_b^{(1)} + m{G}_{21}m{x}_eta^{(1)}$	\triangleright project from Ω_1 to Γ_2
8: $\Delta \boldsymbol{x}_B^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_B^{(2)}; \boldsymbol{x}_b^{(2)}; \boldsymbol{x}_\beta^{(2)}) \setminus \boldsymbol{R}_A^{(2)}(\boldsymbol{x}_B^{(2)}; \boldsymbol{x}_b^{(2)}; \boldsymbol{x}_\beta^{(2)})$	⊳ linear system
9: $oldsymbol{x}_B^{(2)} \leftarrow oldsymbol{x}_B^{(2)} + riangle oldsymbol{x}_B^{(2)}$	
$\left 10: \text{ until } \left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \right)^2 \right]^{1/2} \le \epsilon_{\text{machine}}$	⊳ tight tolerance

Modified Schwarz



Monolithic Schwarz

*A. Mota, I. Tezaur, C. Alleman. "The Schwarz Alternating Method in Solid Mechanics", CMAME 319 (2017), 19-51.

Appendix. Full Schwarz Method



1: $\boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)}$ in $\Omega_{1}, \boldsymbol{x}_{b}^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{1}, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_{1} \triangleright initialize for Ω_1 2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} \triangleright initialize for Ω_2 3: repeat ▷ Schwarz loop $oldsymbol{u}^{(1)} \leftarrow oldsymbol{x}_{\mathrm{D}}^{(1)}$ 4: \triangleright for convergence check $m{x}_eta^{(1)} \leftarrow m{P}_{12} m{x}_B^{(2)} + m{Q}_{12} m{x}_b^{(2)} + m{G}_{12} m{x}_eta^{(2)}$ 5: \triangleright project from Ω_2 to Γ_1 \triangleright Newton loop for Ω_1 6: $\triangle \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)}) \backslash \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)})$ 7: \triangleright linear system $oldsymbol{x}_B^{(1)} \leftarrow oldsymbol{x}_B^{(1)} + riangle oldsymbol{x}_B^{(1)}$ 8: until $|| \triangle \boldsymbol{x}_{B}^{(1)} || / || \boldsymbol{x}_{B}^{(1)} || \leq \epsilon_{\text{machine}}$ 9: \triangleright tight tolerance $oldsymbol{y}^{(2)} \leftarrow oldsymbol{x}^{(2)}_B$ 10: \triangleright for convergence check $oldsymbol{x}_eta^{(2)} \leftarrow oldsymbol{P}_{21}oldsymbol{x}_B^{(1)} + oldsymbol{Q}_{21}oldsymbol{x}_b^{(1)} + oldsymbol{G}_{21}oldsymbol{x}_eta^{(1)}$ 11: \triangleright project from Ω_1 to Γ_2 12: \triangleright Newton loop for Ω_2 repeat $\triangle \bm{x}_B^{(2)} \leftarrow -\bm{K}_{AB}^{(2)}(\bm{x}_B^{(2)};\bm{x}_b^{(2)};\bm{x}_\beta^{(2)}) \backslash \bm{R}_A^{(2)}(\bm{x}_B^{(2)};\bm{x}_b^{(2)};\bm{x}_\beta^{(2)})$ 13: \triangleright linear system $oldsymbol{x}_{R}^{(2)} \leftarrow oldsymbol{x}_{R}^{(2)} + riangle oldsymbol{x}_{R}^{(2)}$ 14: until $|| \triangle \boldsymbol{x}_{B}^{(2)} || / || \boldsymbol{x}_{B}^{(2)} || \leq \epsilon_{\text{machine}}$ 15: \triangleright tight tolerance 16: until $\left[\left(||\boldsymbol{y}^{(1)} - \boldsymbol{x}_{B}^{(1)}||/||\boldsymbol{x}_{B}^{(1)}||\right)^{2} + \left(||\boldsymbol{y}^{(2)} - \boldsymbol{x}_{B}^{(2)}||/||\boldsymbol{x}_{B}^{(2)}||\right)^{2}\right]^{1/2} \leq \epsilon_{\text{machine}}$ \triangleright tight tolerance



Appendix. Inexact Schwarz Method

Classical algorithm originally proposed by Schwarz with *outer Schwarz loop* and *inner Newton loop*, with Newton step converged to a *loose tolerance*.

