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# **Peridynamic Modeling of the Failure of Heterogeneous Solids**

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# Need for a new theory of solid mechanics



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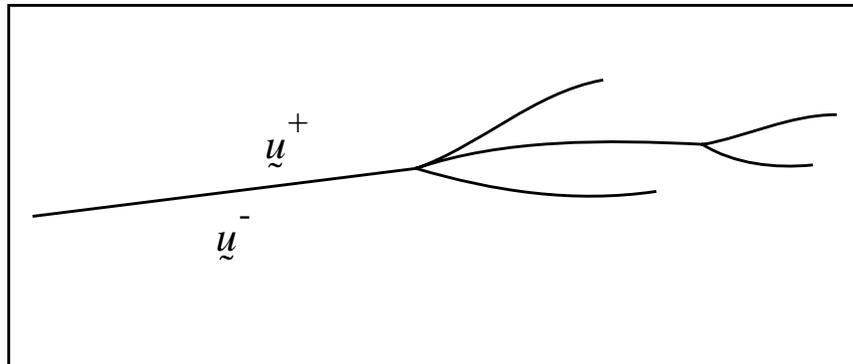
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- Classical formulation uses partial differential equations.
- The necessary spatial derivatives may not exist everywhere in the body.
- Example: Fracture ( $\underline{u}$  is discontinuous)



- Special techniques (of which there are many) are needed to model cracks in the classical theory.

## Goal

Develop a model in which exactly the same equations hold everywhere, regardless of any discontinuities.

- ◆ To do this, get rid of spatial derivatives.

# Basic idea of the peridynamic theory

- Equation of motion:

$$\rho \ddot{\underline{u}} = \underline{L}_u + \underline{b}$$

where  $\underline{L}_u$  is a functional.

- A useful special case:

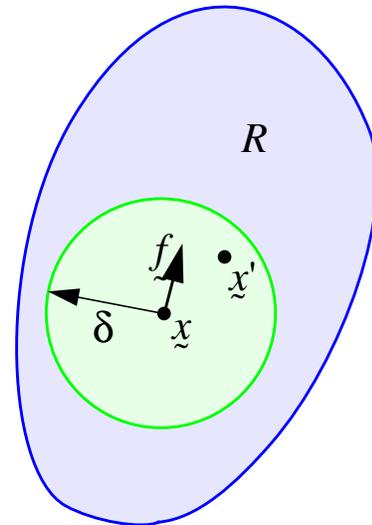
$$\underline{L}_u(\underline{x}, t) = \int_R \underline{f}(\underline{u}(\underline{x}', t) - \underline{u}(\underline{x}, t), \underline{x}' - \underline{x}) dV_{\underline{x}'}$$

where  $\underline{x}$  is any point in the reference configuration, and  $\underline{f}$  is a vector-valued function.

More concisely:

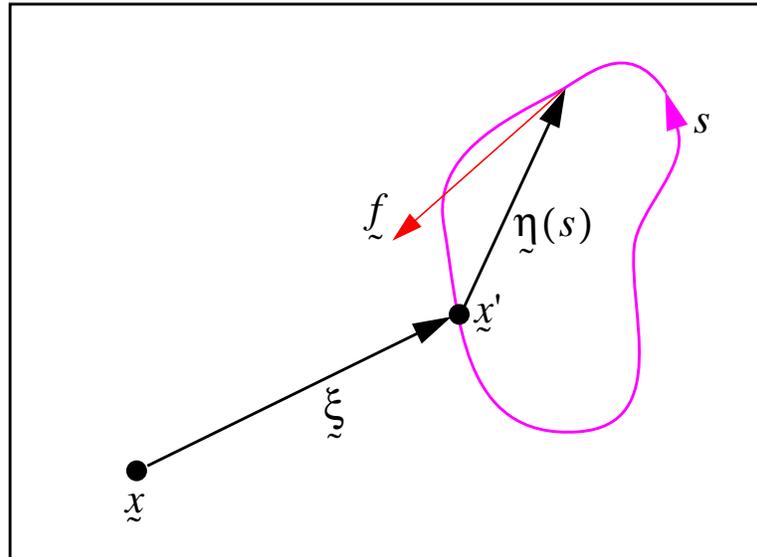
$$\underline{L} = \int_R \underline{f}(\underline{u}' - \underline{u}, \underline{x}' - \underline{x}) dV'.$$

- $\underline{f}$  is the pairwise force function. It contains all constitutive information.
- It is convenient to assume that  $\underline{f}$  vanishes outside some horizon  $\delta$ .



# Microelastic materials

A material is microelastic if, holding any  $\underline{x}$  fixed, the work done by  $\underline{f}$  in moving any  $\underline{x}'$  around a closed path is 0.



In this case, Stokes' Theorem implies:

- ◆ There exists a scalar-valued function  $w$ , called the micropotential, such that

$$\underline{f}(\underline{\eta}, \underline{\xi}) = \frac{\partial w}{\partial \underline{\eta}}(\underline{\eta}, \underline{\xi})$$

where

$$\underline{\xi} = \underline{x}' - \underline{x}$$

$$\underline{\eta} = \underline{u}' - \underline{u}$$

Can further show:

- ◆ There exists a scalar function  $H$  such that

$$\underline{f}(\underline{\eta}, \underline{\xi}) = (\underline{\xi} + \underline{\eta}) H(|\underline{\xi} + \underline{\eta}|, \underline{\xi})$$

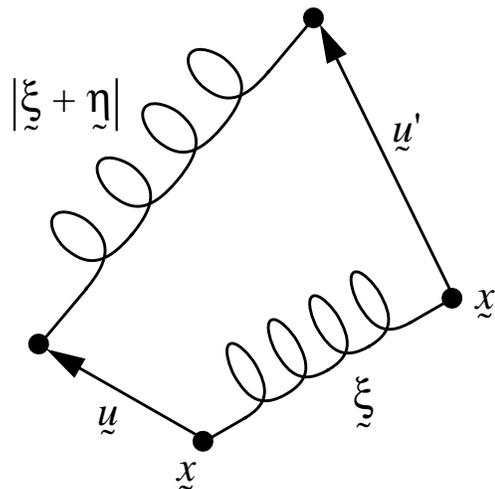
# Interpretation of microelasticity

So the micropotential can depend only on:

- the current separation distance  $|\xi + \eta|$
- the reference separation vector  $\xi$ .

Meaning: any two points  $\tilde{x}$  and  $\tilde{x}'$  are connected by a (possibly nonlinear) spring.

The spring properties can depend on the reference separation vector.



