

The Schwarz alternating method for concurrent multiscale coupling in solid mechanics

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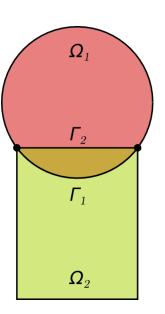
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Outline

- 1. Motivation
- 2. Schwarz Alternating Method: Background & History
- 3. Schwarz Alternating Method for Concurrent Multiscale Coupling in Quasistatics
 - Four Variants: Full Schwarz, Inexact Schwarz, Modified Schwarz, Monolithic Schwarz
 - Implementations: MATLAB, Albany
- 4. Numerical Examples
- 5. Summary
- 6. Future Work
- 7. References
- 8. Appendix





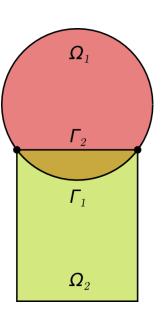


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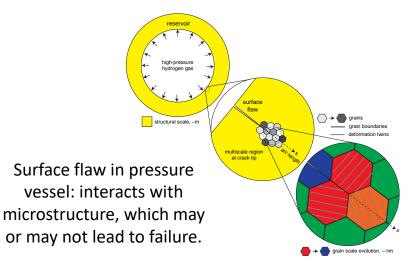
Motivation for Concurrent Multiscale Coupling

- Large scale structural failure frequently originates from small scale phenomena such as defects, microcracks, inhomogeneities and more, which grow quickly in unstable manner.
- Failure occurs due to *tightly coupled interaction* between small scale (stress concentrations, material instabilities, cracks, etc.) and large scale (vibration, impact, high loads and other perturbations).

Concurrent multiscale methods are essential for understanding and prediction of behavior of engineering systems when a small scale failure determines the performance of the entire system.



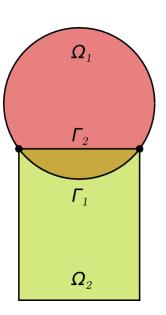
Roof failure of Boeing 737 aircraft due to fatigue cracks. From *imechanica.org*



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Schwarz Alternating Method for Domain Decomposition

Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Simple idea: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.

H. Schwarz (1843 – 1921)

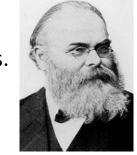
Initialize:

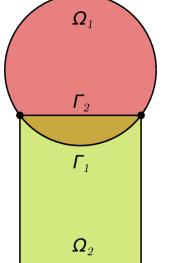
Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Schwarz Alternating Method

Iterate until convergence:

- Solve PDE by any method (can be different than for Ω_1) on Ω_2 w/ Dirichlet BCs on Γ_2 that are the values just obtained for Ω_1 .
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .





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Schwarz Alternating Method

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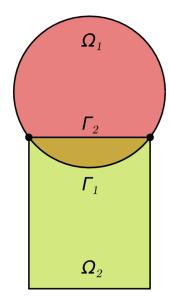
• Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Iterate until convergence:

Requirement for convergence: $\Omega_1 \cap \Omega_2 \neq \emptyset$

- Solve PDE by any method (can be different than for Ω₁) on Ω₂ w/ Dirichlet BCs on Γ₂ that are the values just obtained for Ω₁.
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .





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Schwarz Alternating Method

Initialize:

 Ω_1

 Γ_2

 Γ_1

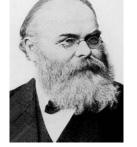
 Ω_2

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Iterate until convergence:

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- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .
- Schwarz alternating method most commonly used as a *preconditioner* for Krylov iterative methods to solve linear algebraic equations.





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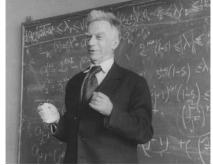
 <u>S. L. Sobolev (1936)</u>: posed Schwarz method for *linear elasticity* in variational form and *proved method's convergence* by proposing a convergent sequence of energy functionals.



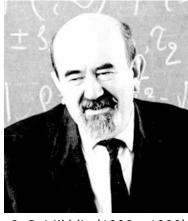
S. L. Sobolev (1908 - 1989)

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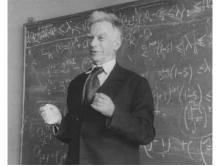


S. G. Mikhlin (1908 – 1990)

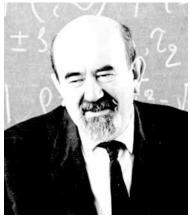
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- A. Mota, I. Tezaur, C. Alleman (2017)*: derived a proof of convergence of the alternating Schwarz method for the finite deformation quasi-static nonlinear PDEs (with energy functional Φ[φ] defined below), and determined a geometric convergence rate for the finite deformation quasi-static problem.

$$\boldsymbol{\Phi}[\boldsymbol{\varphi}] = \int_{B} W(\boldsymbol{F}, \boldsymbol{Z}, T) \, dV - \int_{B} \boldsymbol{B} \cdot \boldsymbol{\varphi} \, dV - \int_{\partial_{T}B} \overline{\boldsymbol{T}} \cdot \boldsymbol{\varphi} \, dS$$
$$\nabla \cdot \boldsymbol{P} + \boldsymbol{B} = \boldsymbol{0}$$



S. L. Sobolev (1908 – 1989)



S. G. Mikhlin (1908 – 1990)

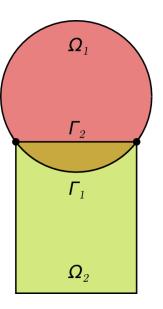


A. Mota, I. Tezaur, C. Alleman

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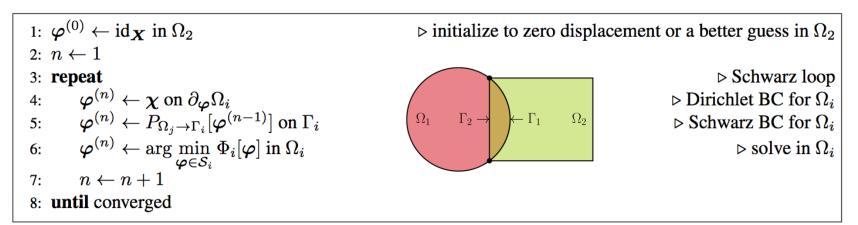
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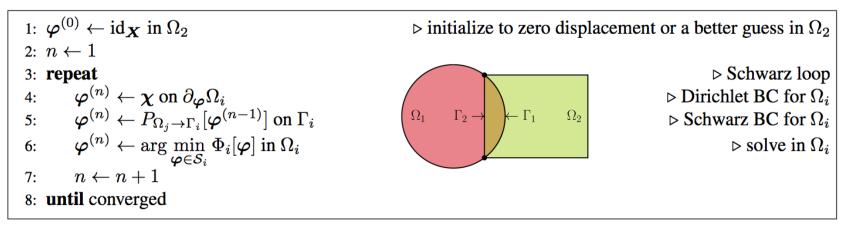


Schwarz Alternating Method for Multiscal Coupling in Quasistatics

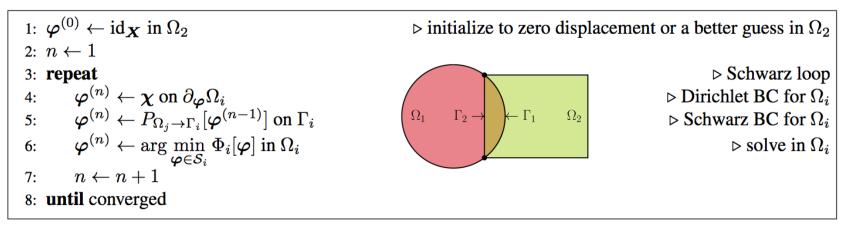


Advantages:

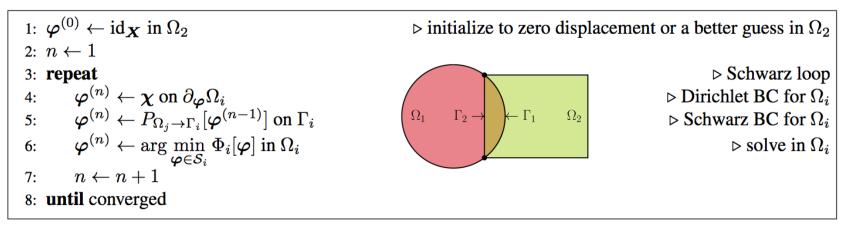
Conceptually very *simple*.



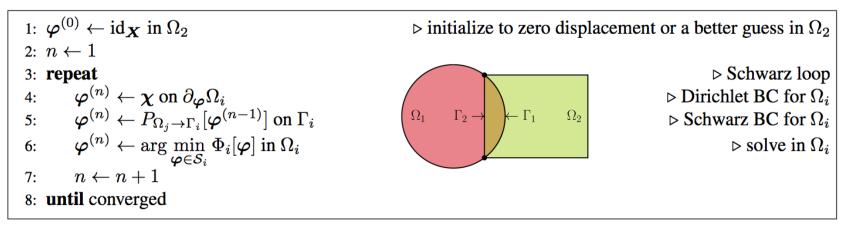
- Conceptually very *simple*.
- Allows the coupling of regions with *different non-conforming meshes*, *different element types*, and *different levels of refinement*.



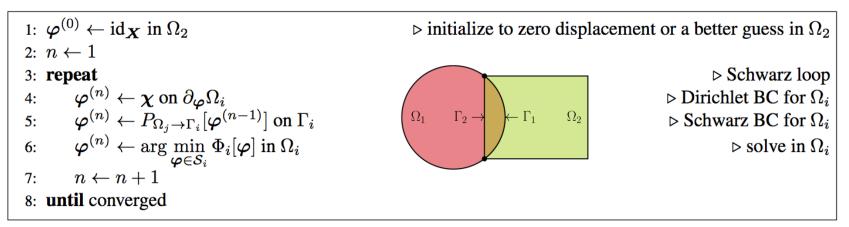
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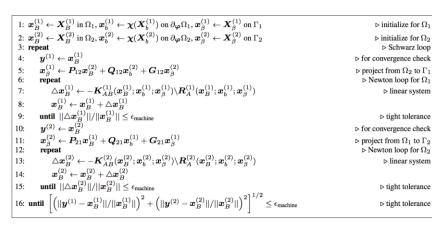


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- Different material models can be coupled provided that they are compatible in the overlap region.

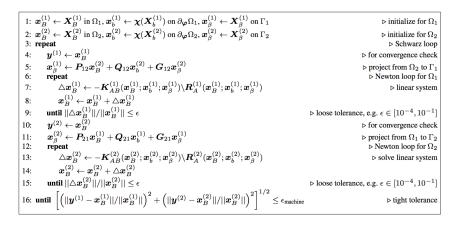


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- Information is exchanged among two or more regions, making coupling concurrent.
- *Different solvers* can be used for the different regions.
- Different material models can be coupled provided that they are compatible in the overlap region.
- Simplifies the task of *meshing complex geometries* for the different scales.

Four Variants* of the Schwarz Alternating ^{Sandia} Method

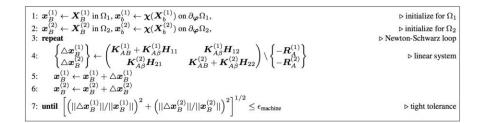


Full Schwarz



1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ or	
2: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{X}_B^{(2)}$ in $\Omega_2, \boldsymbol{x}_b^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(2)})$ of 3: repeat	$ n \partial_{\varphi} \Omega_2, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)} \text{ on } \Gamma_2 $ \triangleright initialize for Ω_2 \triangleright Newton-Schwarz loop
4: $\boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{P}_{12} \boldsymbol{x}_{B}^{(2)} + \boldsymbol{Q}_{12} \boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12}$	$\boldsymbol{x}_{\beta}^{(2)} \triangleright \text{ project from } \Omega_2 \text{ to } \Gamma_1$
5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})^{\prime}$	$\boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$ \triangleright linear system
6: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{x}_B^{(1)} + \Delta \boldsymbol{x}_B^{(1)}$	
7: $\boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{P}_{21} \boldsymbol{x}_{B}^{(1)} + \boldsymbol{Q}_{21} \boldsymbol{x}_{b}^{(1)} + \boldsymbol{G}_{21}$	$\boldsymbol{x}_{\beta}^{(1)} \triangleright \text{ project from } \Omega_1 \text{ to } \Gamma_2$
8: $\Delta m{x}_B^{(2)} \leftarrow -m{K}_{AB}^{(2)}(m{x}_B^{(2)};m{x}_b^{(2)};m{x}_\beta^{(2)})^{\prime}$	$\mathbf{R}_A^{(2)}(\boldsymbol{x}_B^{(2)}; \boldsymbol{x}_b^{(2)}; \boldsymbol{x}_\beta^{(2)})$ \triangleright linear system
9: $oldsymbol{x}_B^{(2)} \leftarrow oldsymbol{x}_B^{(2)} + riangle oldsymbol{x}_B^{(2)}$	
10: until $\left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(1)} \right)^2 \right]$	$\stackrel{2)}{ } \boldsymbol{x}_{B}^{(2)} \Big)^{2} \Big]^{1/2} \leq \epsilon_{\text{machine}} \qquad \qquad \triangleright \text{ tight tolerance}$

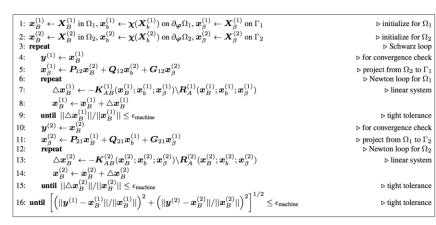
Modified Schwarz



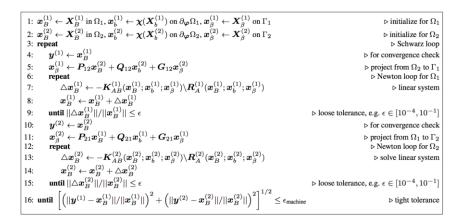
Monolithic Schwarz

Inexact Schwarz

Four Variants* of the Schwarz Alternating Method Least-intrusive variant: by-passes Schwarz iteration,



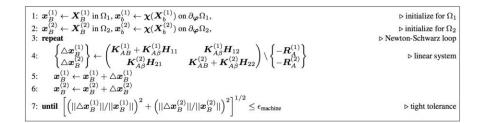
Full Schwarz



no need for block solver.