1: $\boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)}$ in $\Omega_{1}, \boldsymbol{x}_{b}^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{1}, \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)}$ on Γ_{1} \triangleright initialize for Ω_1 2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} \triangleright initialize for Ω_2 3: repeat ▷ Schwarz loop $oldsymbol{y}^{(1)} \leftarrow oldsymbol{x}^{(1)}_{\scriptscriptstyle B}$ 4: \triangleright for convergence check $m{x}_eta^{(1)} \leftarrow m{P}_{12} m{x}_B^{(2)} + m{Q}_{12} m{x}_b^{(2)} + m{G}_{12} m{x}_B^{(2)}$ 5: \triangleright project from Ω_2 to Γ_1 6: \triangleright Newton loop for Ω_1 repeat $\triangle \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)}) \backslash \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)})$ 7: \triangleright linear system $\boldsymbol{x}_{P}^{(1)} \leftarrow \boldsymbol{x}_{P}^{(1)} + \bigtriangleup \boldsymbol{x}_{P}^{(1)}$ 8: until $|| riangle oldsymbol{x}_B^{(1)} || / || oldsymbol{x}_B^{(1)} || \le \epsilon$ 9: \triangleright loose tolerance, e.g. $\epsilon \in [10^{-4}, 10^{-1}]$ $oldsymbol{y}^{(2)} \leftarrow oldsymbol{x}_{\scriptscriptstyle P}^{(2)}$ 10: \triangleright for convergence check $m{x}_eta^{(2)} \leftarrow m{P}_{21} m{x}_B^{(1)} + m{Q}_{21} m{x}_b^{(1)} + m{G}_{21} m{x}_eta^{(1)}$ 11: \triangleright project from Ω_1 to Γ_2 12: \triangleright Newton loop for Ω_2 repeat $\triangle \boldsymbol{x}_{B}^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{B}^{(2)}) \backslash \boldsymbol{R}_{A}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{B}^{(2)})$ 13: \triangleright solve linear system $oldsymbol{x}_{\scriptscriptstyle P}^{(2)} \leftarrow oldsymbol{x}_{\scriptscriptstyle P}^{(2)} + riangle oldsymbol{x}_{\scriptscriptstyle P}^{(2)}$ 14: until $|| \triangle \boldsymbol{x}_{B}^{(2)} || / || \boldsymbol{x}_{B}^{(2)} || \leq \epsilon$ \triangleright loose tolerance, e.g. $\epsilon \in [10^{-4}, 10^{-1}]$ 15: 16: **until** $\left[\left(||\boldsymbol{y}^{(1)} - \boldsymbol{x}_{B}^{(1)}||/||\boldsymbol{x}_{B}^{(1)}||\right)^{2} + \left(||\boldsymbol{y}^{(2)} - \boldsymbol{x}_{B}^{(2)}||/||\boldsymbol{x}_{B}^{(2)}||\right)^{2}\right]^{1/2} \leq \epsilon_{\text{machine}}$ \triangleright tight tolerance



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Appendix. Monolithic Schwarz Method

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Combines Schwarz and Newton loop into *since Newton-Schwarz loop*, with *elimination of Schwarz boundary DOFs*, and tight convergence tolerance.