- Can prove: “microelastic implies macroelastic”:
  - Work done by external forces is stored in a recoverable form

# Some material models

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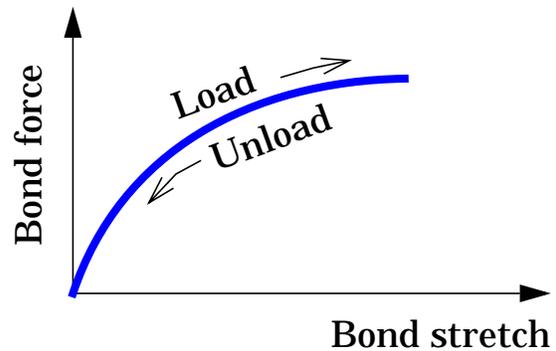
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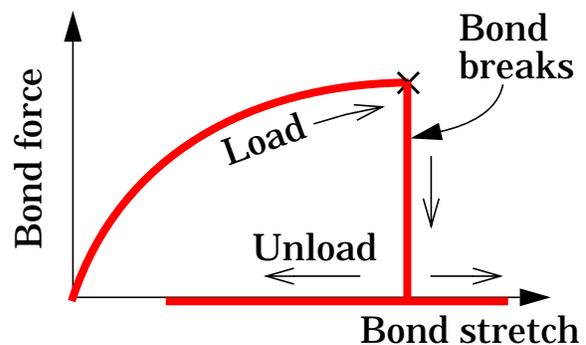
- Microelastic

- ◆ Each pair of particles is connected by a spring.
- ◆ Linear
- ◆ Bilinear
- ◆ etc.



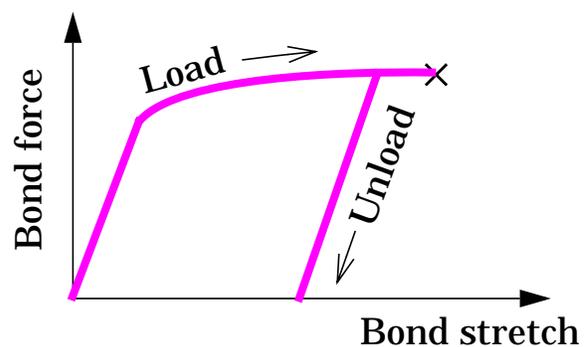
- Brittle microelastic

- ◆ Springs break irreversibly



- Microplastic

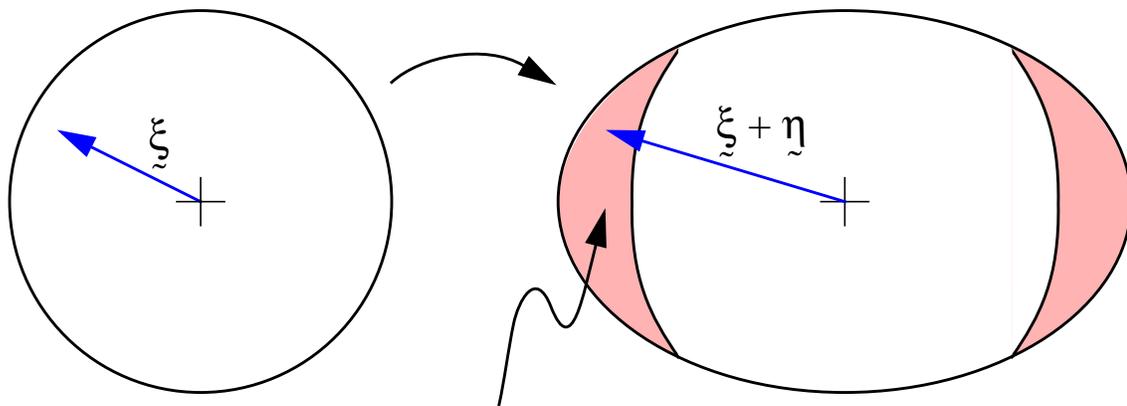
- ◆ Permanent bond deformation upon unloading.



- All of the above can have explicit rate dependence.

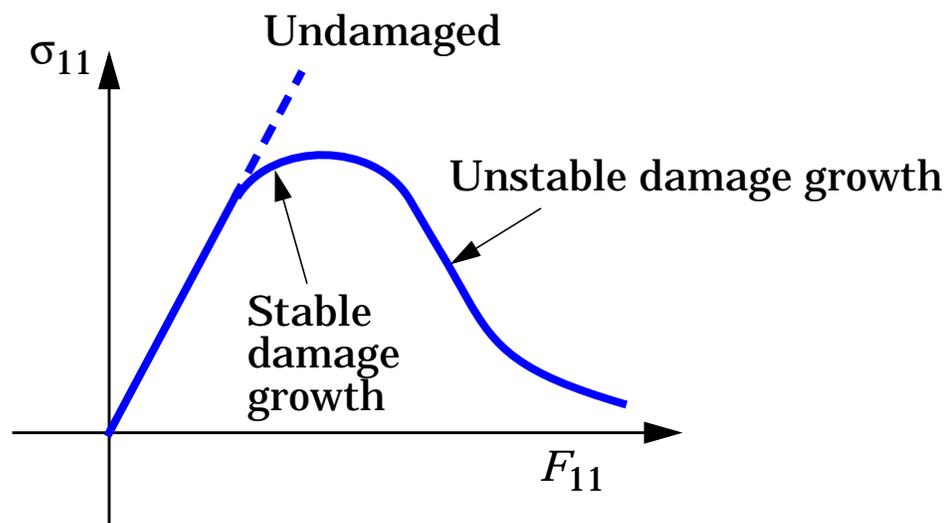
# How bond breakage leads to material fracture

- Continuum damage is caused by deformation:



$$\text{Broken springs: } |\xi + \eta| - |\xi| \geq \epsilon$$

- This causes a change in the “stress-strain” curve:



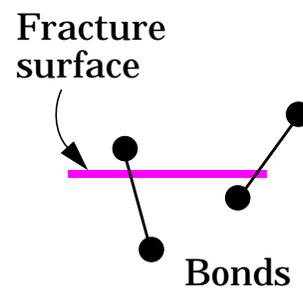
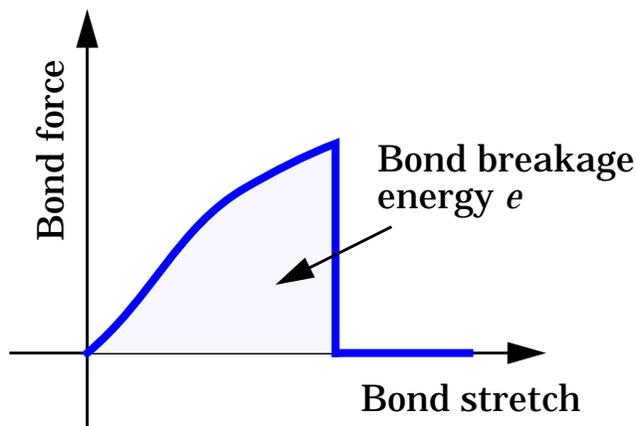
- Need to understand the mathematical conditions under which discontinuities can emerge.