1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_1, \boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ on Γ_1	\triangleright initialize for Ω_1
2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2}	\triangleright initialize for Ω_2
3: repeat	Newton-Schwarz loop
4: $\boldsymbol{x}_{eta}^{(1)} \leftarrow \boldsymbol{P}_{12} \boldsymbol{x}_{B}^{(2)} + \boldsymbol{Q}_{12} \boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12} \boldsymbol{x}_{eta}^{(2)}$	\triangleright project from Ω_2 to Γ_1
5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$	⊳ linear system
6: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{x}_B^{(1)} + \bigtriangleup \boldsymbol{x}_B^{(1)}$	
7: $\boldsymbol{x}_{\beta}^{(2)} \leftarrow \bar{\boldsymbol{P}}_{21} \boldsymbol{x}_{B}^{(1)} + \bar{\boldsymbol{Q}}_{21} \boldsymbol{x}_{b}^{(1)} + \boldsymbol{G}_{21} \boldsymbol{x}_{\beta}^{(1)}$	\triangleright project from Ω_1 to Γ_2
8: $\Delta \boldsymbol{x}_B^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_B^{(2)}; \boldsymbol{x}_b^{(2)}; \boldsymbol{x}_\beta^{(2)}) \setminus \boldsymbol{R}_A^{(2)}(\boldsymbol{x}_B^{(2)}; \boldsymbol{x}_b^{(2)}; \boldsymbol{x}_\beta^{(2)})$	⊳ linear system
9: $\boldsymbol{x}_B^{(2)} \leftarrow \boldsymbol{x}_B^{(2)} + riangle \boldsymbol{x}_B^{(2)}$	
$10: \text{ until } \left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \right)^2 \right]^{1/2} \le \epsilon_{\text{machine}}$	⊳ tight tolerance

Modified Schwarz



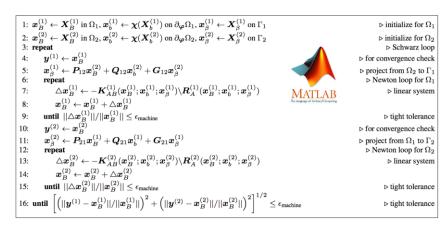
Monolithic Schwarz

Inexact Schwarz

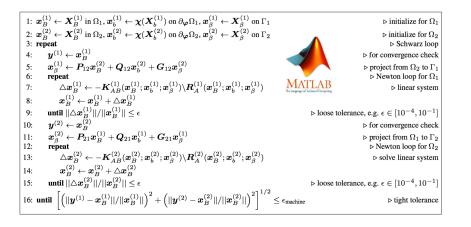
Implementations*



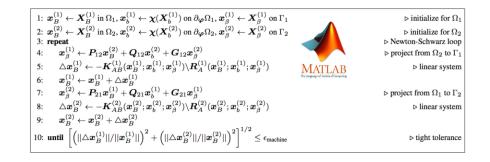
• All *four variants* implemented in *3D MATLAB* code.



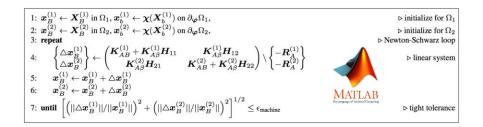
Full Schwarz



Inexact Schwarz



Modified Schwarz

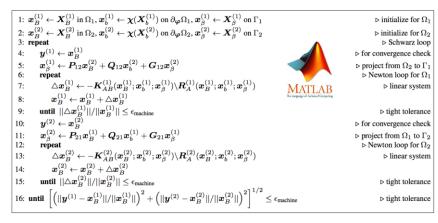


Monolithic Schwarz

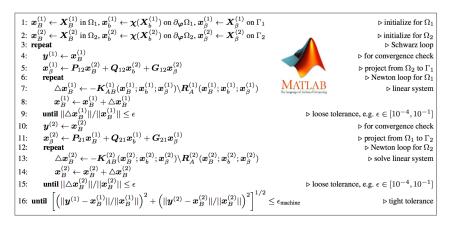
Implementations*

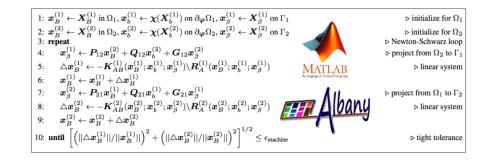


- All *four variants* implemented in *3D MATLAB* code.
- Modified & monolithic Schwarz variants implemented in parallel C++ Albany code.

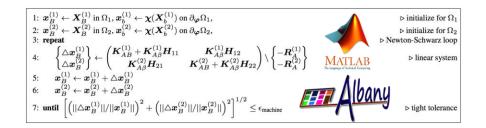


Full Schwarz





Modified Schwarz



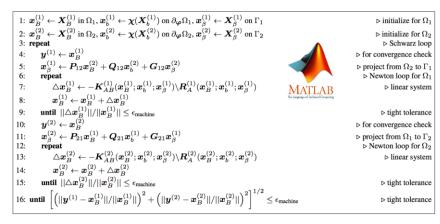
Inexact Schwarz

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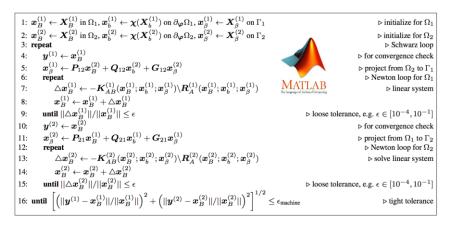
Implementations*

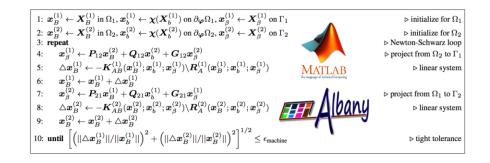


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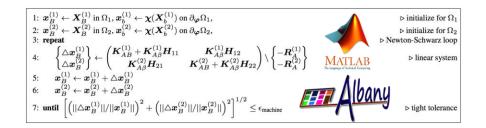


Full Schwarz





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Schwarz Alternating Method in Albany Code

Modified & monolithic Schwarz versions have been implemented within the *LCM project* in Sandia's open-source parallel, C++, multi-physics, finite element code, *Albany*.

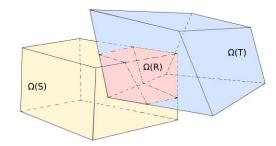
- Component-based design for rapid development of capabilities.
- Extensive use of libraries from the open-source *Trilinos* project.
 - Use of the *Phalanx* package to decompose complex problem into simpler problems with managed dependencies.
 - Use of the *Sacado* package for *automatic differentiation*.
 - Use of *Teko* package for *block preconditioning*.
- Parallel implementation of Schwarz alternating method uses the Data Transfer Kit (DTK).
- All software available on *GitHub*.

Albany

https://github.com/gahansen/Albany



https://github.com/trilinos/trilinos

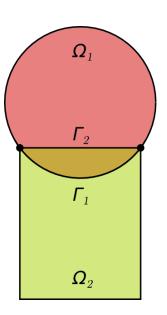


https://github.com/ORNL-CEES/DataTransferKit

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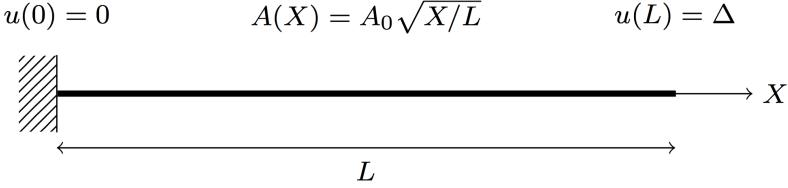






Example #1: Foulk's Singular Bar

- 1D proof of concept problem:
 - **1D bar** with area proportional to square root of length.
 - Strong *singularity* on left end of bar.
 - Simple *hyperelestic* material model with no damage.
 - MATLAB implementation.



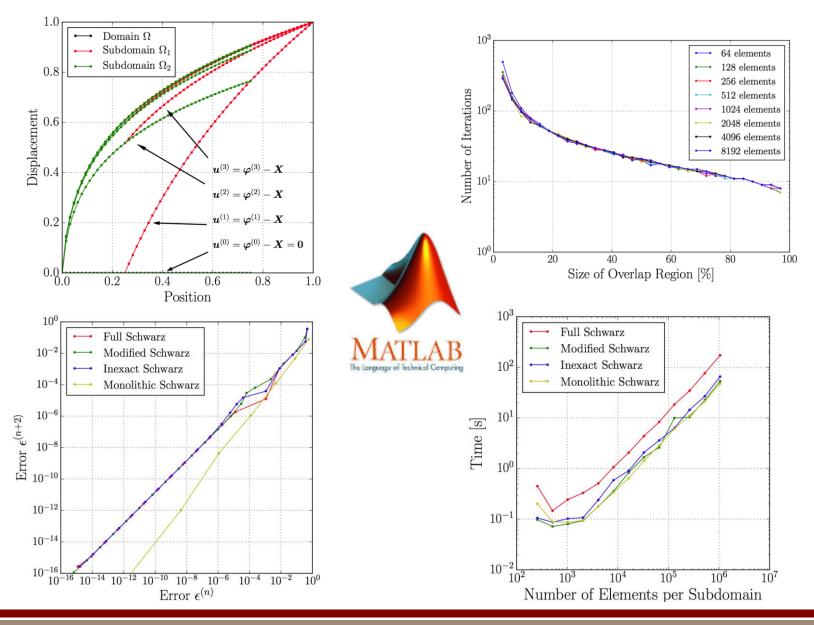
- Problem goals:
 - Explore *viability* of *4 variants* of the Schwarz alternating method.
 - Test *convergence* and compare with literature (Evans, 1986).
 - Expect *faster convergence* in *fewer iterations* with *increased overlap*.





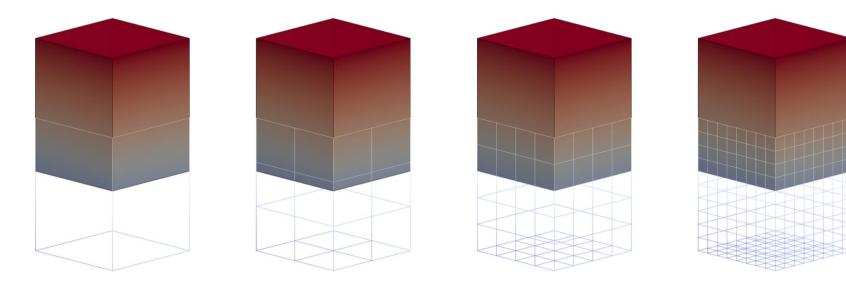
Singular Bar and Schwarz Variants



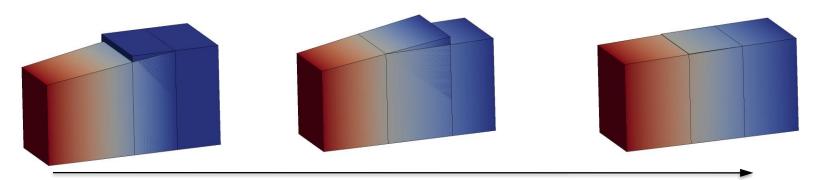


Example #2: Cuboid Problem



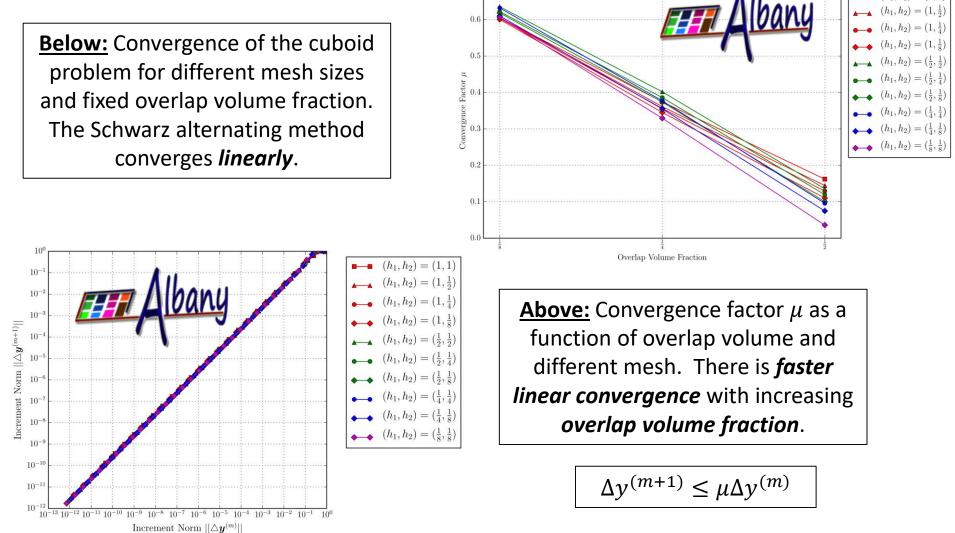


- Coupling of *two cuboids* with square base (above).
- *Neohookean*-type material model.