 $\begin{array}{ll} 1: \ \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)} \text{ in } \Omega_{1}, \ \boldsymbol{x}_{b}^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(1)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_{1}, & \triangleright \text{ initialize for } \Omega_{1} \\ 2: \ \boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)} \text{ in } \Omega_{2}, \ \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_{2}, & \triangleright \text{ initialize for } \Omega_{2} \\ 3: \ \mathbf{repeat} & \triangleright \text{ Newton-Schwarz loop} \\ 4: \quad \left\{ \begin{array}{c} \Delta \boldsymbol{x}_{B}^{(1)} \\ \Delta \boldsymbol{x}_{B}^{(2)} \end{array} \right\} \leftarrow \left(\begin{array}{c} \boldsymbol{K}_{AB}^{(1)} + \boldsymbol{K}_{A\beta}^{(1)} \boldsymbol{H}_{11} & \boldsymbol{K}_{A\beta}^{(1)} \boldsymbol{H}_{12} \\ \boldsymbol{K}_{A\beta}^{(2)} \boldsymbol{H}_{21} & \boldsymbol{K}_{A\beta}^{(2)} + \boldsymbol{K}_{A\beta}^{(2)} \boldsymbol{H}_{22} \end{array} \right) \setminus \left\{ \begin{array}{c} -\boldsymbol{R}_{A}^{(1)} \\ -\boldsymbol{R}_{A}^{(1)} \end{array} \right\} & \triangleright \text{ linear system} \\ 5: \quad \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{x}_{B}^{(1)} + \Delta \boldsymbol{x}_{B}^{(1)} \\ 6: \quad \boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{x}_{B}^{(2)} + \Delta \boldsymbol{x}_{B}^{(2)} \\ 7: \ \text{ until } \left[\left(||\Delta \boldsymbol{x}_{B}^{(1)}||/||\boldsymbol{x}_{B}^{(1)}|| \right)^{2} + \left(||\Delta \boldsymbol{x}_{B}^{(2)}||/||\boldsymbol{x}_{B}^{(2)}|| \right)^{2} \right]^{1/2} \leq \epsilon_{\text{machine}} \qquad \triangleright \text{ tight tolerance} \end{array} \right.$

Advantages:

By-passes Schwarz loop.

Disadvantages:

• Off-diagonal coupling terms → block linear solver is needed.

Appendix. Modified Schwarz Method



Combines Schwarz and Newton loop into *since Newton-Schwarz loop*, with *Schwarz boundaries* at *Dirichlet boundaries* and tight convergence tolerance.

1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_1, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_1	\triangleright initialize for Ω_1
2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} 3: repeat $(1) \qquad \boldsymbol{\Sigma} \qquad (2) \qquad \boldsymbol{\Omega} \qquad (2) \qquad \boldsymbol{\Omega} \qquad (2)$	▷ initialize for Ω_2 ▷ Newton-Schwarz loop
4: $\boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{P}_{12} \boldsymbol{x}_{B}^{(2)} + \boldsymbol{Q}_{12} \boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12} \boldsymbol{x}_{\beta}^{(2)}$	\triangleright project from Ω_2 to Γ_1
5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$	⊳ linear system
6: $oldsymbol{x}_B^{(1)} \leftarrow oldsymbol{x}_B^{(1)} + riangle oldsymbol{x}_B^{(1)}$	
7: $m{x}_{eta}^{(2)} \leftarrow m{P}_{21}m{x}_{B}^{(1)} + m{Q}_{21}m{x}_{b}^{(1)} + m{G}_{21}m{x}_{eta}^{(1)}$	\triangleright project from Ω_1 to Γ_2
8: $\Delta m{x}_B^{(2)} \leftarrow -m{K}_{AB}^{(2)}(m{x}_B^{(2)};m{x}_b^{(2)};m{x}_\beta^{(2)}) ackslash m{R}_A^{(2)}(m{x}_B^{(2)};m{x}_b^{(2)};m{x}_\beta^{(2)})$	⊳ linear system
9: $oldsymbol{x}_B^{(2)} \leftarrow oldsymbol{x}_B^{(2)} + riangle oldsymbol{x}_B^{(2)}$	
10: until $\left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \right)^2\right]^{1/2} \le \epsilon_{\text{machine}}$	⊳ tight tolerance

Advantages:

- By-passes Schwarz loop.
- No diagonal coupling (conventional linear solver can be used in each subdomain).

Least-intrusive variant: by-passes Schwarz iteration, no need for block solver.

Appendix. Foulk's Singular Bar

- **1D proof of concept** problem:
 - **1D bar** with area proportional to square root of length.
 - Strong *singularity* on left end of bar.
 - Simple *hyperelestic* material model with no damage.
 - MATLAB implementation.



- Problem goals:
 - Explore *viability* of *4 variants* of the Schwarz alternating method.
 - Test *convergence* and compare with literature (Evans, 1986).
 - Expect *faster convergence* in *fewer iterations* with *increased overlap*.





Appendix. Singular Bar and Schwarz Varia



Appendix. Multiscale Modeling of Localization







Strain localization can cause *localized necking* (left) and ultimately *fracture* (above).

Goals:

- Connect *physical length scales* to *engineering scale models*.
- Investigate importance of *microstructural detail.*
- Develop bridging technologies for *spatial multiscale/ multiphysics*.