# Determination of constitutive parameters



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- Linear microelastic:
  - Basic:
    - ◆ Spring constant is fit to wave speed data.
  - Advanced:
    - ◆ Can fit wave dispersion data if available.
    - ◆ Bond properties can depend on initial bond length.
- Microplastic:
  - Fit to uniaxial stress-strain curve.
- Bond breakage properties:
  - Fit breakage stretch to fracture toughness data.



Sum over bonds that are broken by the fracture:

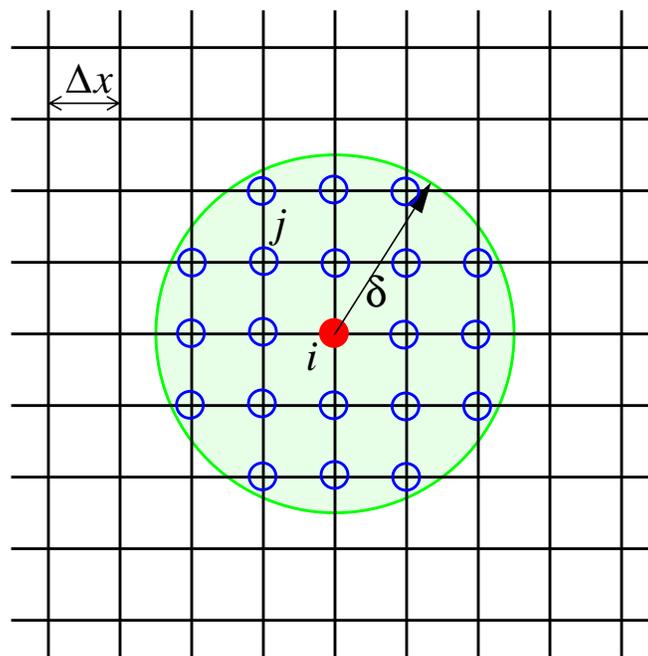
$$\text{Fracture energy} = \sum \text{bond breakage energy}$$

# Numerical solution method for dynamic problems



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- Theory lends itself to mesh-free numerical methods.
  - No elements.
  - Changing connectivity.
- Brute-force integration in space.



$$\rho \ddot{u}^i = \sum_{|x^j - x^i| < \delta} f(\underline{u}^j - \underline{u}^i, \underline{x}^j - \underline{x}^i) (\Delta x)^3$$

- Solution method has been found to scale well (almost linear speedup) when run on the Intel Teraflops computer at Sandia.
- Stable time step does not depend on mesh spacing (!)

# Peridynamic fracture model is “autonomous”



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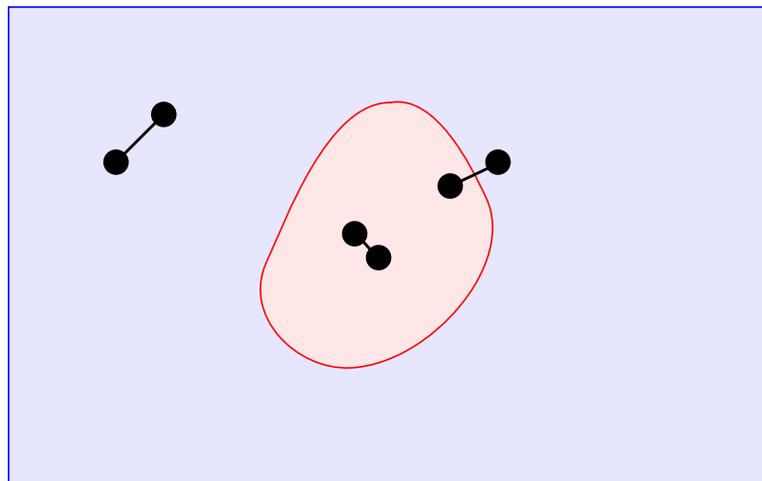
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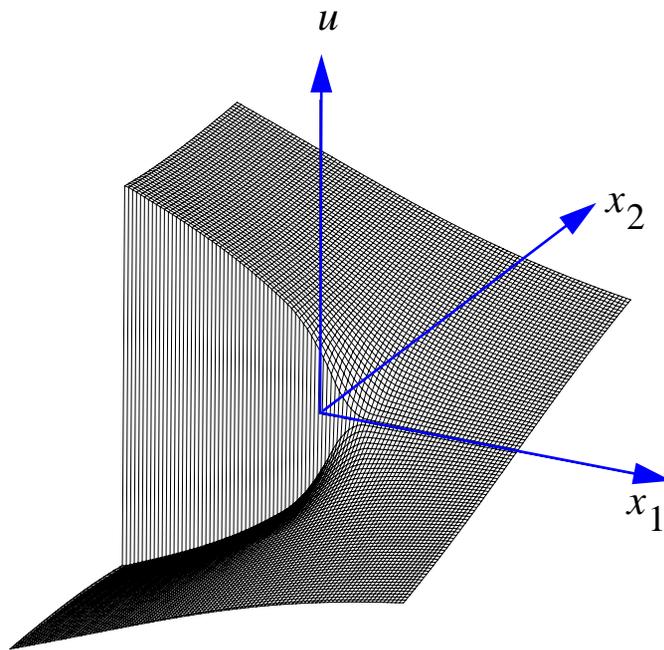
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- Cracks grow when and where it is energetically favorable for them to do so.
- Path, growth rate, arrest, branching, mutual interaction are predicted by the constitutive model and equation of motion (alone).
  - No need for any externally supplied relation controlling these things.
- Any number of cracks can occur and interact.
- Interfaces between materials have their own bond properties.

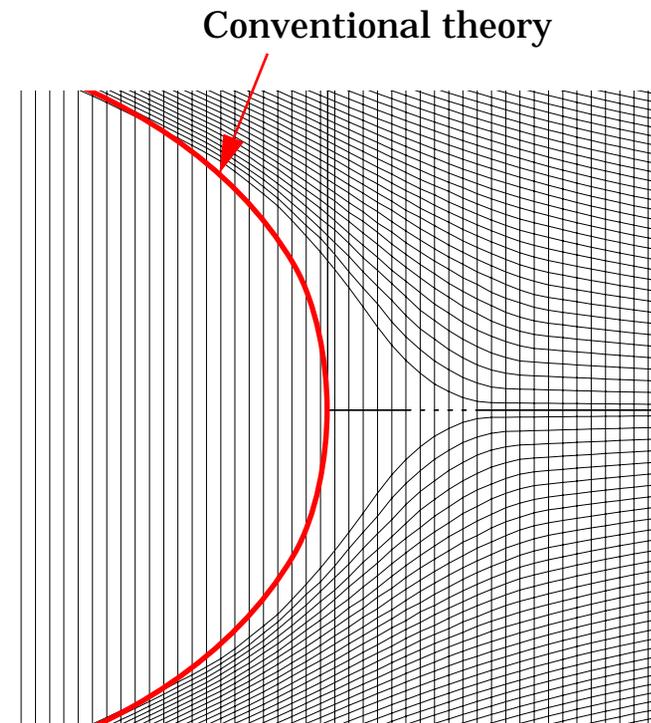


# Mode-III crack tip field

- Same equations are applied everywhere.
- Crack faces have cusp shape near tip.
  - No need for additional hypotheses (e.g. Barenblatt).