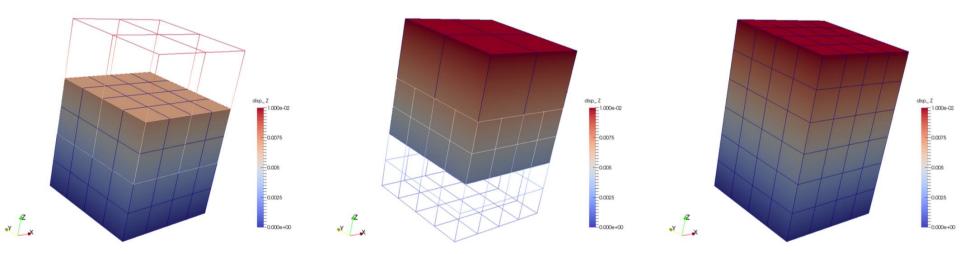
Combined Newton-Schwarz Iteration





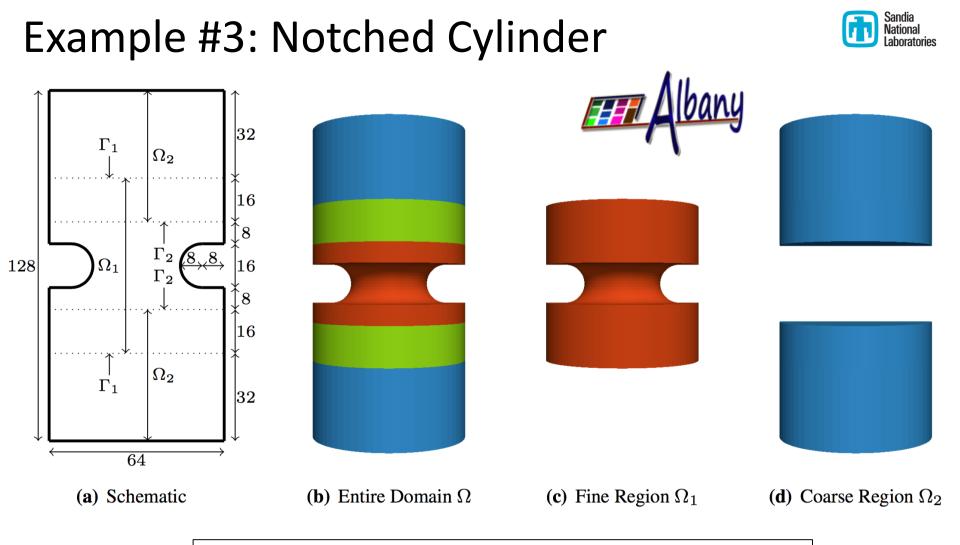
Cuboid Problem: Schwarz Error





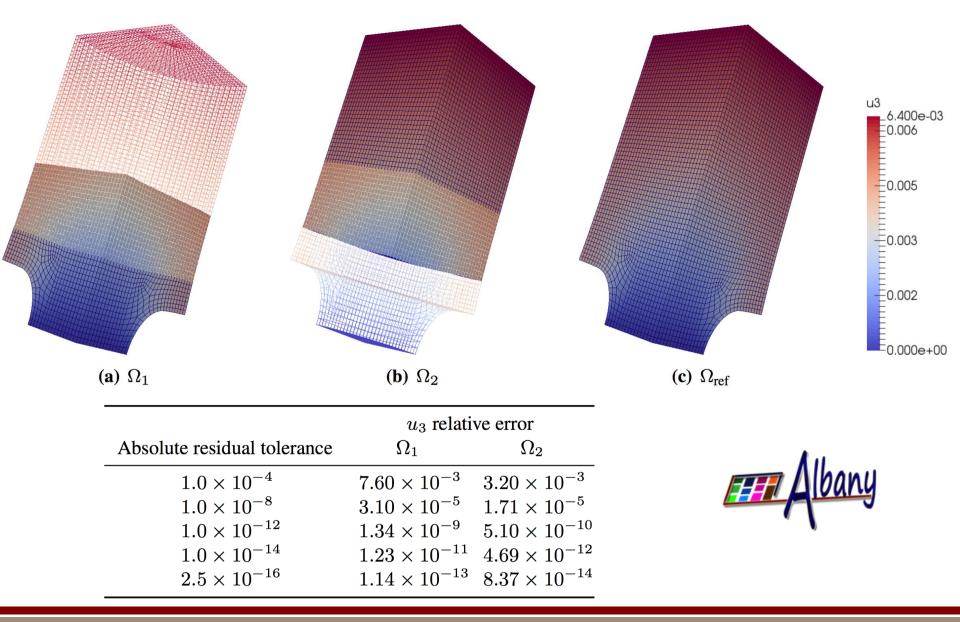
Subdomain	u_3 relative error	σ_{33} relative error	
$\Omega_1 \ \Omega_2$	1.24×10^{-14} 7.30×10^{-15}	$\begin{array}{c} 2.31 \times 10^{-13} \\ 3.06 \times 10^{-13} \end{array}$	



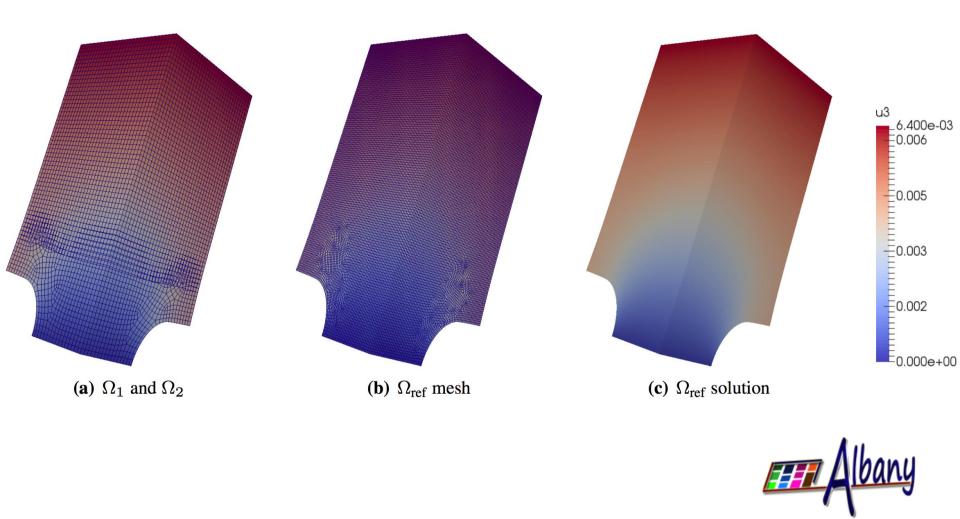


- Notched cylinder that is stretched along its axial direction.
- Domain decomposed into *two subdomains*.
- Neohookean-type material model.

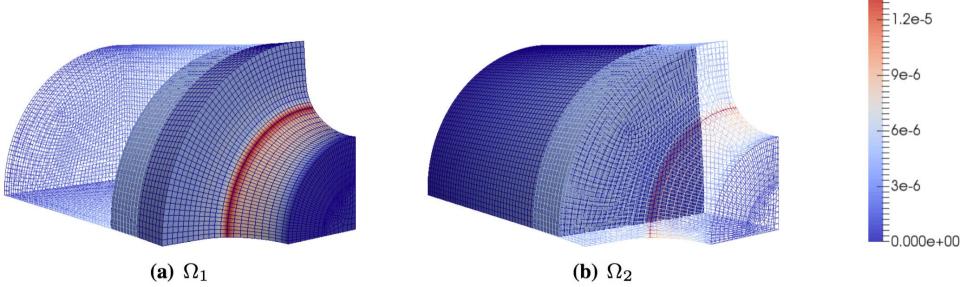
Notched Cylinder: Conformal HEX-HEX Coupling



Notched Cylinder: Nonconformal HEX-HEX in Sandia Coupling



Notched Cylinder: Nonconformal HEX-HEX Sandia Coupling



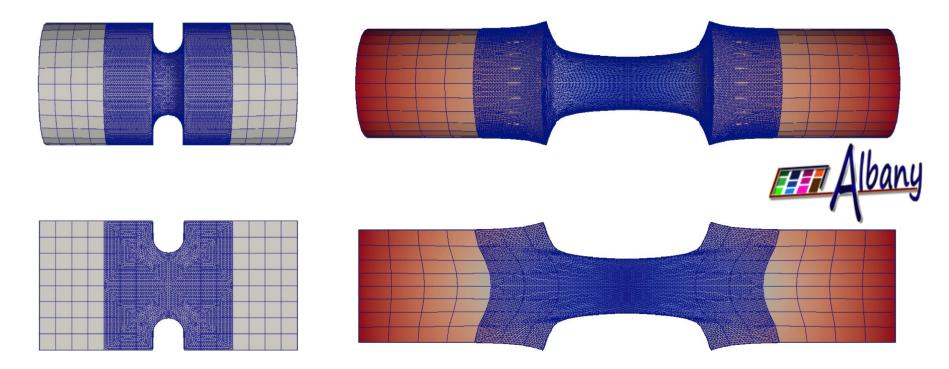
	u_3 relative error	
Absolute residual tolerance	Ω_1	Ω_2
1.0×10^{-8}	1.31×10^{-3}	4.45×10^{-4}
$1.0 imes 10^{-12}$	1.30×10^{-3}	4.43×10^{-4}
1.0×10^{-14}	1.30×10^{-3}	4.43×10^{-4}
2.5×10^{-16}	1.30×10^{-3}	4.43×10^{-4}



Notched Cylinder: TET-HEX Coupling

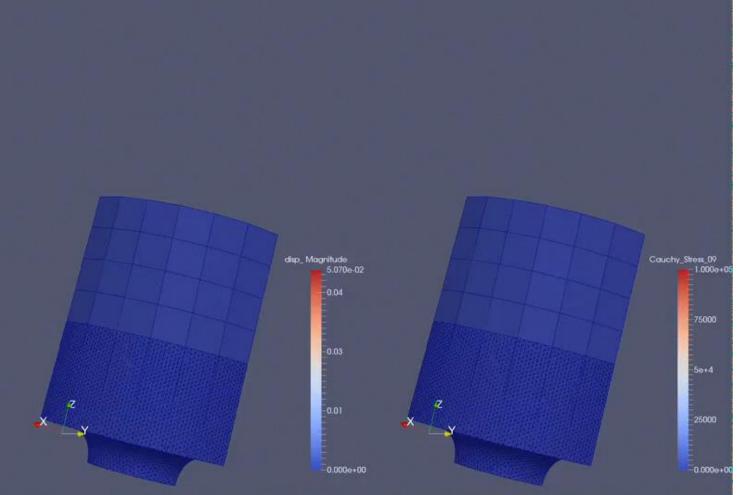


- The Schwarz alternating method is capable of coupling *different mesh topologies*.
- The notched region, where stress concentrations are expected, is *finely* meshed with *tetrahedral* elements.
- The top and bottom regions, presumably of less interest, are meshed with *coarser hexahedral* elements.



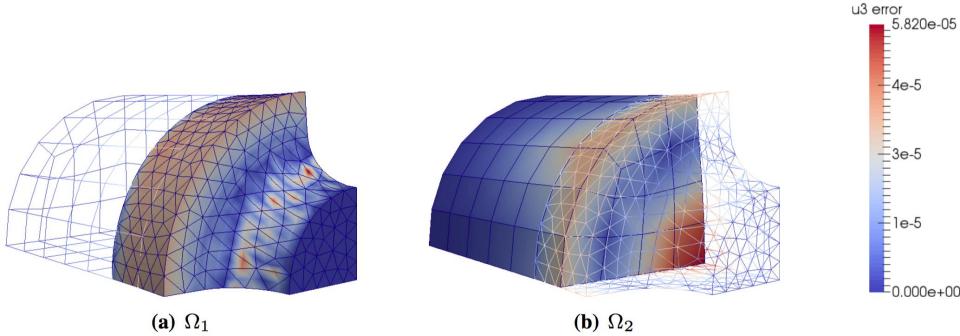
Notched Cylinder: TET-HEX Coupling





Notched Cylinder: Conformal TET-HEX Coupling





	u_3 relative error	
Absolute residual tolerance	Ω_1	Ω_2
$1.0 imes 10^{-14}$	$9.27 imes 10^{-3}$	3.70×10^{-3}



Notched Cylinder: Coupling Different Materials

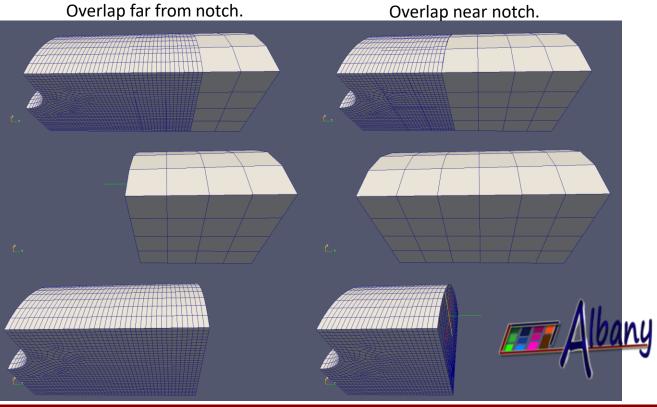
The Schwarz method is capable of coupling regions with *different material models*.