Appendix. Parallelization via DTK: Weak 🙃 Scaling on Cubes Problem



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Appendix. Parallelization via DTK: Strong Scaling on Cubes Problem



Appendix. Notched Cylinder: HEX-HEX Coupling



Appendix. Notched Cylinder: Nonconformal HEX - HEX Coupling



Appendix. Notched Cylinder: Nonconformal HEX - HEX Coupling



	u_3 relati	ve error
Absolute residual tolerance	Ω_1	Ω_2
$egin{array}{c} 1.0 imes10^{-8}\ 1.0 imes10^{-12}\ 1.0 imes10^{-14}\ 2.5 imes10^{-16} \end{array}$	$\begin{array}{c} 1.31 \times 10^{-3} \\ 1.30 \times 10^{-3} \\ 1.30 \times 10^{-3} \\ 1.30 \times 10^{-3} \end{array}$	$\begin{array}{c} 4.45 \times 10^{-4} \\ 4.43 \times 10^{-4} \\ 4.43 \times 10^{-4} \\ 4.43 \times 10^{-4} \end{array}$

A	bany
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1.446e-05

Appendix. Rubiks Cube Problem





Appendix. Tensile Bar





ASTM tensile geometry
Appendix. Tensile Bar: Meso-Macroscale Coupling

Mesoscale

SPARKS-generated microstructure (F. Abdeljawad)



ensembles of

cubic elastic constant : $C_{11} = 204.6$ GPa cubic elastic constant : $C_{12} = 137.7$ GPa cubic elastic constant : $C_{44} = 126.2$ GPa reference shear rate : $\dot{\gamma}_0 = 1.0 \ 1/s$

rate sensitivity factor : m = 20

hardening rate parameter : $\dot{q}_0 = 2.0 \times 10^4 \text{ 1/s}$

initial hardness : $g_0 = 90$ MPa

saturation hardness : $g_s = 202 \text{ MPa}$

saturation exponent : $\omega = 0.01$

Fix microstructure, investigate ensembles





Load microstructural ensembles in uniaxial stress Fit flow curves with a macroscale J₂ plasticity model



$$\sigma_y = \sigma_0 + H\epsilon_p + S(1 - e^{-\alpha\epsilon_p})$$

Appendix. Tensile Bar: Results





Appendix. Schwarz Alternating Method for Dynamics

- In the literature the Schwarz method is applied to dynamics by using *space-time discretizations*.
- This was deemed *unfeasible* given the design of our current codes and size of simulations.



Overlapping non-matching meshes and time steps in dynamics.

Appendix. A Schwarz-like Time Integrator 🛅 National Laboratories

- We developed an *extension of Schwarz coupling* to *dynamics* using a governing time stepping algorithm that controls time integrators within each domain.
- Can use *different integrators* with *different time steps* within each domain.
- 1D results show *smooth coupling without numerical artifacts* such as spurious wave reflections at boundaries of coupled domains.



Schwarz Alternating Method for Dynamic Multiscale Coupling: Theory

- Our dynamic alternating Schwarz algorithm can be *interpreted* as applying the traditional Schwarz alternating iterations in *space-time* between $\Omega_1 \times I_k$ and $\Omega_2 \times I_k$
- This interpretation *does not* require the method to be implemented using a spacetime framework.



Appendix. Sierra/SM: Dynamic Schwarz Coupling



begin sierra schwarz_tension begin material aluminum density = 2700.0begin parameters for model fefp youngs modulus = 70.0e+09poissons ratio = 0.36 yield stress = 250.0e+06 hardening model = linear hardening modulus = 0.7e+09end end begin function applied_disp_top_ramp type = analytic evaluate expression = " $0.01 \times \cos ramp(x, 0.0, 1.0)$ " end begin function applied_disp_bottom_ramp type = analyticevaluate expression = $"-0.01 * \cos ramp(x, 0.0, 1.0)"$ end begin total lagrange section gauge_formulation strain incrementation = logarithmic_mapping cubature degree = 3volume average J = onelement jacobian = element probe end begin total lagrange section ends formulation strain incrementation = logarithmic mapping cubature degree = 3volume average J = onelement jacobian = element_probe end begin finite element model gauge_mesh database name = gauge.g database type = exodusII begin parameters for block gauge material = aluminum model = fefpsection = gauge_formulation end end begin finite element model ends mesh database name = ends.g database type = exodusII **begin** parameters for block ends material = aluminum model = fefpsection = ends_formulation end end {include("schwarz_funcs.i")}