Surface into which a plane deforms

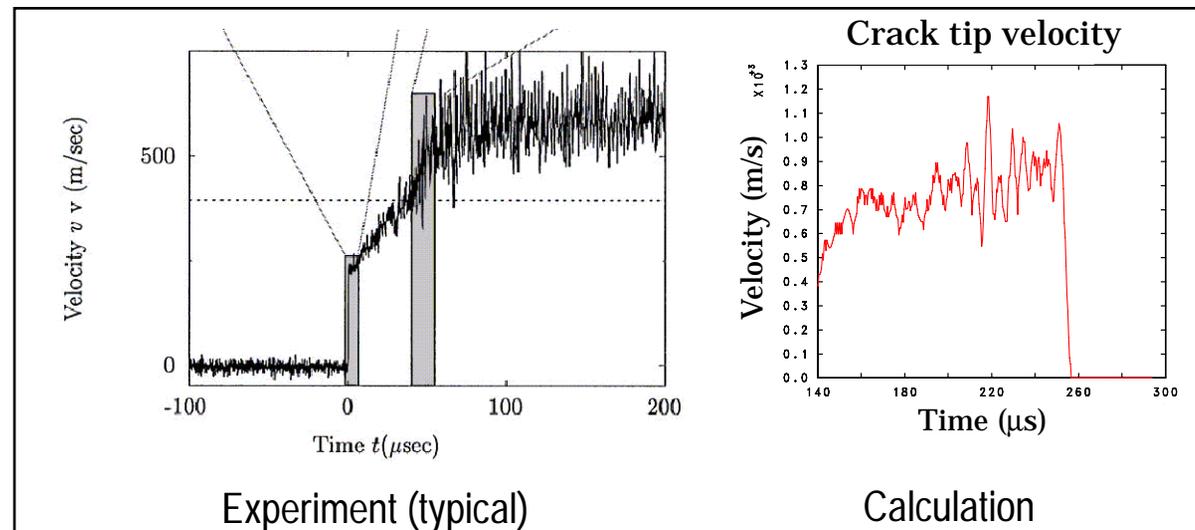
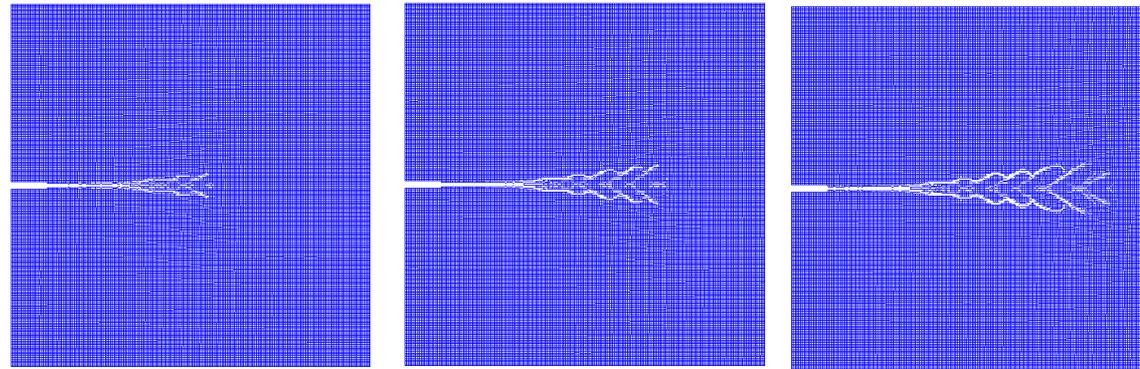
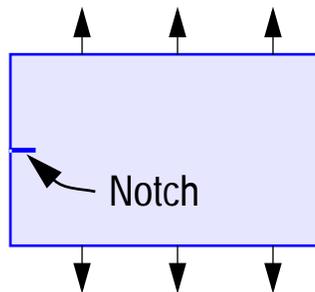


Enlarged view of crack tip

# Dynamic brittle fracture

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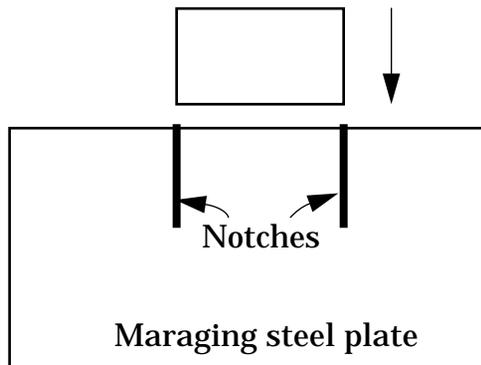
- Stretching of a PMMA plate<sup>1</sup>



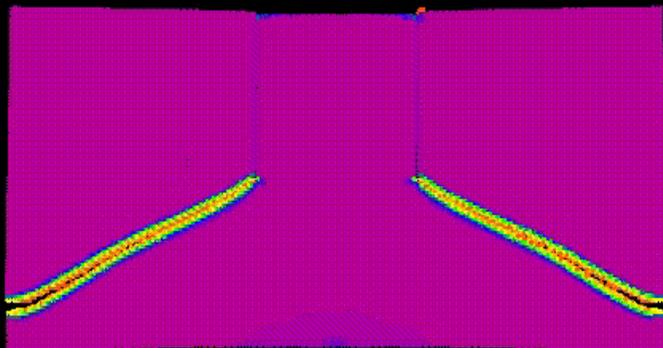
1. J. Fineberg and M. Marder, *Physics Reports* 313 (1999) 1-108

# Dynamic fracture in a tough steel: Kalthoff-Winkler experiment

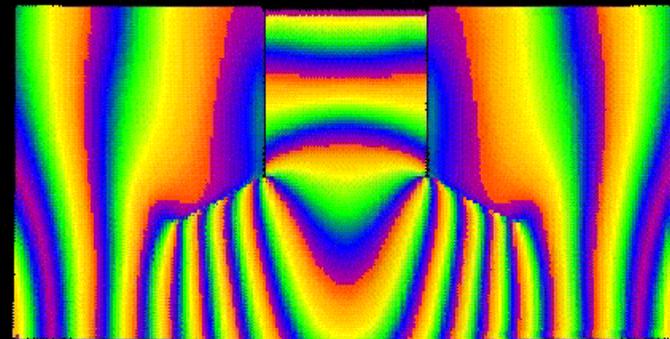
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- Code predicts correct crack angles.