- Notched cylinder subjected to tensile load with an *elastic* and *J2 elasto-plastic* regions.
- *Coarse* region is *elastic* and *fine* region is *elasto-plastic*.
- The overlap region in the first mesh is nearer the notch, where plastic behavior is expected.

Coupled regions

Coarse, elastic region

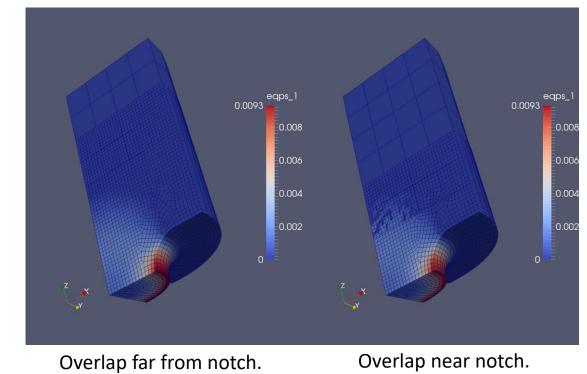
Fine, elasto-plastic region



Notched Cylinder: Coupling Different Materials

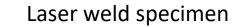
Need to be careful to do domain decomposition so that material models are *consistent* in overlap region.

- When the *overlap* region is *far from the notch*, no plastic deformation exists in it: the coarse and fine regions predict the *same behavior*.
- When the *overlap* region is *near the notch*, plastic deformation spills onto it and the two models predict different behavior, affecting convergence *adversely*.

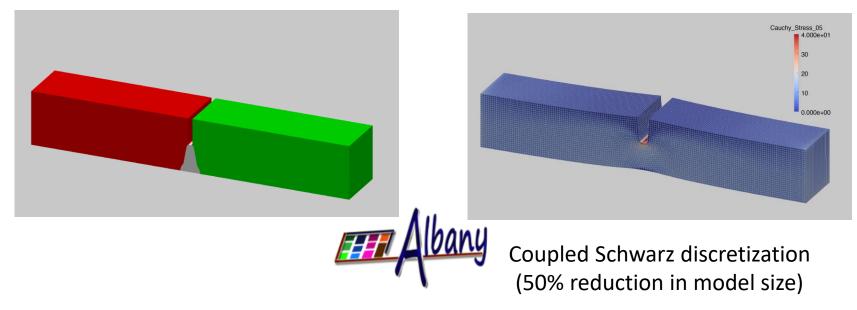




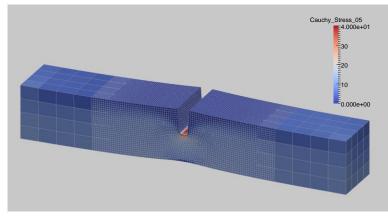
Example #4: Laser Weld with 3 Subdomain 5 Sandia Laboratories



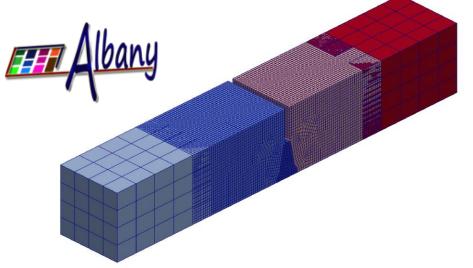
Single domain discretization



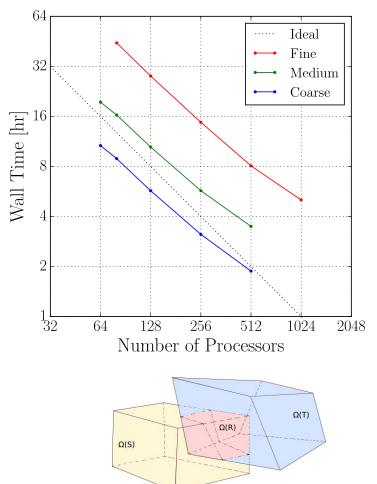
- Problem of *practical scale (~200K dofs).*
- *Isotropic elasticity* and *J2 plasticity* with linear isotropic hardening.
- *Identical parameters* for weld and base materials for proof of concept, to become independent models.



Laser Weld: Strong Scalability of Parallel Schwarz with DTK



• *Near-ideal linear speedup* (64-1024 cores).

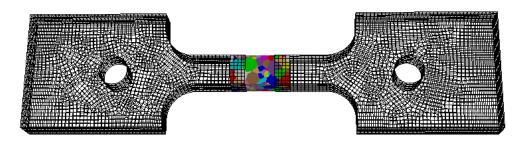


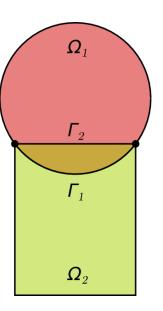
Sandia National

Data Transfer Kit (DTK)

Outline

- 1. Motivation
- 2. Schwarz Alternating Method: Background & History
- 3. Schwarz Alternating Method for Concurrent Multiscale Coupling in Quasistatics
 - Four Variants: Full Schwarz, Inexact Schwarz, Modified Schwarz, Monolithic Schwarz
 - Implementations: MATLAB, Albany
- 4. Numerical Examples
- 5. Summary
- 6. Future Work
- 7. References
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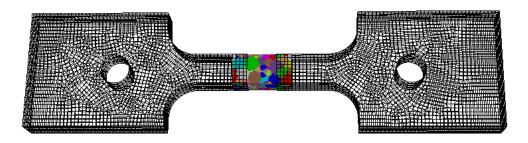
Summary

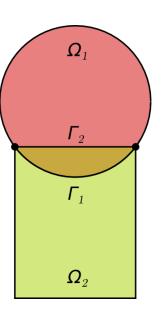


- We have proposed the Schwarz alternating method as a means of concurrent multiscale coupling in finite deformation quasistatic solid mechanics.
- We have developed *four variants* of the Schwarz alternating method (*Full Schwarz, Modified Schwarz, Inexact Schwarz, Monolithic Schwarz*).
- We have *proven* that the Full Schwarz variant converges geometrically for the solid mechanics problem.
- We have *demonstrated numerically* that the *convergence* of the Schwarz method in its four variants is *linear*.
- We have demonstrated *coupling* of *conformal* and *non-conformal meshes*, meshes with *different levels of refinement*, meshes with different *element topologies*, and > *two subdomains* via the proposed method.
- We have demonstrated that the *error* in the coupling can be decreased up to *numerical precision* provided that no other sources of error exist.
- We have developed a *parallel* implementation of the *Modified Schwarz* method in the *Albany code* and demonstrated that the *strong scalability* of our implementation is close to *ideal*.

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Future Work

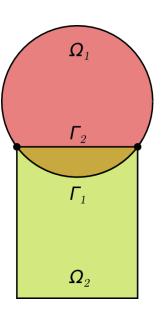


- Extension of the methods presented herein to *transient dynamics (hyperbolic)* problems with the ability to use *different time steps* and *time integrators* for each subdomain.
- Development of a *multi-physics coupling framework* based on variational formulations and the Schwarz alternating method.
- Analysis of the convergence for the other Schwarz variants introduced herein, namely Modified Schwarz, Inexact Schwarz, and Monolithic Schwarz.
- Using the Schwarz alternating method with *different solvers* in different domains.
- Develop a *hybrid FOM-ROM* (full-order-model reduced-order-model) framework using the Schwarz alternating method.
- Introduction of *pervasive multi-threading* into our *Albany* implementation of the Schwarz alternating method using the *Kokkos* framework.
- Multiscale coupling using the proposed Schwarz alternating method in *other applications*.

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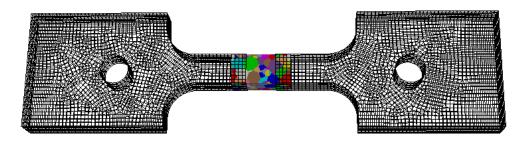
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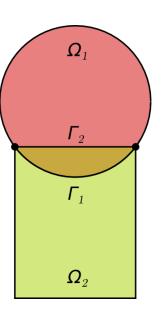
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Appendix. Previous Work

Comput Mech (2014) 54:803-820 DOI 10.1007/s00466-014-1034-0

ORIGINAL PAPER

A multiscale overlapped coupling formulation for large-deformation strain localization

WaiChing Sun · Alejandro Mota

Received: 18 September 2013 / Accepted: 7 April 2014 / Published online: 3 May 2014 © Springer-Verlag Berlin Heidelberg 2014

Abstract We generalize the multiscale overlapped domain framework to couple multiple rate-independent standard dissipative material models in the finite deformation regime across different length scales. We show that a fully coupled multiscale incremental boundary-value problem can be recast as the stationary point that optimizes the partitioned incremental work of a three-field energy functional. We also establish inf-sup tests to examine the numerical stability issues that arise from enforcing weak compatibility in the three-field formulation. We also devise a new block solver for the domain coupling problem and demonstrate the performance of the formulation with one-dimensional numerical examples. These simulations indicate that it is sufficient to introduce a localization limiter in a confined region of interest to regularize the partial differential equation if loss of ellipticity occurs.

strain localization may lead to the eventual failure of materials, this phenomenon is of significant importance to modern engineering applications.

The objective of this work is to introduce concurrent coupling between sub-scale and macro-scale simulations for inelastic materials that are prone to strain localization. Since it is not feasible to conduct sub-scale simulations on macroscopic problems, we use the domain coupling method such that computational resources can be efficiently allocated to regions of interest [14,23,24,30]. To the best of our knowledge, this is the first work focusing on utilizing the domain coupling method to model strain localization in inelastic materials undergoing large deformation.

Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious meshdependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

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Three-field multiscale coupling formulation with compatibility enforced weakly using *Lagrange multipliers*.

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Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious meshdependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

Appendix. Full Schwarz Method

Classical algorithm originally proposed by Schwarz with **outer Schwarz loop** and **inner Newton loop**, each converged to a **tight tolerance** ($\epsilon_{machine}$).

1: $\boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)}$ in $\Omega_{1}, \boldsymbol{x}_{b}^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{1}, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_{1} \triangleright initialize for Ω_1 2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} \triangleright initialize for Ω_2 3: repeat ▷ Schwarz loop $oldsymbol{u}^{(1)} \leftarrow oldsymbol{x}_{\mathrm{D}}^{(1)}$ 4: \triangleright for convergence check $m{x}_eta^{(1)} \leftarrow m{P}_{12} m{x}_B^{(2)} + m{Q}_{12} m{x}_b^{(2)} + m{G}_{12} m{x}_eta^{(2)}$ 5: \triangleright project from Ω_2 to Γ_1 \triangleright Newton loop for Ω_1 6: $\triangle \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)}) \backslash \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)}; \boldsymbol{x}_{b}^{(1)}; \boldsymbol{x}_{\beta}^{(1)})$ 7: \triangleright linear system $oldsymbol{x}_B^{(1)} \leftarrow oldsymbol{x}_B^{(1)} + riangle oldsymbol{x}_B^{(1)}$ 8: until $|| \triangle \boldsymbol{x}_{B}^{(1)} || / || \boldsymbol{x}_{B}^{(1)} || \leq \epsilon_{\text{machine}}$ 9: \triangleright tight tolerance $oldsymbol{y}^{(2)} \leftarrow oldsymbol{x}^{(2)}_B$ 10: \triangleright for convergence check $oldsymbol{x}_eta^{(2)} \leftarrow oldsymbol{P}_{21}oldsymbol{x}_B^{(1)} + oldsymbol{Q}_{21}oldsymbol{x}_b^{(1)} + oldsymbol{G}_{21}oldsymbol{x}_eta^{(1)}$ 11: \triangleright project from Ω_1 to Γ_2 12: \triangleright Newton loop for Ω_2 repeat $\triangle \bm{x}_B^{(2)} \leftarrow -\bm{K}_{AB}^{(2)}(\bm{x}_B^{(2)};\bm{x}_b^{(2)};\bm{x}_\beta^{(2)}) \backslash \bm{R}_A^{(2)}(\bm{x}_B^{(2)};\bm{x}_b^{(2)};\bm{x}_\beta^{(2)})$ 13: \triangleright linear system $oldsymbol{x}_{\mathrm{D}}^{(2)} \leftarrow oldsymbol{x}_{\mathrm{D}}^{(2)} + riangle oldsymbol{x}_{\mathrm{D}}^{(2)}$ 14: until $|| \triangle \boldsymbol{x}_{B}^{(2)} || / || \boldsymbol{x}_{B}^{(2)} || \leq \epsilon_{\text{machine}}$ 15: \triangleright tight tolerance 16: until $\left[\left(||\boldsymbol{y}^{(1)} - \boldsymbol{x}_{B}^{(1)}||/||\boldsymbol{x}_{B}^{(1)}||\right)^{2} + \left(||\boldsymbol{y}^{(2)} - \boldsymbol{x}_{B}^{(2)}||/||\boldsymbol{x}_{B}^{(2)}||\right)^{2}\right]^{1/2} \leq \epsilon_{\text{machine}}$ \triangleright tight tolerance



Appendix. Inexact Schwarz Method

Classical algorithm originally proposed by Schwarz with *outer Schwarz loop* and *inner Newton loop*, with Newton step converged to a *loose tolerance*.