tension.i

begin procedure schwarz_procedure

begin solution control description schwarz_loop use system main begin system main begin transient schwarz_solve begin nonlinear schwarz_iteration advance ends_region transfer from_ends_to_gauge advance gauge_region transfer from_gauge_to_ends end nonlinear schwarz_iteration end transient schwarz solve end system main begin parameters for transient schwarz_solve start time = 0.0 termination time = 1.0 begin parameters for adagio region ends_region time increment = 0.01 end begin parameters for adagio region gauge_region time increment = 0.01end begin parameters for nonlinear schwarz_iteration converged when "gauge_region.schwarz_norm < 1.0e-06 && ends_region.schwarz_norm < 1.0e-06 end end solution control description schwarz loop begin transfer from ends to gauge interpolate volume nodes from ends_region to gauge_region send block ends to schwarz_upper_transition schwarz_lower_transition send field displacement state new to schwarz_disp state none send field velocity state new to velocity state new send field acceleration state new to acceleration state new end begin transfer from_gauge_to_ends interpolate volume nodes from gauge_region to ends_region send block gauge to schwarz_upper_gauge schwarz_lower_gauge send field displacement state new to schwarz_disp state none send field velocity state new to velocity state new send field acceleration state new to acceleration state new end {include("gauge.i")}
{include("ends.i")} end procedure schwarz_procedure

begin feti equation solver feti
residual norm tolerance = 1e-3
maximum orthogonalization = 5000
end

end sierra schwarz_tension

Appendix. Sierra/SM: Dynamic Schwarz Coupling





use finite element model ends_mesh begin implicit dynamics end begin adaptive time stepping cutherk factor = 0 5

region ends_regio

cutback factor = 0.5 growth factor = 1.1 maximum failure cutbacks = 16 end

{include("schwarz_vars.i")}

begin results output database name = ends.e database type = exodusII at time 0, increment = 0.01 element variables = stress as str element variables = log_strain element variables = log_strain element variables = von_mises nodal variables = von_mises nodal variables = volocity nodal variables = acceleration nodal variables = coordinates global variables = timestep end

{include("solver.i")}

begin fixed displacement
 node set = upper_grip
 component = x
end

begin fixed displacement
 node set = upper_grip
 component = z
end

begin prescribed displacement node set = upper_grip function = applied_disp_top_ramp component = y end

begin fixed displacement
 node set = lower_grip
 component = x
end

begin fixed displacement
 node set = lower_grip
 component = z
end

begin prescribed displacement
node set = lower_grip
function = applied_disp_bottom_ramp
component = y
end

begin prescribed displacement schwarz_bc_x node set = schwarz_upper_gauge schwarz_lower_gauge component = x function = schwarz_disp_x end

begin prescribed displacement schwarz_bc_y node set = schwarz_upper_gauge schwarz_lower_gauge component = y function = schwarz_disp_y end

begin prescribed displacement schwarz_bc_z node set = schwarz_upper_gauge schwarz_lower_gauge component = z function = schwarz_disp_z end

and adagio region ends_region

ends.i



begin implicit dynamics end begin adaptive time stepping cutback factor = 0.5

egin adagio region gauge_region use finite element model gauge mesh

growth factor = 1.1 maximum failure cutbacks = 16 end

{include("schwarz_vars.i")}

begin results output database name = gauge.e database type = exodusII at time 0, increment = 0.01 element variables = stress as stress element variables = cauchy_stress element variables = cauchy_stress element variables = cauchy_stress nodal variables = von_mises nodal variables = volcity nodal variables = acceleration nodal variables = timestep end