Crack paths

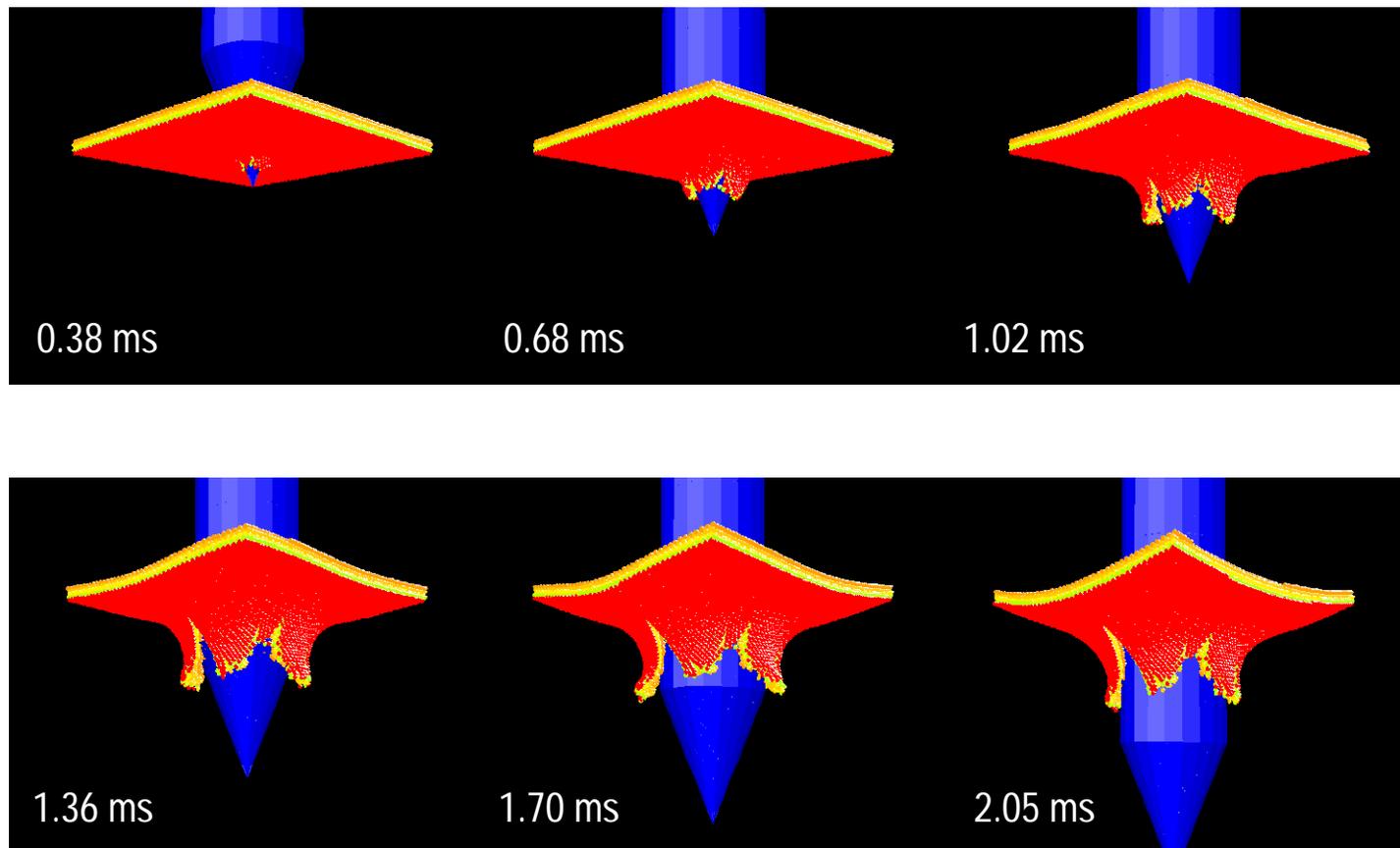


Simulated Moire fringes

# Perforation of thin ductile targets

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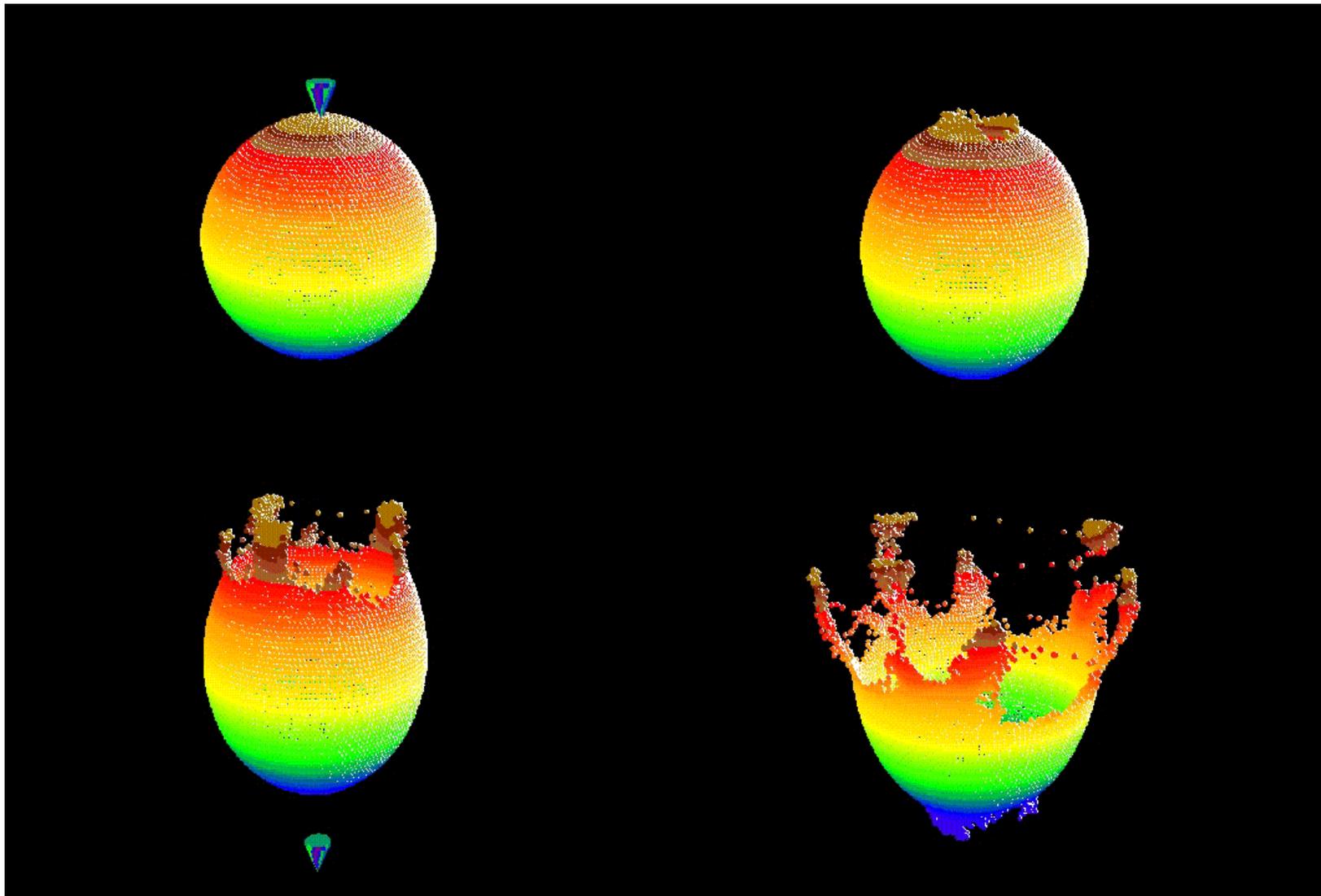
- Peak force occurs at about 0.4ms (end of drilling phase):<sup>1</sup>