Appendix. Monolithic Schwarz Method



Combines Schwarz and Newton loop into *since Newton-Schwarz loop*, with *elimination of Schwarz boundary DOFs*, and tight convergence tolerance.

$$\begin{array}{ll} 1: \ \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{X}_{B}^{(1)} \text{ in } \Omega_{1}, \ \boldsymbol{x}_{b}^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(1)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_{1}, & \triangleright \text{ initialize for } \Omega_{1} \\ 2: \ \boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)} \text{ in } \Omega_{2}, \ \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)}) \text{ on } \partial_{\boldsymbol{\varphi}} \Omega_{2}, & \triangleright \text{ initialize for } \Omega_{2} \\ 3: \ \mathbf{repeat} & \triangleright \text{ Newton-Schwarz loop} \\ 4: \quad \left\{ \begin{array}{c} \Delta \boldsymbol{x}_{B}^{(1)} \\ \Delta \boldsymbol{x}_{B}^{(2)} \end{array} \right\} \leftarrow \left(\begin{array}{c} \boldsymbol{K}_{AB}^{(1)} + \boldsymbol{K}_{A\beta}^{(1)} \boldsymbol{H}_{11} & \boldsymbol{K}_{A\beta}^{(1)} \boldsymbol{H}_{12} \\ \boldsymbol{K}_{A\beta}^{(2)} \boldsymbol{H}_{21} & \boldsymbol{K}_{AB}^{(2)} + \boldsymbol{K}_{A\beta}^{(2)} \boldsymbol{H}_{22} \end{array} \right) \setminus \left\{ \begin{array}{c} -\boldsymbol{R}_{A}^{(1)} \\ -\boldsymbol{R}_{A}^{(2)} \end{array} \right\} & \triangleright \text{ linear system} \\ 5: \quad \boldsymbol{x}_{B}^{(1)} \leftarrow \boldsymbol{x}_{B}^{(1)} + \Delta \boldsymbol{x}_{B}^{(1)} \\ 6: \quad \boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{x}_{B}^{(2)} + \Delta \boldsymbol{x}_{B}^{(2)} \\ 7: \ \text{ until } \left[\left(||\Delta \boldsymbol{x}_{B}^{(1)}||/||\boldsymbol{x}_{B}^{(1)}|| \right)^{2} + \left(||\Delta \boldsymbol{x}_{B}^{(2)}||/||\boldsymbol{x}_{B}^{(2)}|| \right)^{2} \right]^{1/2} \leq \epsilon_{\text{machine}} \qquad \triangleright \text{ tight tolerance} \end{array}$$

Advantages:

By-passes Schwarz loop.

Disadvantages:

• Off-diagonal coupling terms \rightarrow block linear solver is needed.

Appendix. Modified Schwarz Method



Combines Schwarz and Newton loop into *since Newton-Schwarz loop*, with *Schwarz boundaries* at *Dirichlet boundaries* and tight convergence tolerance.

1: $\boldsymbol{x}_B^{(1)} \leftarrow \boldsymbol{X}_B^{(1)}$ in $\Omega_1, \boldsymbol{x}_b^{(1)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_b^{(1)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_1, \boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{X}_{\beta}^{(1)}$ on Γ_1 \triangleright initialize for Ω	21
2: $\boldsymbol{x}_{B}^{(2)} \leftarrow \boldsymbol{X}_{B}^{(2)}$ in $\Omega_{2}, \boldsymbol{x}_{b}^{(2)} \leftarrow \boldsymbol{\chi}(\boldsymbol{X}_{b}^{(2)})$ on $\partial_{\boldsymbol{\varphi}}\Omega_{2}, \boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{X}_{\beta}^{(2)}$ on Γ_{2} 3: repeat \triangleright Newton-Schwarz loc	~
4: $\boldsymbol{x}_{\beta}^{(1)} \leftarrow \boldsymbol{P}_{12}\boldsymbol{x}_{\beta}^{(2)} + \boldsymbol{Q}_{12}\boldsymbol{x}_{b}^{(2)} + \boldsymbol{G}_{12}\boldsymbol{x}_{\beta}^{(2)}$ \triangleright project from Ω_2 to I	^
5: $\Delta \boldsymbol{x}_{B}^{(1)} \leftarrow -\boldsymbol{K}_{AB}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)}) \setminus \boldsymbol{R}_{A}^{(1)}(\boldsymbol{x}_{B}^{(1)};\boldsymbol{x}_{b}^{(1)};\boldsymbol{x}_{\beta}^{(1)})$ \triangleright linear system in the system of the system	m
6: $oldsymbol{x}_B^{(1)} \leftarrow oldsymbol{x}_B^{(1)} + riangle oldsymbol{x}_B^{(1)}$	
7: $\boldsymbol{x}_{\beta}^{(2)} \leftarrow \boldsymbol{P}_{21}\boldsymbol{x}_{B}^{(1)} + \boldsymbol{Q}_{21}\boldsymbol{x}_{b}^{(1)} + \boldsymbol{G}_{21}\boldsymbol{x}_{\beta}^{(1)}$ \triangleright project from Ω_{1} to I	$\mathbf{\overline{2}}$
8: $\Delta \boldsymbol{x}_{B}^{(2)} \leftarrow -\boldsymbol{K}_{AB}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{\beta}^{(2)}) \setminus \boldsymbol{R}_{A}^{(2)}(\boldsymbol{x}_{B}^{(2)};\boldsymbol{x}_{b}^{(2)};\boldsymbol{x}_{\beta}^{(2)})$ \triangleright linear system in the system of the system	m
9: $oldsymbol{x}_B^{(2)} \leftarrow oldsymbol{x}_B^{(2)} + riangle oldsymbol{x}_B^{(2)}$	
10: until $\left[\left(\triangle \boldsymbol{x}_B^{(1)} / \boldsymbol{x}_B^{(1)} \right)^2 + \left(\triangle \boldsymbol{x}_B^{(2)} / \boldsymbol{x}_B^{(2)} \right)^2 \right]^{1/2} \le \epsilon_{\text{machine}}$ \triangleright tight tolerand	ce

Advantages:

- By-passes Schwarz loop.
- No diagonal coupling (conventional linear solver can be used in each subdomain).

Least-intrusive variant: by-passes Schwarz iteration, no need for block solver.

Appendix. Convergence Proof



A. Moto, I. Tezaut, C. Alleman Schwarz Alternating Method in Solid Mechanics

2 Formulation of the Schwarz Alternating Method

We start by defining the standard finite deformation variational formulation to establish notation before presenting the formulation of the coupling method.

2.1 Variational Formulation on a Single Domain

2.1 Variations formations on a single to omain factor is a single to make the strength of the strength o

 $\Phi[\varphi] := \int_{\Omega} A(\mathbf{F}, \mathbf{Z}) dV - \int_{\Omega} R\mathbf{B} \cdot \varphi dV - \int_{\Omega \cap \Omega} \mathbf{T} \cdot \varphi dS$ (1)

in which A(F, Z) is the Helmholtz free-mergy density and Z is a collection of internal variables. The weak form of the problem is obtained by minimizing the energy functional $\Phi(a)$ over the Soboter space $W_1^1(0)$ that is comprised on all have square integrable in the averagement first derivatives. Define $S := \{\varphi \in W_2^1(\Omega) : \varphi = \chi \text{ on } \partial_\varphi \Omega\}$ (2)

and $\mathcal{V} := \{ \xi \in W_2^1(\Omega) : \xi = 0 \text{ on } \partial_{\varphi} \Omega \}$

(3) ν := γξ ∈ W₂(x), ζ = 0 on φ₂(x) = ζ = 0 on φ₂(x) = ζ = φ₁(φ₂) = φ₁(φ + cξ) for all quarks where ξ ∈ V is a test function. The potential energy is minimized if and only if Φ₁(φ) ≤ Φ₁(φ + cξ) for all ξ ∈ V and c ∈ R. It is straightforward to show that the minimum of Φ₁(φ) is the mapping φ ∈ S that satisfies

 $D\Phi[\varphi](\xi) = \int_{\Omega} \mathbf{P} : \text{Grad} \, \xi \, dV - \int_{\Omega} R\mathbf{B} \cdot \xi \, dV - \int_{\partial \gamma \cdot \Omega} \mathbf{T} \cdot \xi \, dS = 0,$ (4)

where $P = \partial A/\partial F$ denotes the first Piola Kirchhoff stress. The Euler-Lagrange equation corresponding to the variational statement (4) is Div P + RB = 0, in Ω ,

PN = T, on $\partial_T \Omega$, $\varphi = \chi$, on $\partial_{\varphi} \Omega$. (5)

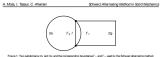
2.2 Coupling Two or More Subdomains via the Schwarz Alternating Method In this section, we describe the Schwarz advantation method for coupling multiple workpropus advantation of the schwarz advantation method advantation of the schwarz advantation method with the schwarz advantation of the schwarz advantation method with the schwarz advantation method method with the schwarz advantation of the schwarz advantati

 $n \in \mathbb{N}^{0} = \{0, 1, 2, ...\}, \quad i = 2 - n + 2 \lfloor \frac{n}{2} \rfloor \in \{1, 2\}, \quad j = n + 1 - 2 \lfloor \frac{n}{2} \rfloor \in \{1, 2\},$ (6) 5

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A Proof of Convergence of th Finite-Deformation Inelasti		dethod for the
In this section we give a proof of Theorem 1. remarks. Assume properties 1–5 enumerated in		es, presented below as
Remark 1 By the coercivity of $\Phi[\varphi]$, it follows this functional over S exists, i.e., the minimization		a unique minimizer to
Remark 2 By the Stampacchia theorem, the mithat		
$\langle \Psi _{S}$ for all $\xi \in S$.	φ], $\xi - \varphi$) ≥ 0	(51)
Remark 3 Recall that the strict convexity prop		
$\Phi[\psi_2] - \Phi[\psi_1]$	$-(\Phi'[\psi_1], \psi_2 - \psi_1) \ge 0,$	(52)
$\forall \psi_1, \psi_2 \in S$. From (36), remark that if $\Phi[\varphi]$ is st an $\alpha_R > 0$ such that $\forall \psi_1, \psi_2 \in K_R$ we have	rictly convex over $S \forall R \in \mathbb{R}$ such the	at $R < \infty$, we can find
$\Phi[\psi_2] - \Phi[\psi_1] - (\Phi'[\psi$	$ \psi_1 , \psi_2 - \psi_1 \ge \sigma_R \psi_2 - \psi_1 ^2.$	(53)
Remark 4 By property 5, the uniform continuit	ty of $\Phi'[\varphi]$, there exists a modulus of	continuity $\omega > 0$, with
$\omega : K_R \rightarrow K_R$, such that $ \Phi'(\psi_1) - \Phi' $	$ \psi_2 \le \omega(\psi_1 - \psi_2),$	(54)
$\forall \psi_1, \psi_2 \in K_R$. By definition, $\omega(\epsilon) \rightarrow 0$ as $\epsilon \rightarrow 0$		
Remark 5 It was shown in [35] that in the cas	ie $\Omega_1 \cap \Omega_2 \neq \emptyset, \forall \varphi \in S$, there exist	$\zeta_1 \in S_1$ and $\zeta_2 \in S_2$
	$= \zeta_1 + \zeta_2$,	(55)
and $\max(\zeta_1 $	$, \zeta_2) \le C_0 \varphi ,$	(56)
for some $C_0 > 0$ independent of φ .		
Remark 6 Note that (39) can be written as		
$(\Phi'[\tilde{\varphi}^{(n)}], \xi^{(i)}) = 0,$	for $\tilde{\varphi}^{(u)} \in \tilde{S}_n$, $\forall \xi^{(i)} \in S_i$,	(57)
for $i \in \{1, 2\}$ and $n \in \{0, 1, 2,\}$ (recall from (6 of the solution to each minimization problem over \hat{S}_n .) the relation between i and n). This is \hat{S}_n and the definition of $\hat{\varphi}^{(n)}$ as the n	s due to the uniqueness ninimizer of $\Phi[\varphi]$ over
Remark 7 Let $\hat{\varphi}^{(n)} \in \hat{S}_n$, and let $\xi \in S$. By R	temark 5, there exist $\zeta_1 \in S_1$ and ζ_2	$\in S_2$ such that