{include("solver.i")}

begin prescribed displacement schwarz_bc_x node set = schwarz_upper_transition schwarz_lower_transition component = x function = schwarz_disp_x end

begin prescribed displacement schwarz_bc_y node set = schwarz_upper_transition schwarz_lower_transition component = y function = schwarz_disp_y end

begin prescribed displacement schwarz_bc_z node set = schwarz_upper_transition schwarz_lower_transition component = z function = schwarz_disp_z

d adagio region gauge_region

end

gauge.i

Appendix. Sierra/SM: Dynamic Schwarz Coupling



#		
# Functions for transfer of Schwarz BCs #		
<pre>begin function schwarz_disp_x type = analytic</pre>	schwarz tuncs.	
expression variable: dx = nodal schwarz_disp(x)		
end		
<pre>begin function schwarz_disp_y type = analytic</pre>		<pre>begin user variable schwarz_disp type = node vector</pre>
expression variable: dy = nodal schwarz_disp(y) evaluate expression = "dy" ord		end
begin function schwarz_disp_z		<pre>begin user output compute nodal delta_x as funct</pre>
<pre>type = analytic expression variable: dz = nodal schwarz_disp(z) evaluate expression = "dz" end</pre>		compute nodal delta_y as funct compute nodal delta_z as funct compute global delta_x_norm as
# # Define functions for building the Schwarz convergence criterion		compute global delta_y_norm as compute global delta_z_norm as
# begin function x_prev type = analytic	schwarz vars.i	compute global curr_y_norm as compute global curr z norm as
expression variable: x_prev = nodal coordinates(x) evaluate expression = "x_prev" end	—	compute global schwarz_norm as compute at every step end
<pre>begin function y_prev type = analytic</pre>		begin user output
<pre>expression variable: y_prev = nodal coordinates(y) evaluate expression = "y_prev" end</pre>		compute nodal x_prev as functi compute nodal y_prev as functi compute nodal z prev as functi
<pre>begin function z_prev type = analytic</pre>		compute at every step
expression variable: z_prev = nodal coordinates(z) evaluate expression = "z_prev" end		
begin function delta_x		<pre>\${maximum_iterations = 150} \${tangent_iterations = 125}</pre>
expression variable: x_curr = nodal coordinates(x) expression variable: x_prev = nodal x_prev		<pre>\${smooth_iterations = maximum_it howin column</pre>
evaluate expression = "x_curr - x_prev" end		begin solver
<pre>begin function delta_y type = analytic expression variable: y_curr = nodal coordinates(y) expression variable: y_prev = nodal y_prev evaluate expression = "y curr - y new"</pre>		<pre>begin cg iteration print = 100 maximum iterations = {maximum minimum iterations = 2</pre>
end		target relative residual
<pre>begin function delta_z type = analytic expression variable: z_curr = nodal coordinates(z)</pre>	solver.l	target residual acceptable residual acceptable relative residual
expression variable: z_prev = nodal z_prev evaluate expression = "z_curr - z_prev" end		begin full tangent preconditi
begin function schwarz_norm		linear solver = feti
expression variable: xn = global curr_x_norm		small number of iterations
expression variable: zn = global curr_z_norm		end end cg
expression variable: dx = global delta_x_norm expression variable: dy = global delta_y_norm		
expression variable: dz = global delta_z_norm evaluate expression = "sgrt(dx*dx+dv*dv+dz*dz)/sgrt(xn*xn+vn*vn+zn*zn)"		end solver
end		

type = node vector gin user output compute nodal delta_x as function delta_x compute nodal delta_y as function delta_y compute nodal delta_z as function delta_z compute global delta_x_norm as l2norm of nodal delta_x compute global delta_y_norm as l2norm of nodal delta_y compute global delta_z_norm as l2norm of nodal delta_z compute global curr_x_norm as l2norm of nodal coordinates(x) compute global curr_y_norm as l2norm of nodal coordinates(y) compute global curr_z_norm as l2norm of nodal coordinates(z) compute global schwarz_norm as function schwarz_norm compute at every step gin user output compute nodal x_prev as function x_prev compute nodal y_prev as function y_prev compute nodal z_prev as function z_prev compute at every step aximum_iterations = 150} angent_iterations = 125} smooth iterations = maximum iterations - tangent iterations} ain solver begin cg iteration print = 100 maximum iterations = {maximum_iterations} minimum iterations = 2 = 1e-6 = 1e-8 target residual = 1e-3 acceptable relative residual = 1e-3**begin** full tangent preconditioner conditioning = no_check small number of iterations = 100 end end cg

Appendix. Dynamic Singular Bar



- Inelasticity masks problems by introducing *energy dissipation*.
- Schwarz does not introduce numerical artifacts.
- Can couple domains with *different time integration schemes* (*Explicit-Implicit* below).