1. Not all of the target is shown.

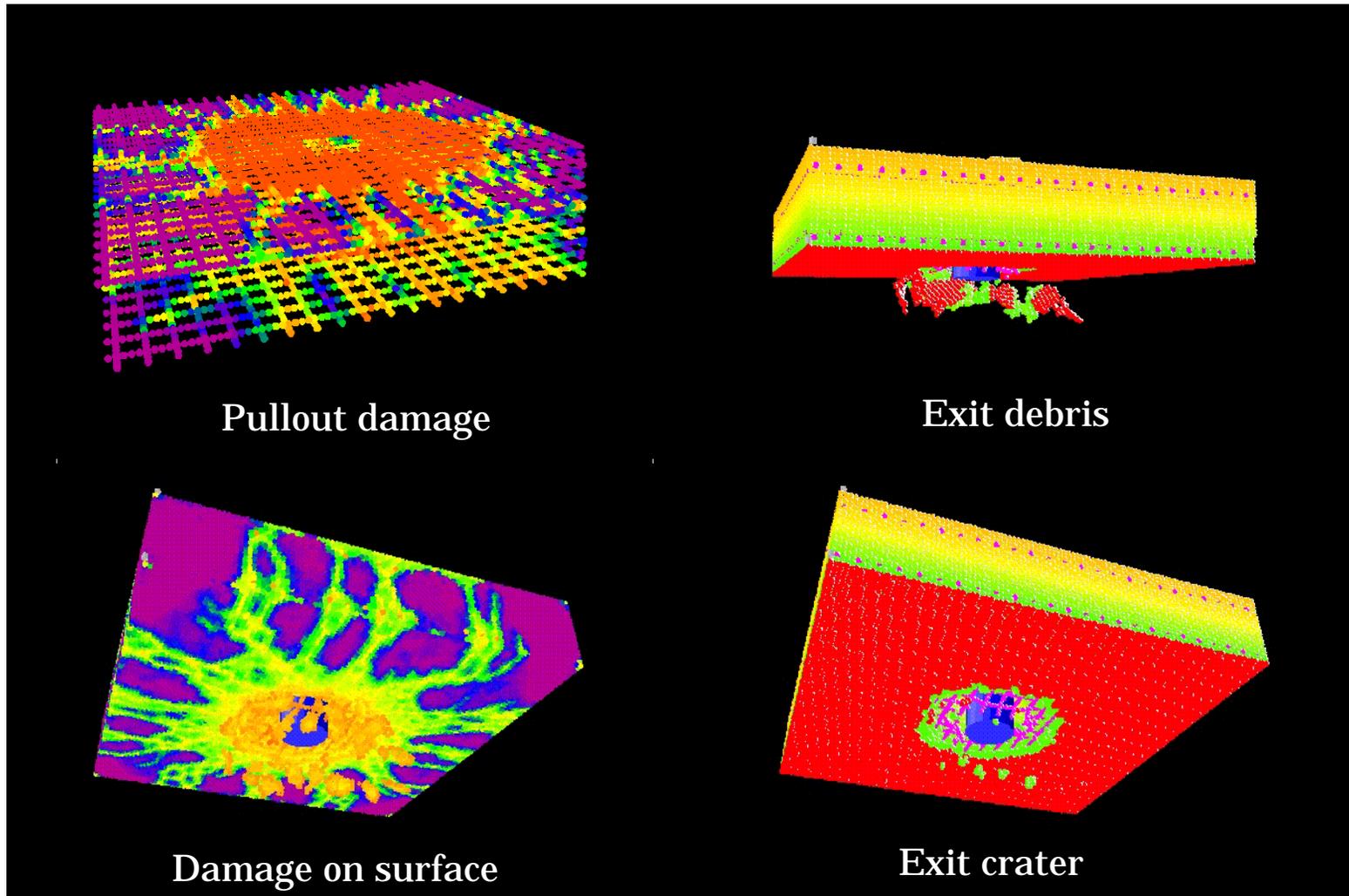
# Dynamic fracture in a balloon

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# Nonhomogeneous materials: Perforation of reinforced concrete

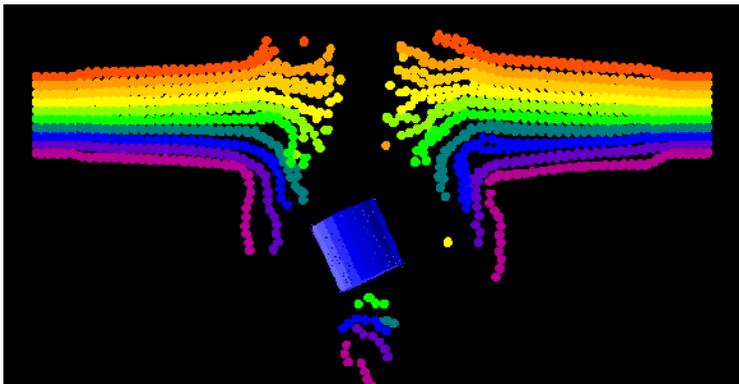
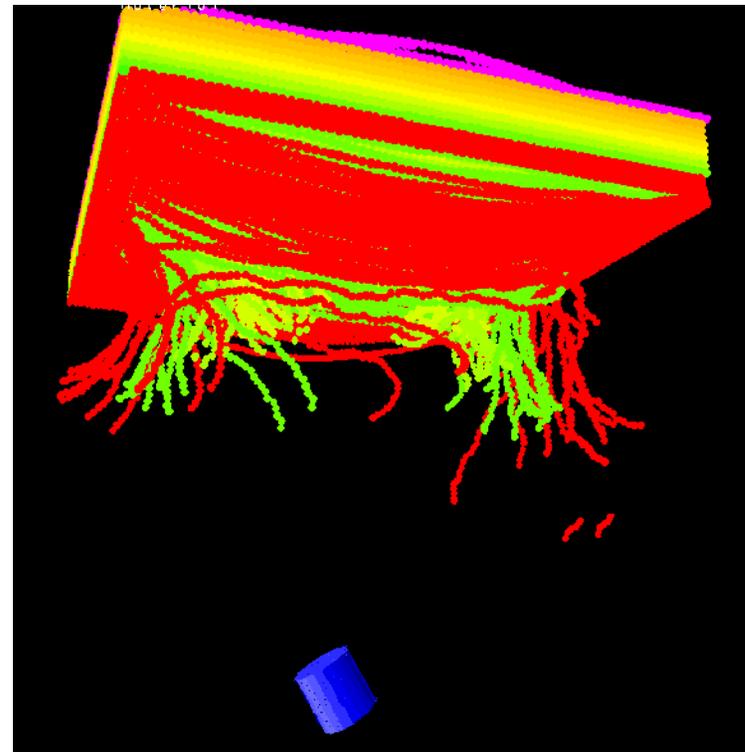
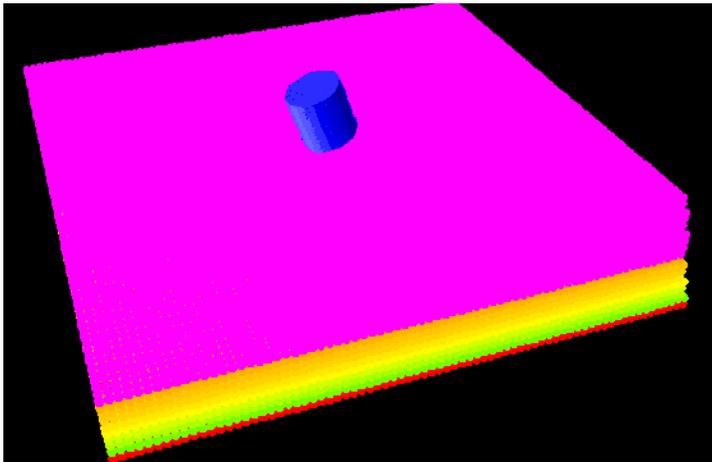
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# Nonhomogeneous materials: Perforation of a fabric