 $(\Phi'[\hat{\omega}^{(n)}], \ell) = (\Phi'[\hat{\omega}^{(n)}], \ell_1 + \ell_2),$ (58)

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the list i + rad (- 2 lin / i cost, and (- 2 ma) (- 1 lin / i cost), threadown the following definitions for induced as the doctory of th			
$\label{eq:second} \begin{array}{l} \bullet \mbox{ Order}_{k} \in \mathfrak{B}(\mathbf{y}) \in \mathfrak{B}(\mathbf{x}), \ \mbox{ Bin}\\ \bullet \mbox{ Nameral module}_{k}, \ \mbox{ Bin} \in \mathfrak{B}(\mathbf{x}), \ \mbox{ Bin}, \ Bi$		r	
$\label{eq:second} \begin{split} & (\operatorname{second}_{\mathcal{C}}(\operatorname{second}_{\mathcal{C}}), & (\operatorname{second}_{\mathcal{C}}), & ($	 Closure: (0) >= (0) / (0) 		
$\label{eq:second} \begin{array}{l} \cdot \operatorname{Brane transmission} (f_i = 0; i) \\ Rescalar bar with the excession \\ \operatorname{Rescalar bar with the excession \\ \operatorname{$	Dirichlet boundary: @ III) := @ IIII III).		
Note the three detectors we guarantee that $@(z) = 0$, $@(z) = 1$, $@(z) = 1$, $(z) = 1$	 Neumann boundary: @ ID := @ ID I ID. 		
Now define the species $ \begin{array}{l} (\boldsymbol{g}_{n} \in I \times \boldsymbol{g}_{n}^{(1)}(\boldsymbol{g}_{n})) := \boldsymbol{g}_{n} \in \boldsymbol{g}_{n} \in \boldsymbol{g}_{n}^{(1)}, \ [\boldsymbol{f}_{n}^{(1)} \in \boldsymbol{g}_{n}^{(1)} \in \boldsymbol{g}_{n}^{(1)} \in \boldsymbol{g}_{n}^{(1)}, \ [\boldsymbol{f}_{n}^{(1)} \in \boldsymbol{g}_{n}^{(1)} $	 Schwarz boundary: Γ_i := @Q \ IQ. 		
and $V_{i} \approx \left\{ e \in W_{i}^{i}(\alpha) : i = 0 \Rightarrow 0 \oplus 2\pi/f_{i}$ (0) where the quester P_{ij} , e_{ij}^{i} ($\frac{1}{2}$ divers the projection for the absolution By constrained for the field boundary f_{i} . The projection is operating the spectra of the field boundary diverse the straining emperation of the strained operation o	Note that with these definitions we guarantee that $\otimes \boxtimes 1 \otimes \boxtimes 2 = ;$, $\otimes \boxtimes 1 \Gamma_i = ;$ and $\otimes \boxtimes 1 \Gamma_i = ;$ Now define the spaces		
$ \begin{array}{c} V_{1} = \left\{ -2 \; \delta_{1}^{2}(0) : -e^{2} \circ 0 \in 0 \; 0 \; (1 \; , \ (0) \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\$	$S_i := \{ 2 W_2^i(\mathbb{Z}) : 2 = \chi \text{ on } \otimes \mathbb{Z}_{2^{-1}} = P_{\mathbb{Z}_2^{-1}} : \{ (\mathbb{Z}_2^i) \} \text{ on } \Gamma_i \}$, (7)	,	
where the querter $P_{q_{1}}$, (1) draws the projection from the absorburs fly on the Ghraux boundary f_{1} . The projection querter prices are strategistically as a strategistic of the Shraux boundary of the distribution	and		
The pspeciation spacing rules a contrast relation the Schwarz alternating relation. In from an entraphysication of the state of the Schwarz alternating relation, the Schwarz alternation of the Schwarz alternat	$V_i := \{ \leftarrow 2 \ W_2^1(X) \} : \leftarrow = 0 \text{ on } @ X \setminus [\Gamma_i],$ (8)	,	
where id_X is the identity map that maps X onto itself (i.e. zero displacement), and Z Z	This projection operator plays a contral role in the Schwarz alternating method. Its form and implementation are discussed in subsequent sections. For the moment It is sufficient to assume that the operator is able to project a field ⁴ . The mon esubdomain to the Schwarz boundary of the other subdomain. The Schwarz alternating method solves a sequence of problems on Sig- and Sig. The solution ⁴ (4) for the add the other in the sub-		
	where id_X is the identity map that maps X onto itself (i.e. zero displacement), and Z Z Z Z)	

A better guess, if available, may be used to initialize $'^{(0)}$ on \boxtimes rather than the identity map id_X. The minimization of the functional (10) leads to a variational formulation of the form (4)-(5) for each subdomain 7 $D\Phi_i(\tau^{(n)})(e^{(j)}) = \bigcup_{i=0}^{L} P: \text{Grad} e^{(j)} dV - \bigcup_{i=0}^{L} RB \cdot e^{(j)} dV - \bigcup_{i=0}^{L} T \cdot e^{(j)} dS = 0,$ (11)

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By (57), $(\Phi'[\hat{\omega}^{(n-1)}], \hat{c}_0) = 0$. Hence,		
1		
$(\Phi'[\phi^{(n)}], \zeta_1 + \zeta_2) = (\Phi'[\phi^{(n)}], \zeta_1 + \zeta_2) - (\Phi'[\phi^{(n-1)}], \zeta_2 + \zeta_2) - (\Phi'[\phi^{(n-1)}], \zeta_$		
since $(\Phi^{i}[\hat{\varphi}^{(n)}], \zeta_{1}) = 0$, also by (57). By the Cauchy-Schwarz inequality,		
$(\Phi'[\varphi^{(n)}], \zeta_2) - (\Phi'[\varphi^{(n-1)}], \zeta_2) = (\Phi'[\varphi^{(n)}] - \Phi'[\varphi^{(n-1)}], \zeta_2)$	$ \Phi^{(1)} , \zeta_2) \le \Phi'[\phi^{(n)}] - \Phi'[\phi^{(n-1)}] \cdot \zeta_2 .$ (60)	
Again using (57) and also (58) in (60) leads to		
$(\Phi'[\tilde{\varphi}^{(n)}] - \Phi'[\tilde{\varphi}^{(n-1)}], \zeta_2) = (\Phi'[\tilde{\varphi}^{(n)}], \xi)$	$\leq \Phi'[\tilde{\varphi}^{(n)}] - \Phi'[\tilde{\varphi}^{(n-1)}] \cdot \zeta_2 ,$ (61)	
and substituting (56) into (61) we finally obtain that		
$\langle \Psi'[\hat{\varphi}^{(n)}], \xi \rangle \le C_0 \Psi'[\hat{\varphi}^{(n)}] $	$-\Phi'[\hat{\varphi}^{(n-1)}] \cdot \xi ,$ (62)	
$\forall \xi \in S.$		
Remark 8 For part (d) of Theorem 1, recall the definition	n of geometric convergence:	
$E_{n+1} \le CI$	S ₀ , (63)	
$\forall n \in \{0, 1, 2,\}$ for some $C > 0$, where		
$E_n := \phi^{(n+1)} $	$-\phi^{(n)} $. (64)	
Remark 9. Recall from the definition of continuity that if $\Phi^*(\varphi)$ is Lipshitz continuous at $\tilde{\varphi}^{(\alpha)}$ near φ , then there exists a constant $K \ge 0$ such that		
$\frac{ \Phi'[\phi^{(\alpha)}] - \Phi'[\phi}{ \phi^{(\alpha)} - \phi }$	$\frac{111}{100} \le K.$ (65)	
Considering that $\Phi'[\phi]=0$ since ϕ is the minimizer of $\Phi[g$	o], (65) is equivalent to	
$ \Phi'[\phi^{(n)}] \le K \phi $	$ y^{(n)} - \varphi $. (66)	
Proof of Theorem 1		
Proof of (a). Let $\overline{\psi}^{(1)} = \arg \min_{\psi \in S_0^+} \Phi[\psi]$. By (40), $\overline{\psi}^{(1)} \in \overline{S_2}$. Let $\overline{\psi}^*$ be the minimizer of $\Phi[\varphi]$ service $\overline{S_2}$ and suppose $\Phi[\psi^*] > \Phi[\psi^{(1)}]$. But this is a contradiction, since we can take $\overline{\psi}^* = \overline{\psi}^{(1)}$. Hence, it cannot be that $\Phi[\overline{\psi}^{(1)}] = \Phi[\overline{\psi}^{(1)}] = \arg \min_{\psi \in S_2^+} \Phi[\psi]$. It follows by induction that		
$\Phi[\hat{\varphi}^{(n)}] \le \Phi[$	$\hat{\rho}^{(n-1)}$] (67)	
for $n \in \{1, 2, 3,\}$. Now let φ be the minimizer of Φ unique. Hence $\Phi[\varphi] \leq \Phi[\varphi^{(n)}]$ for all $n \in \{1, 2, 3,\}$.	$ \varphi $ over S. Since the problem is well-posed φ is	

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$ \left\{ \begin{array}{c} \left\{ \begin{array}{c} 1 \\ 1 \\ - \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \\ - \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \\ - \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \right\} \left\{ \begin{array}{c} 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \right\} \left\{ 1 \end{array} \right\{ 1 \\ \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \left\{ 1 \end{array} \left$	to linear syste
$ \begin{array}{l} & \left\{ \begin{array}{l} \Delta x_{0}^{(1)} \\ \Delta x_{0}^{(2)} \end{array} \right\} \leftarrow \left(\begin{array}{l} \mathbf{K}_{0d}^{(1)} + \mathbf{K}_{0d}^{(1)} \mathbf{H}_{11} \\ \mathbf{K}_{0d}^{(2)} \mathbf{H}_{11} \\ \mathbf{K}_{0d}^{(2)} \mathbf{K}_{0d}^{(2)} + \mathbf{K}_{0d}^{(1)} \mathbf{H}_{2d} \\ \end{array} \right) \setminus \left\{ \begin{array}{l} -\mathbf{R}_{0}^{(1)} \\ -\mathbf{R}_{0}^{(1)} \\ -\mathbf{R}_{0}^{(1)} \\ -\mathbf{R}_{0}^{(1)} \end{array} \right\} \\ & s \mathbf{x}_{0}^{(1)} \leftarrow \mathbf{x}_{0}^{(1)} + \Delta \mathbf{x}_{0}^{(1)} \\ \mathbf{K}_{0d}^{(2)} \mathbf{H}_{2d} \\ \mathbf{K}_{0d}^{(2)} + \mathbf{K}_{0d}^{(2)} \mathbf{H}_{2d} \\ -\mathbf{K}_{0d}^{(2)} \\ -\mathbf{K}_{0d}^{(2)} \end{array} \right) \\ & s \mathbf{K}_{0d}^{(2)} \mathbf{K}_{0d}^{(2)} \mathbf{H}_{2d} \\ \end{array} $	
$\begin{array}{c} & & \sum_{m} (1 - x_{B}^{(1)} + i - i - x_{B}^{(1)} \\ & & x_{B}^{(1)} + e x_{B}^{(2)} + i - x_{B}^{(1)} \\ & & \text{undif} \left[\left((\Delta x_{B}^{(1)}) / x_{B}^{(1)} \right)^{2} + \left(\Delta x_{B}^{(1)} / x_{B}^{(2)} \right)^{2} \right]^{1/2} \leq \epsilon_{\text{matrix}} \end{array}$	

[35, 34, 4]. Although we do not provide here formal convergence proofs for the remaining variants of the Schwarz method, we offer some manerical results illustrating their convergence in Section 4. Consider the energy functional $\Phi[\phi]$ defined in (1). We will denote by $\langle \cdot, \cdot \rangle$ the usual L^2 inner product m = 0 then is the second section of the se over Ω , that is, $(\psi_1, \psi_2) := \int_{U} \psi_1 \cdot \psi_2 \, dV,$ (35) for $\psi_1, \psi_2 \in W_2^1(\Omega)$, with corresponding nerm $\|\cdot\|_1$ proof of the convergence of the Schwarz alternating method requires that the functional $\Phi[\varphi]$ satisfy the following properties over the space S defined in (2): 1. $\Phi[\varphi]$ is coercive. 2. $\Phi[\varphi]$ is Fréchet differentiable, with $\Psi[\varphi]$ denoting its Fréchet derivative. Φ[φ] is strictly convex. 4. $\Phi[\varphi]$ is lower semi-contin 5. $\Phi'[\phi]$ is uniformly continuous on $\mathcal{K}_{\mathrm{R}},$ where $\mathcal{K}_R := \{ \varphi \in S : \Phi[\varphi] < R, R \in \mathbb{R}, R < \infty \}$. (36) It can be shown that the energy functional $\vartheta(\varphi)$ defined in (1) is strictly convex in S (property 3) provided that the Helmholtz free-energy density $\mathcal{A}(F, Z)$ is a quasi-convex function of F [26]. Properties 1, 2, 4 and 5 follow from the strict convexity of $\vartheta(\varphi)$. Next, then we conditional sets of spaces $\hat{S}_{n} := \left\{ \varphi \in S : \varphi = P_{\Omega_{j} \rightarrow \Gamma_{1}} [\varphi^{(n-1)}(\Omega_{j})] \text{ on } \Gamma_{i}, \varphi = \varphi^{(n-1)} \text{ on } \Omega \setminus \Omega_{i} \right\},$ (37) and $\tilde{V}_i := \{ \xi \in S : \xi = 0 \text{ in } \Omega \setminus \Omega_i \}$, (38)

where i = 1 and j = 2 if n is odd, and i = 2 and j = 1 if n is even for $n \in \{1, 2, ...\}$ as given by (6) and with the function $q^{(0)} \in S$ an initial gauss. Note that the spaces S_i in (37) are extensions of the spaces S_i in (7) to the entities domain 10. With this touttain in place, the solution start place $n \in [0, 1]$ are recarst one of the recarst S_i in $(n \in [0, 1])$. $\bar{\varphi}^{(n)} = \begin{cases} \mathrm{id}_{\boldsymbol{X}}, & \text{for } n = 0; \\ \arg\min_{\boldsymbol{w} \in \mathcal{N}} \Phi[\varphi], & \text{for } n > 0. \end{cases}$ (39)