Appendix. Elastic Wave: Energy Conservation



- For clamped beam problem, total energy (TE = $0.5x^TKx + 0.5\dot{x}^TM\dot{x}$) should be conserved.
- Total energy is calculated in 2 ways: with most of contribution from Ω_0 and from Ω_1 .

Max / Avg # Schwarz iterations during run

Albani

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54

Max / Avg # Schwarz iterations during run

Albany

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54

As Schwarz tolerance is **tightened**, number of Schwarz **iterations goes up** (as expected)

Max / Avg # Schwarz iterations during run

Albany

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54



Left: time vs. # Schwarz iterations (per time step) – behavior is as expected given nature of solution/domain decomposition

Max / Avg # Schwarz iterations during run

Albani

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54





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Max / Avg # Schwarz iterations during run

Albany

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54





lia onal ratories

Max / Avg # Schwarz iterations during run

FFF Albany

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54





Max / Avg # Schwarz iterations during run

Albany

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54





Max / Avg # Schwarz iterations during run

Albani

	Impl-impl dt=1e-6, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-6	Impl-impl dt=1e-7, Schwarz tol=1e-10	Impl-impl dt=1e-7, Schwarz tol=1e-15	Expl-impl dt=1e-7 Schwarz tol=1e-6	Expl-impl dt=1e-8 Schwarz tol=1e-6
Hex-Hex	3 / 2.23	3 / 2.08		4 / 2.83	2 / 2.0	
Tet-Hex					2 / 2.0	
Nonconformal Hex-Hex		2 / 2.0	3 / 2.36		2 / 2.0	2 / 1.54







- Left figure shows *# of iterations* as a function of *overlap region size* for 2 subdomains. The method does not converge for 0% overlap. If the overlap is 100% then the single-domain solution is recovered for each of the subdomains.
- Right figure shows *linear convergence rate* of dynamic Schwarz implementation (for small overlap fraction of 0.2%).

Appendix. Torsion Problem

- Nonlinear elastic bar (Neohookean material model) subjected to a high degree of *torsion*.
- The *domain* is $\Omega = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.5, 0.5).$
- We evaluate *dynamic Schwarz* with 2 subdomains: $\Omega_0 = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.5, 0.25), \Omega_1 = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.25, 0.5).$
- **Time-discretizations:** Newmark (implicit, explicit) with same Δt .
- *Meshes:* HEX, Composite TET10.





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Appendix. Torsion: HEX - HEX Coupling

- Each subdomain discretized using **uniform HEX mesh** with $\Delta x_i = 0.01$, and advanced in time using implicit Newmark-Beta scheme with $\Delta t = 1e-6$.
- Results compared to single-domain solution on mesh conformal with Schwarz domain meshes.





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 $\Omega_{\rm ref}$

 Ω_0



Appendix. Bolted Joint Problem: Performance



- Schwarz *tolerance* = 1e-6
- Dynamic Schwarz converges in between 2-4 Schwarz iterations per time-step despite the overlap region being very small for this problem.

Appendix. Numerical Example: 1D Dynamic Wave Propagation Problem



• Alternating **Dirichlet-Neumann** Schwarz BCs with **no relaxation** ($\theta = 1$) on Schwarz boundary Γ



θ	Min # Schwarz Iters	Max # Schwarz Iters	Total # Schwarz Iters
1.10	3	9	59,258
1.00	1	4	24,630
0.99	1	5	35,384
0.95	3	6	45,302
0.90	3	8	56,114

> A parameter sweep study revealed $\theta = 1$ gave best performance (min # Schwarz iters)

- All couplings were implicit-implicit with Δt₁ = Δt₂ = ΔT = 10⁻⁷ and Δx₁ = Δx₂ = 10⁻³
 ➤ Time-step and spatial resolution chosen to be small enough to resolve propagating wave
- All reproductive and predictive cases run on the same RHEL8 and RHEL7 machines.
- Model accuracy evaluated w.r.t. analogous FOM-FOM coupling via mean square error (MSE):

$$\varepsilon_{MSE}(\widetilde{\boldsymbol{u}}_i) \coloneqq \frac{\sqrt{\sum_{n=1}^{S} ||\widetilde{\boldsymbol{u}}_i^n - \boldsymbol{u}_i^n||_2^2}}{\sqrt{\sum_{n=1}^{S} ||\boldsymbol{u}_i^n||_2^2}}$$