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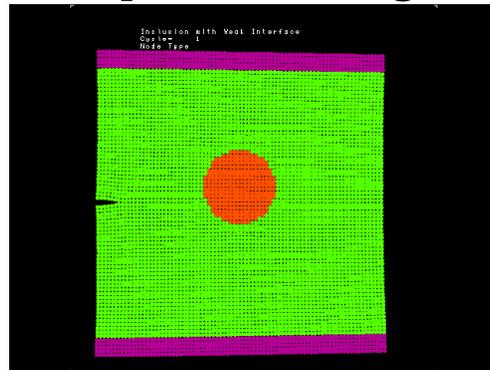
- Model includes fiber breakage, contact, and adhesion:



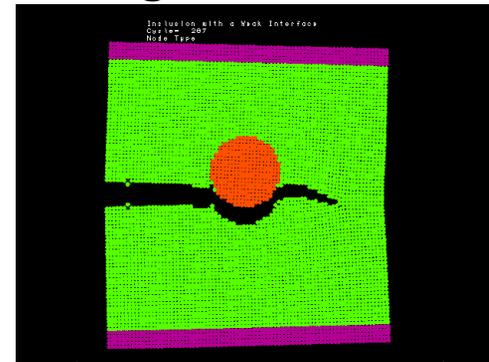
# Nonhomogeneous materials: Fracture in a composite unit cell

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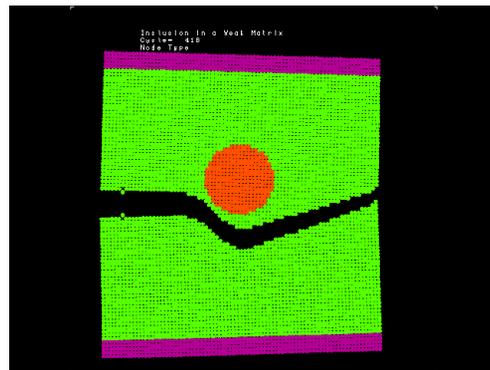
- Crack path, growth, and stability depend only on material properties.
- No need for separate laws governing crack growth.