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Proof of (b). By (a) $\Phi[Q^{(n)}] \rightarrow l$ as $n \rightarrow \infty$ for some	
Private of (b). By (a) $\Phi[\varphi^{n+1}] \rightarrow I$ as $n \rightarrow \infty$ for some bound	1 () F. Now, combining (51) and (53), we have the
$\Phi[\hat{\varphi}^{(n)}] - \Phi[\hat{\varphi}^{(n+1)}] \ge \Phi[\hat{\varphi}^{(n)}] - \Phi[\hat{\varphi}^{(n+1)}] - \left(\Phi'[$	(68)
for all $n \in \{1, 2, 3,\}$. Since $\Phi[\tilde{\varphi}^{(n)}] \rightarrow l$ as $n \rightarrow \infty$ From (68), we have that $\lim_{n \rightarrow \infty} \tilde{\varphi}^{(n)} - \tilde{\varphi} $	
from which we can conclude that $\hat{\varphi}^{(n)} - \hat{\varphi}^{(n+1)} \rightarrow 0$	
$ \varphi - \tilde{\varphi}^{(u)} ^2 \le \frac{1}{\alpha_R} \left\{ \Phi(\varphi) - \Phi \tilde{\varphi} - \Phi \tilde{\varphi} \right\}$	$[\hat{\varphi}^{(n)}] = \left(\Phi'[\hat{\varphi}^{(n)}], \varphi - \hat{\varphi}^{(n)}\right)\right).$ (70)
Since φ is the minimum of $\Phi[\varphi]$, by (a) we have that	$\Phi[\varphi] \le \Phi[\tilde{\varphi}^{(n)}]$. It follows that
$\Phi[\varphi] - \Phi[\hat{\varphi}^{(n)}] - \left(\Phi'[\hat{\varphi}^{(n)}], \varphi - \hat{\varphi}^{(n)}\right) \le - \left(\Phi'$	$[\hat{\varphi}^{(n)}], \varphi - \hat{\varphi}^{(n)}) = (\Phi'[\hat{\varphi}^{(n)}], \hat{\varphi}^{(n)} - \varphi).$ (71)
Substituting (71) into (70) we have	
$ \varphi - \bar{\varphi}^{(n)} ^2 \le \frac{1}{\alpha \nu} \left(\epsilon$	$V[\hat{\varphi}^{(n)}], \hat{\varphi}^{(n)} - \varphi$. (72)
Now by (62) (Remark 7),	
$\left(\Phi'[\hat{\varphi}^{(n)}], \hat{\varphi}^{(n)} - \varphi\right) \le C_0 \Phi'[\hat{\varphi}] $	$^{(n)}] = \Phi'[\tilde{\varphi}^{(n-1)}] \cdot \tilde{\varphi}^{(n)} - \varphi .$ (73)
Substituting (73) into (72) leads to	
$ \hat{\varphi}^{(n)} - \varphi \le \frac{C_0}{\alpha_R} \Phi' $	$[\hat{\varphi}^{(n)}] = \Phi'[\hat{\varphi}^{(n-1)}] .$ (74)
Applying the uniform continuity assumption (54), we of	obsain
$ \hat{\varphi}^{(n)} - \varphi \le \frac{C_0}{\alpha_R} \omega$	$ \tilde{\varphi}^{(n)} - \tilde{\varphi}^{(n-1)} $. (75)
By (69), $ \hat{\varphi}^{(n)} - \hat{\varphi}^{(n-1)} \rightarrow 0$ as $n \rightarrow \infty$. From t $n \rightarrow \infty$.	his we obtain the result, namely that $\hat{\varphi}^{(o)} \rightarrow \varphi$ as
Proof of (c). This follows immediately from (a) and (b).
Proof of (d) . By (b) , for large enough n , there exists s	ome $C_1 > 0$ independent of n such that
$ \hat{\varphi}^{(n)} - \varphi ^2 \le C_1 $	$ \hat{\varphi}^{(n+1)} - \hat{\varphi}^{(n)} ^2$. (76)
Let us choose C_1 such that $C_1 > \alpha_R/K$, where K is (68) with (76) leads to	the Lipshitz continuity constant in (66). Combining
$\frac{1}{\alpha_R} \left(\Phi[\hat{\varphi}^{(n)}] - \Phi[\hat{\varphi}^{(n+1)}] \right) \ge \hat{\varphi}^{(n+1)} $	$ \psi^{(n)} - \hat{\varphi}^{(n)} ^2 \ge \frac{1}{C_1} \hat{\varphi}^{(n)} - \varphi ^2.$ (17)

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Remark that [50] $\hat{\mathcal{S}}_n = \hat{\varphi}^{(n-1)} + \hat{\mathcal{V}}_i \quad \text{for} \quad \hat{\varphi}^{(n-1)} \in \hat{\mathcal{S}}_{n-1} \Rightarrow \hat{\varphi}^{(n-1)} \in \hat{\mathcal{S}}_n.$ (40) **Theorem 1.** Assume that the energy functional $\Phi(\varphi)$ satisfies properties l-5 above. Consider the Schwarz alternating method of Section 2 defined by (9)–(13) and its equivalent form (39). Then $(a) \ \Phi[\hat{\varphi}^{(0)}] \geq \Phi[\hat{\varphi}^{(1)}] \geq \cdots \geq \Phi[\hat{\varphi}^{(n-1)}] \geq \Phi[\hat{\varphi}^{(n)}] \geq \cdots \geq \Phi[\varphi], \text{ where } \varphi \text{ is the minimizer of } \Phi[\varphi]$ (b) the sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in S.

(c) the Schwarz minimum values $\Phi[\vec{\varphi}^{(n)}]$ converge monotonically to the minimum value $\Phi[\varphi]$ in S starting from any initial guess $\vec{\varphi}^{(0)}$.

(d) if Φ[†](φ) is Lipschit; continuous in a neighborhood of φ, then the sequence {φ⁽ⁿ⁾} converges geometrically to the minimizer φ². Proof. See Appendix A. п

Finally, while most of works cited above present their analysis for the specific case of two subdomains extension to multiple subdomains is in general straightforward. The case of multiple subdomains is considered specifically in Liess [53]. Batter [4], and Li-Shan and Faras [34].

4 Numerical Examples

Interactive Learning to the second second

4.1 Implementation

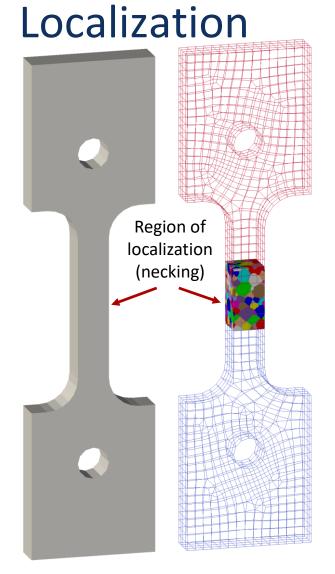
The four variants of the Schwarz alternating method described in Section 2.4 have been implemented in a non-dimensional MATLAN code. The objective is to determine the convergence behavior, efficiency, and performance of each variant. This code has been optimized both in terms of memory usage and essecution record.

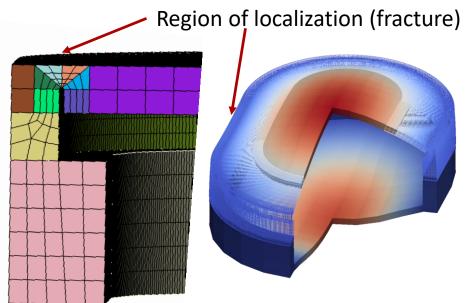
eed. In addition, the Modified variant of the Schwarz alternating method described in Section 2.4 has been uplemented in ALBANY, an open-source multiphysics research platform developed mainly at Sandis National ³See Remark 9 in the Appendix for a definition of geometric corresponce. 15

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Using the identity $\Phi[\hat{\varphi}^{(n)}] - \Phi[\hat{\varphi}^{(n+1)}] = \left(\Phi[\hat{\varphi}^{(n)}] - 0\right)$ that		
$\left(\Phi[\varphi^{(n)}] - \Phi[\varphi]\right) - \left(\Phi[\varphi^{(n+1)}]\right)$	$-\Phi[\varphi] \ge \frac{\alpha_R}{C_1} \varphi^{(n)} - \varphi ^2.$ (78)	
Substituting $\psi_1 = \hat{\varphi}^{(n)}$ and $\psi_2 = \hat{\varphi}$ into (53) and rear	ranging, we obtain	
$\left(\Phi'[\hat{\varphi}^{(n)}], \varphi - \hat{\varphi}^{(n)}\right) \le \left(\Phi'[\hat{\varphi}^{(n)}], \varphi - \hat{\varphi}^{(n)}\right)$		
since $\alpha_B \ge 0$. Now, by the Cauchy-Schwarz inequality followed by the application of the Lipshitz continuity of $\Phi'[\varphi]$ (66) we can write		
$\left(\Phi' \tilde{\varphi}^{(n)} , \varphi - \tilde{\varphi}^{(n)}\right) \le \Phi' \tilde{\varphi}^{(n)} $	$ \varphi - \tilde{\varphi}^{(n)} \le K \varphi - \tilde{\varphi}^{(n)} ^{2}$. (80)	
Hence, from (79), $\Phi[\phi^{(u)}] - \Phi[\phi] \leq$	$K \dot{\varphi}^{(n)} - \varphi ^2$. (81)	
Moreover, by (53) since $\Phi'[\phi] = 0$,		
$\Phi[\hat{\varphi}^{(n)}] - \Phi[\varphi] \ge \epsilon$	$x_R \hat{\varphi}^{(u)} - \varphi ^2$. (82)	
Using (81) and (82) we obtain		
$\left(\Phi[\hat{\varphi}^{(n)}] - \Phi[\varphi]\right) - \left(\Phi[\hat{\varphi}^{(n+1)}] - \Phi[\varphi]\right)$	$\leq K \tilde{\varphi}^{(u)} - \varphi ^2 - \alpha_R \tilde{\varphi}^{(n+1)} - \varphi ^2.$ (83)	
Combining (83) and (78) leads to		
$\frac{\alpha_R}{C_1} \boldsymbol{\varphi}^{(n)} - \boldsymbol{\varphi} ^2 \leq \left(\Phi[\boldsymbol{\varphi}^{(n)}] - \Phi[\boldsymbol{\varphi}] \right) - \left(\Phi[\boldsymbol{\varphi}^{(n+1)}] \right)$	$-\Phi[\varphi] \le K \phi^{(n)} - \varphi ^2 - \alpha_R \phi^{(n+1)} - \varphi ^2.$ (84)	
or $ \hat{\varphi}^{(n+1)} - \varphi \le$	$B \hat{\varphi}^{(n)} - \varphi $ (85)	
with $B := \sqrt{\frac{K}{\alpha_I}}$	$\frac{1}{c} - \frac{1}{C_1}$, (86)	
and $B \in \mathbb{R}$ as we chose $C_1 > \alpha_R/K$. Furthermore, sin the minimizer φ of $\Phi[\varphi]$ by (b) and (c), it follows that can be recast as		
$ \phi^{(\alpha+1)} - \phi^{(\alpha)} \le 0$	$2 \phi^{(n)} - \phi^{(n-1)} $ (87)	
whereapon the claim is proven.	0	
B Analytic Solution for Linear-Elastic Singular Bar		
As reference, herein we provide the solution of the si equilibrium equation is	ngular bar of Section 4.3 for linear elasticity. The	
$P=\sigma(X)A(X)={\rm const.}, \sigma(X)=E\epsilon(X),$	$e(X) := u'(X), A(X) = A_0 \left(\frac{X}{L}\right)^{\frac{1}{2}},$ (88)	
37		

Appendix. Multiscale Modeling of





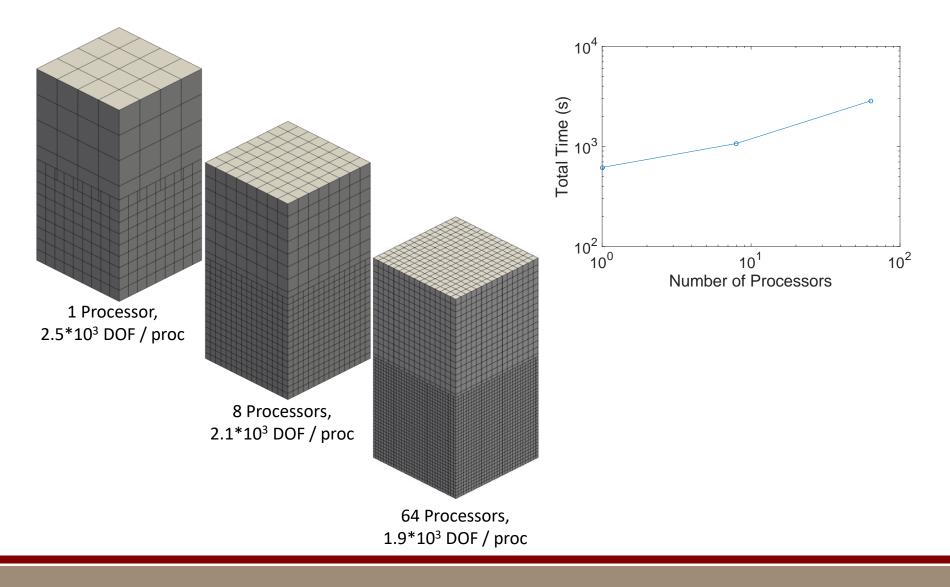


Strain localization can cause *localized necking* (left) and ultimately *fracture* (above).

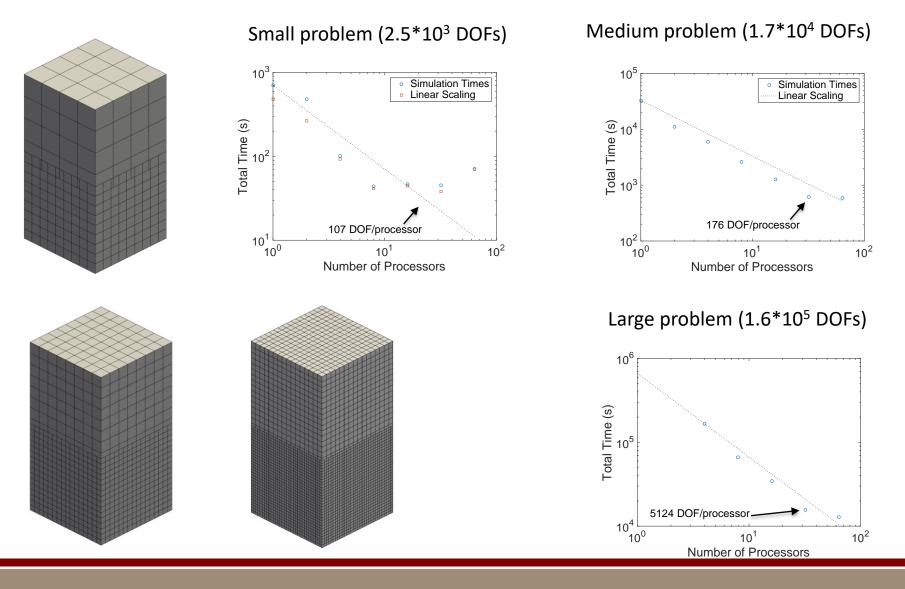
Goals:

- Connect *physical length scales* to *engineering scale models*.
- Investigate importance of *microstructural detail.*
- Develop bridging technologies for *spatial multiscale/ multiphysics*.

Appendix. Parallelization via DTK: Weak Scaling on Cubes Problem

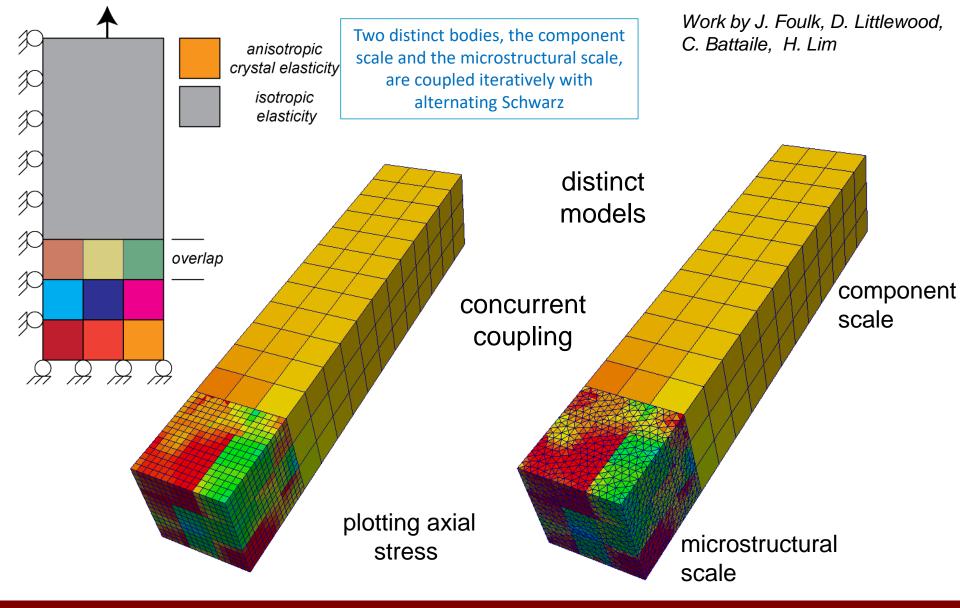


Appendix. Parallelization via DTK: Strong Scaling on Cubes Problem



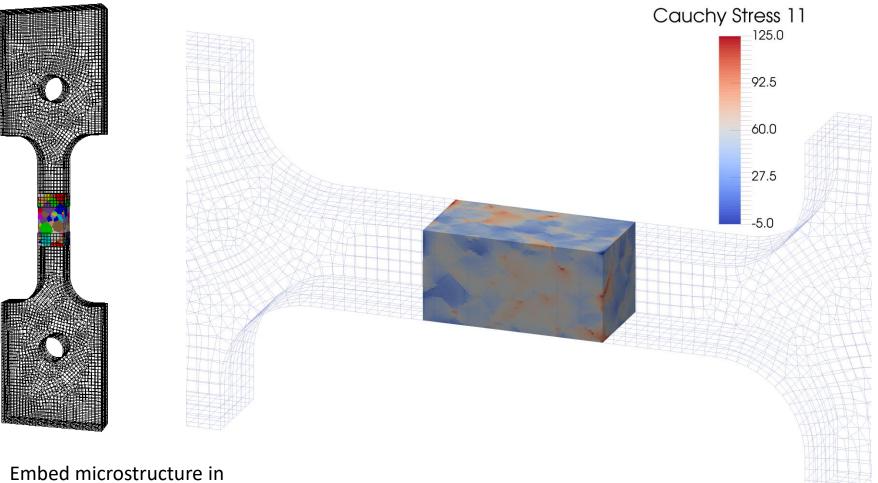
Appendix. Rubiks Cube Problem





Appendix. Tensile Bar





ASTM tensile geometry

Appendix. Tensile Bar: Meso-Macroscale

Mesoscale

SPARKS-generated microstructure (F. Abdeljawad)

cubic elastic constant : $C_{11} = 204.6$ GPa cubic elastic constant : $C_{12} = 137.7$ GPa cubic elastic constant : $C_{44} = 126.2$ GPa

reference shear rate : $\dot{\gamma}_0 = 1.0 \ 1/s$

rate sensitivity factor : m = 20

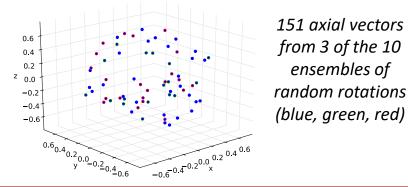
hardening rate parameter : $\dot{g}_0 = 2.0 \times 10^4 \text{ 1/s}$

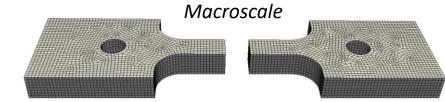
initial hardness : $g_0 = 90 \text{ MPa}$

saturation hardness : $g_s = 202 \text{ MPa}$

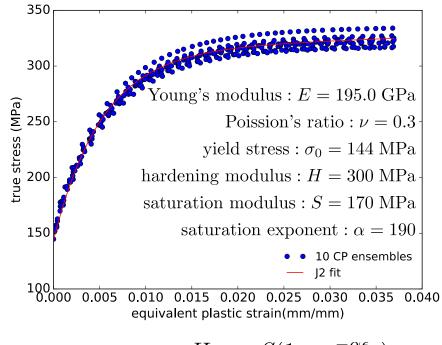
saturation exponent : $\omega = 0.01$

Fix microstructure, investigate ensembles





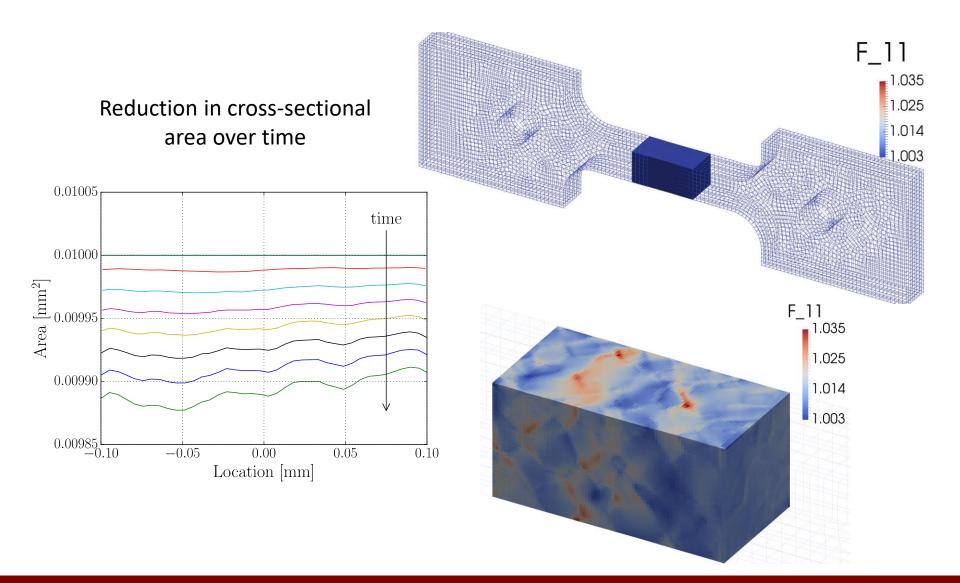
Load microstructural ensembles in uniaxial stress
 Fit flow curves with a macroscale J₂ plasticity model



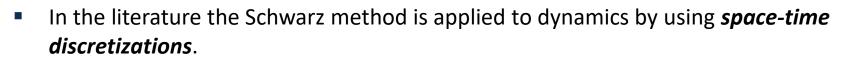
$$\sigma_y = \sigma_0 + H\epsilon_p + S(1 - e^{-\alpha\epsilon_p})$$

Appendix. Tensile Bar: Results

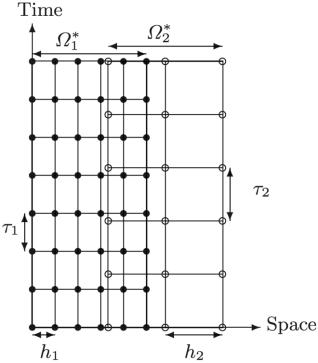




Appendix. Schwarz Alternating Method for Dynamics



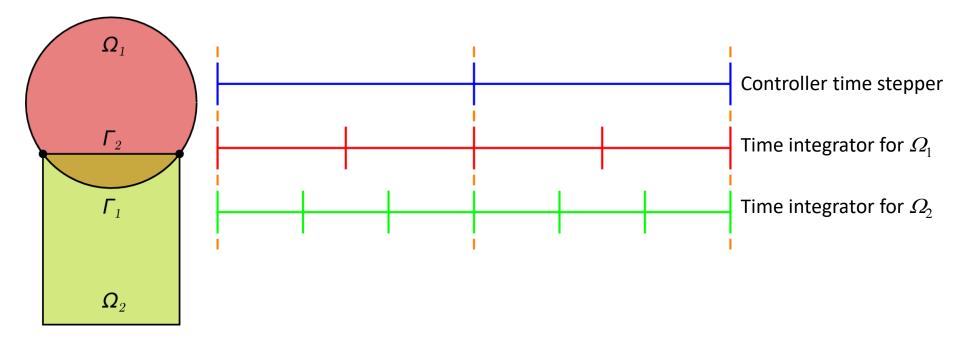
This was deemed *unfeasible* given the design of our current codes and size of simulations.



Overlapping non-matching meshes and time steps in dynamics.

Appendix. A Schwarz-like Time Integrator Distance

- We developed an *extension of Schwarz coupling* to *dynamics* using a governing time stepping algorithm that controls time integrators within each domain.
- Can use *different integrators* with *different time steps* within each domain.
- 1D results show *smooth coupling without numerical artifacts* such as spurious wave reflections at boundaries of coupled domains.



Appendix. Dynamic Singular Bar



- Inelasticity masks problems by introducing *energy dissipation*.
- Schwarz does not introduce numerical artifacts.
- Can couple domains with *different time integration schemes* (*Explicit-Implicit* below).