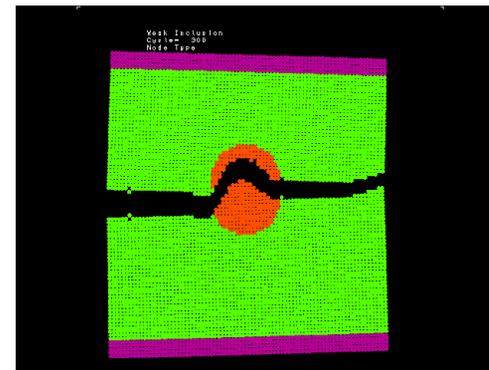
Initial condition



Weak interface



Weak matrix

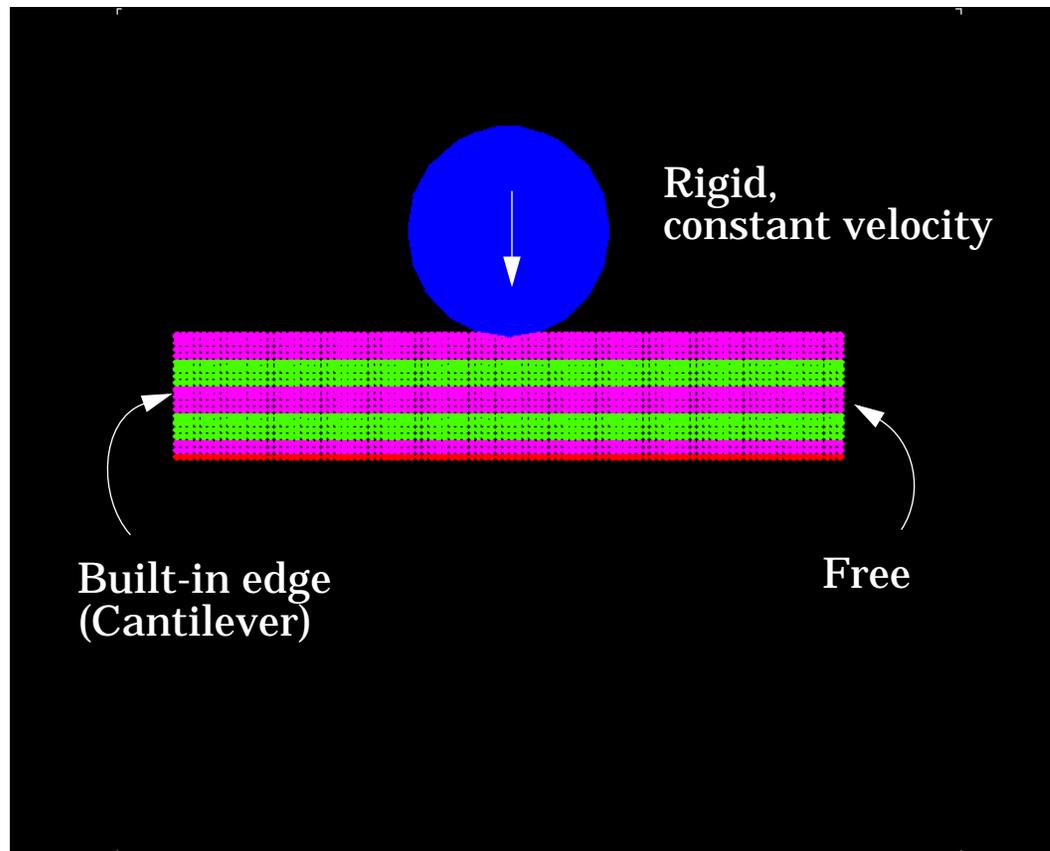


Weak fiber

# Layered material example: Where does failure first occur?

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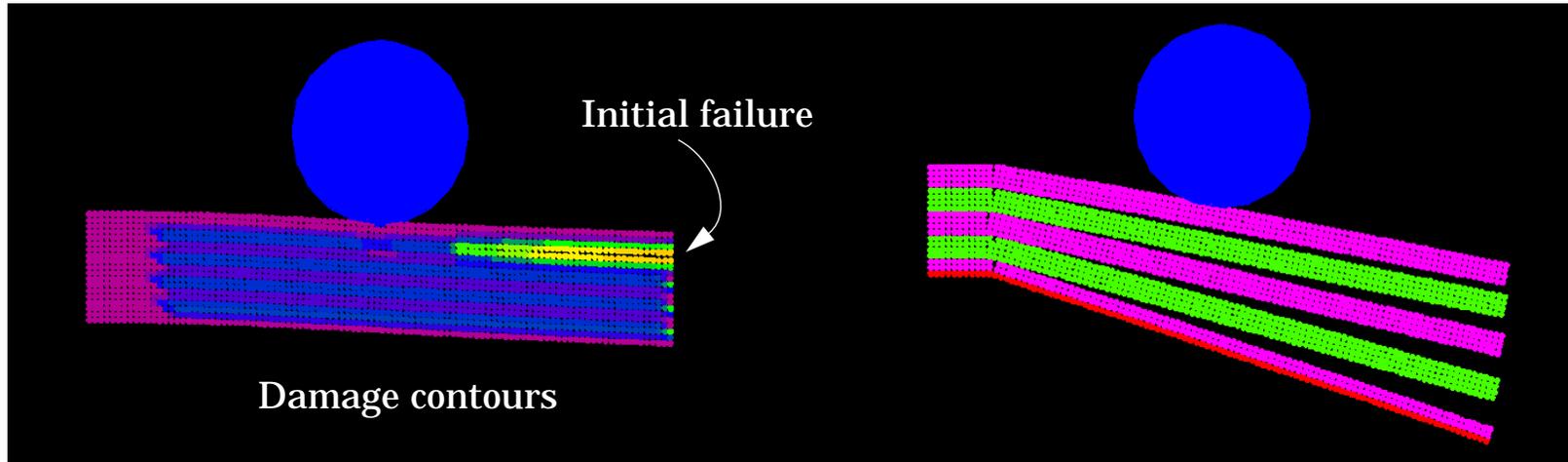
- Layers have identical properties.
- Interfaces have half the strength of the layers.



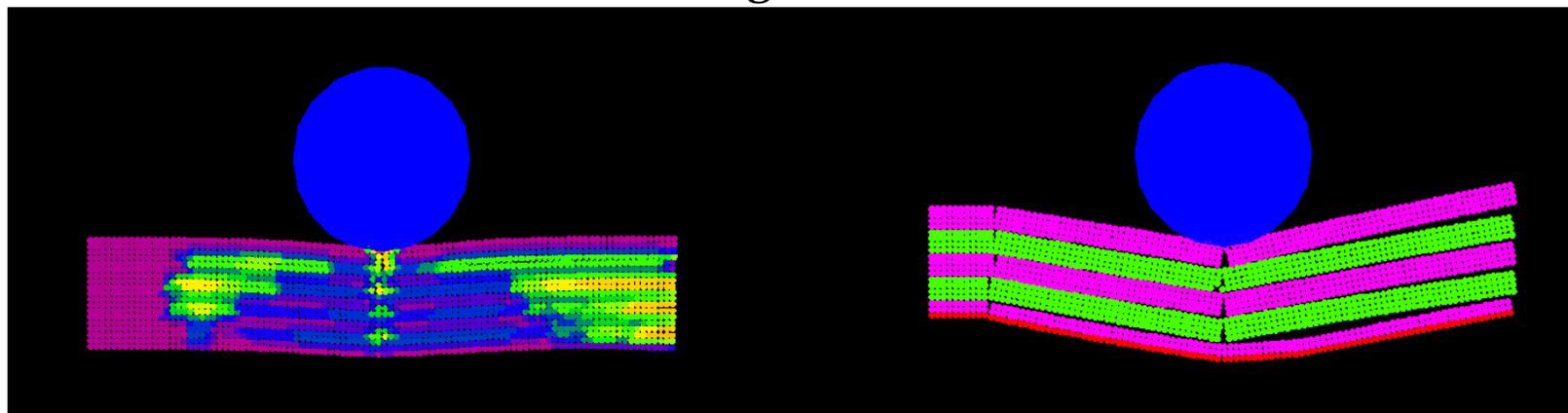
# Initial failure site and mode depends on loading rate

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Low rate

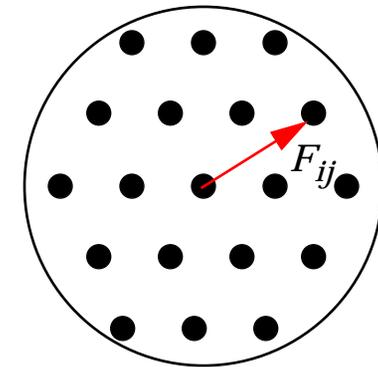


High rate



# Correspondence with atomic-scale physics

- Can a constitutive model be derived rigorously from an atomic-scale physical description?
  - Classical theory: people have been trying for a long time.
  - Peridynamic theory may be a more natural way to do this because of its similarity to molecular dynamics.
    - ◆ This is currently being attempted by Bhattacharya (Caltech) and Abeyaratne (MIT).
  - May lead to a good way to do multiscale modeling.



# Conclusions



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- Method appears to have the potential to model:
  - Heterogeneous materials of great complexity.
  - Complex fracture systems without the need to keep track of each crack.

### **Possible research directions**

- Mechanics of heterogeneous materials
  - Understand how failure progresses from one material to another.
  - Improved material models.
  - Validation against interface crack data.
  - Fatigue cracks.
  - Multiscale modeling.
  - Learn how to do complex material systems.
- Theory and numerical solutions
  - Improved solvers.
    - ◆ Multigrid, iterative, implicit, etc.